

Supplementary Materials for

Plasmonic nanoreactors regulating selective oxidation by energetic electrons and nanoconfined thermal fields

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Section S1. Experimental Section

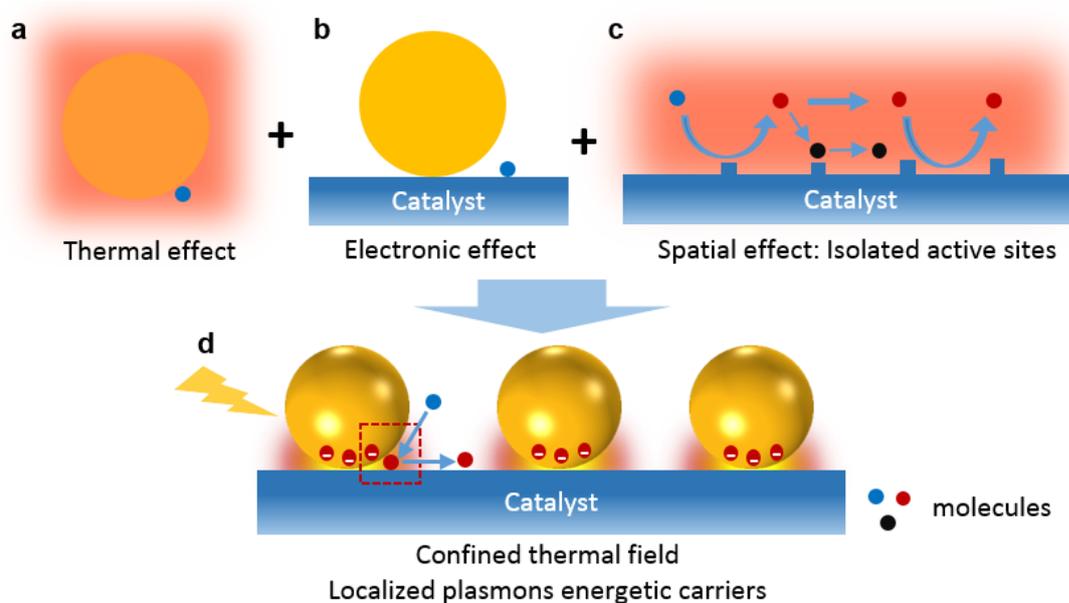


Fig. S1. Summary of plasmonic effects that may be used to improve catalytic performance. a, Thermal effect, which helps overcome an activation barrier. b, Electronic effect, which can directly influence the catalytic process. c, Spatial effect, which can isolate the active site to influence the reaction process. d, SPs are consistent with the design of hierarchical structures, thereby adding new attributes to existing catalysts so as to take advantage of all or most of the three effects a, b and c, described above. And the spatial distributions of temperature and energetic charge carriers induced by SPs is non-uniform and more local, on the nanoscale. The energetic charge carriers yielded by SPs can directly transfer to the reactants to activate them.

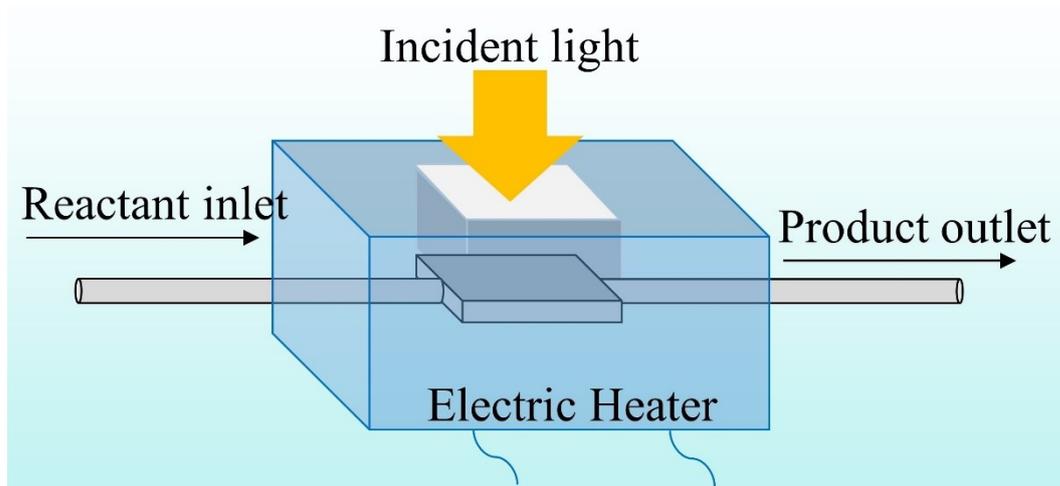


Fig. S2. Schematic of set-up and quartz glass reactor.

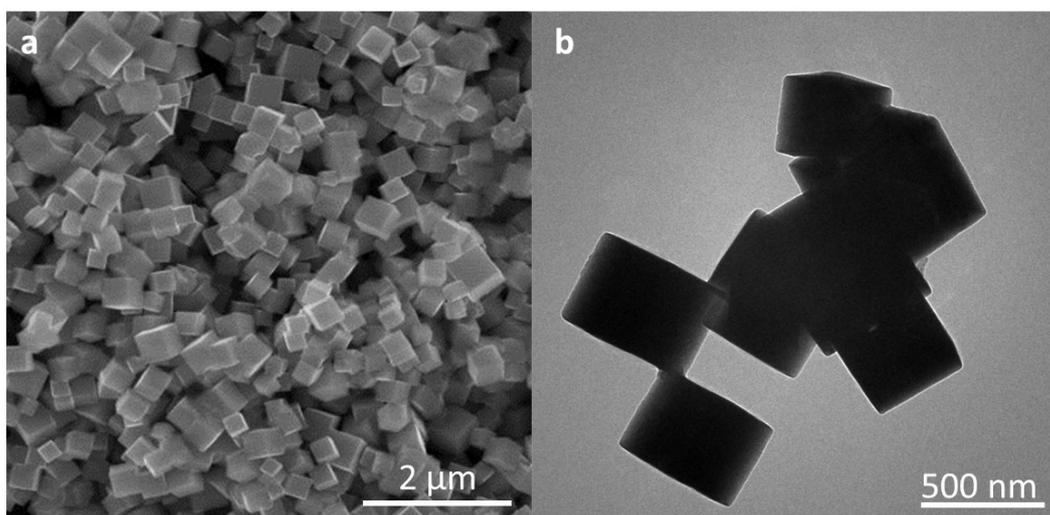


Fig. S3. SEM and TEM images of as-prepared cube Cu_2O . The images well prove the cube Cu_2O we prepared are uniform and crystallized with the (100) crystal plane exposure.

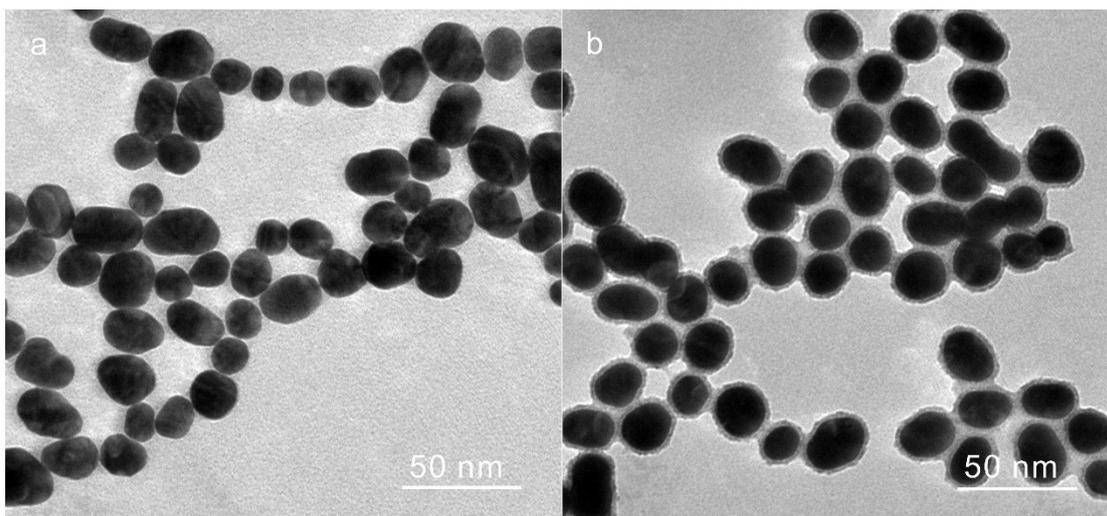


Fig. S4. TEM images of as-prepared 30 nm Au nanoparticles (a) and Au@SiO₂ nanoparticles (b). The SiO₂ shell is about 5 nm.

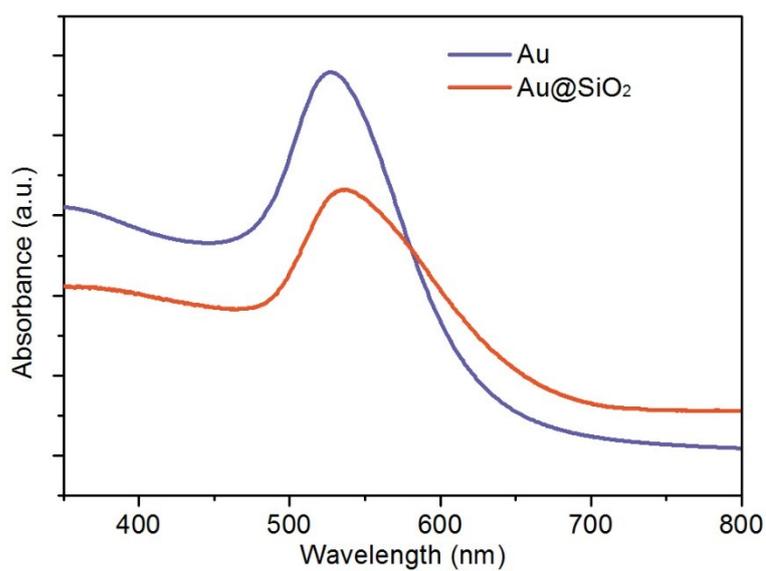


Fig. S5. Absorption spectrum of as-prepared 30 nm Au nanoparticles and Au@SiO₂ nanoparticles.

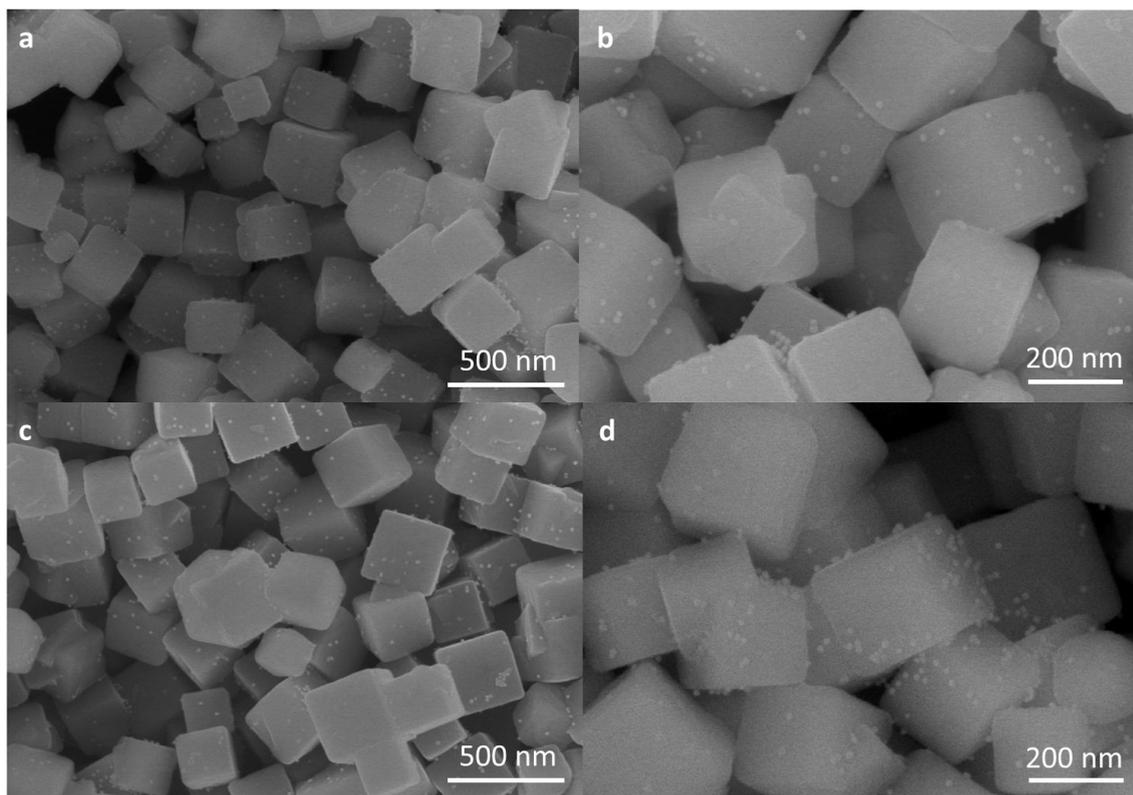


Fig. S6. SEM images of as-prepared Au-Cu₂O (a, b) and Au@SiO₂-Cu₂O (c, d) hierarchical catalysts.

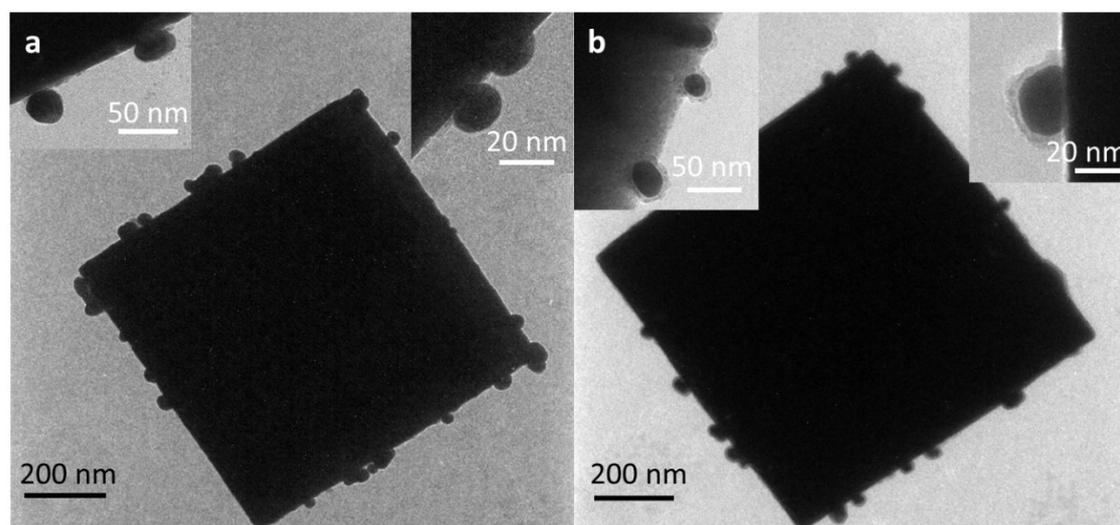


Fig. S7. TEM images of as prepared Au-Cu₂O (a) and Au@SiO₂-Cu₂O (b) hierarchical catalysts.

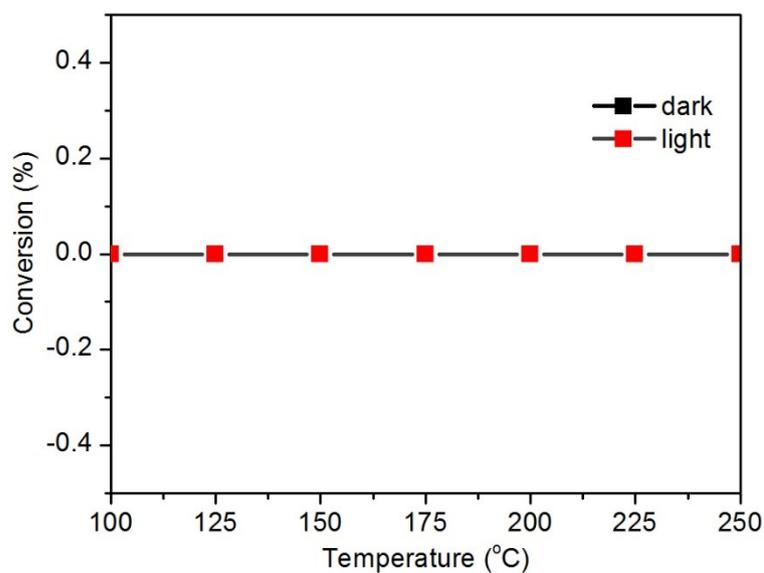


Fig. S8. Catalytic performance of Au-SiO₂ with or without illumination at different temperature. There is no obvious reaction between 100 °C to 250 °C.

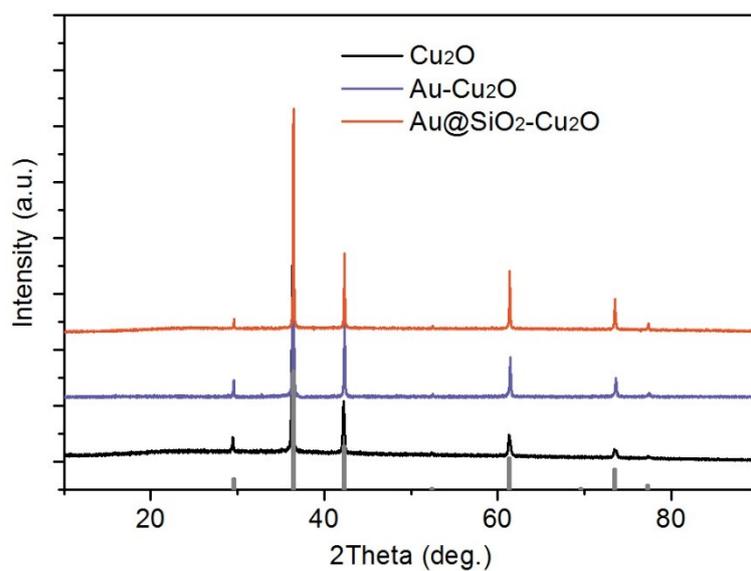


Fig. S9. XRD patterns of as prepared cube Cu₂O, Au-Cu₂O and Au@SiO₂-Cu₂O catalysts. The assemble of Au NPs or Au@SiO₂ NPs on the surface of cube Cu₂O does not affect its crystal phase.

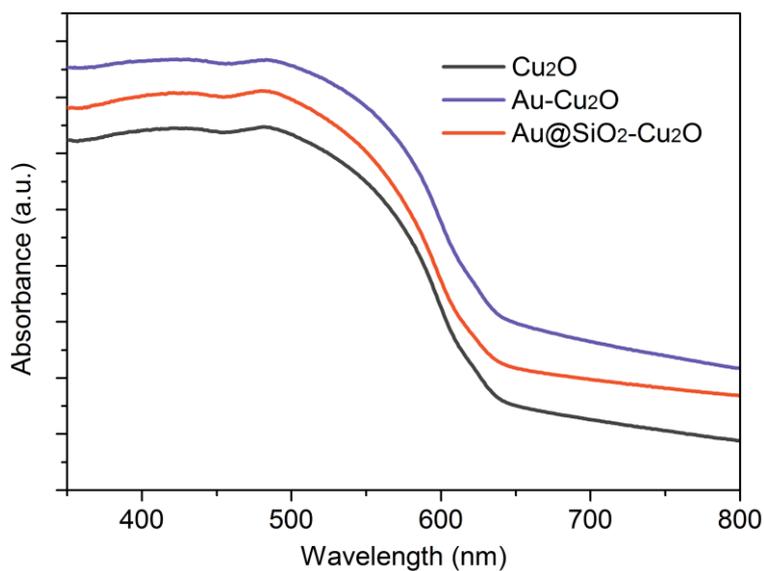


Fig. S10. Absorption spectrum of as prepared cube Cu₂O, Au-Cu₂O and Au@SiO₂-Cu₂O hierarchical catalysts.

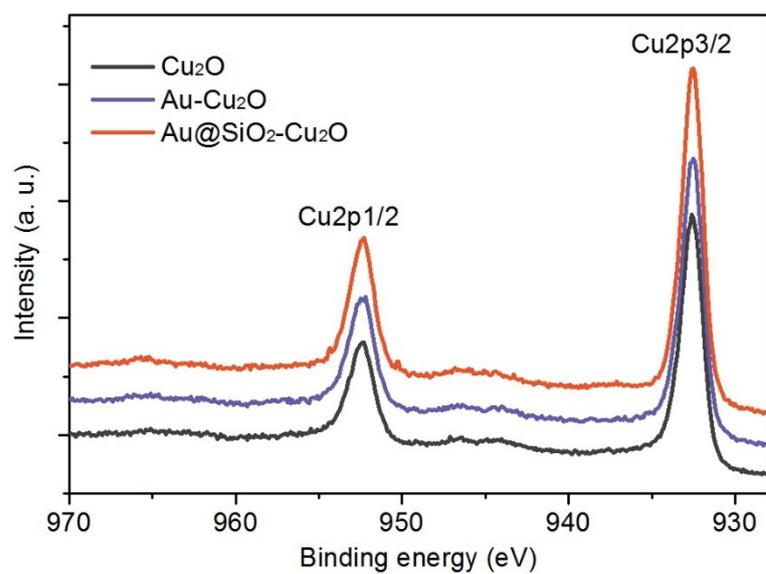


Fig. S11. XPS of Cu 2p of C-Cu₂O, Au-Cu₂O and Au@SiO₂-Cu₂O hierarchical catalysts. The Cu 2p XPS spectrum with no satellite structure to the main peaks indicates no oxidation of surface Cu⁺ (54).

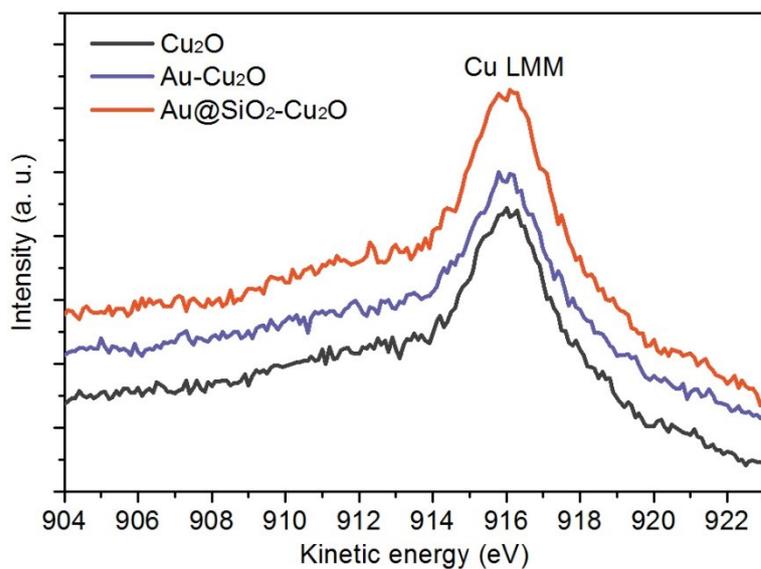


Fig. S12. XPS results of Cu LMM of C-Cu₂O, Au-Cu₂O and Au@SiO₂-Cu₂O hierarchical catalysts. Cu LMM Auger transition kinetic energy at 916.5 eV further improves the existence of Cu⁺ on the surface (54, 55).

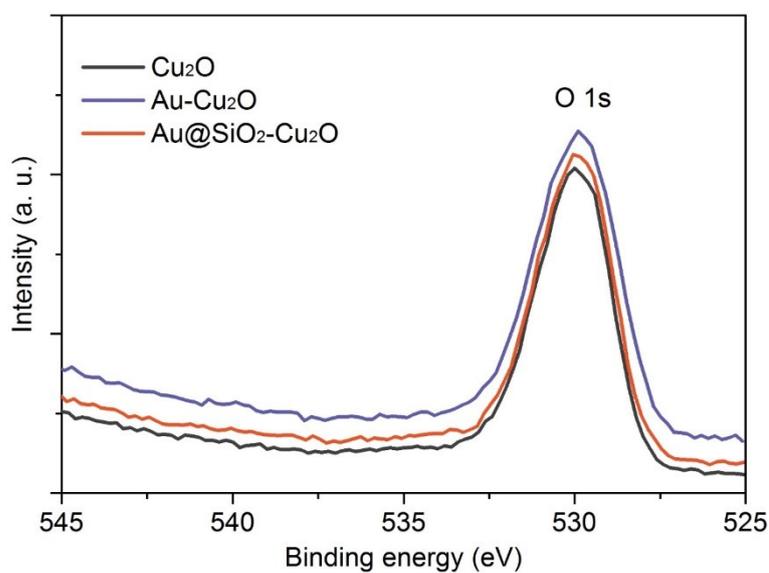


Fig. S13. XPS results of O 1s of C-Cu₂O, Au-Cu₂O and Au@SiO₂-Cu₂O hierarchical catalysts.

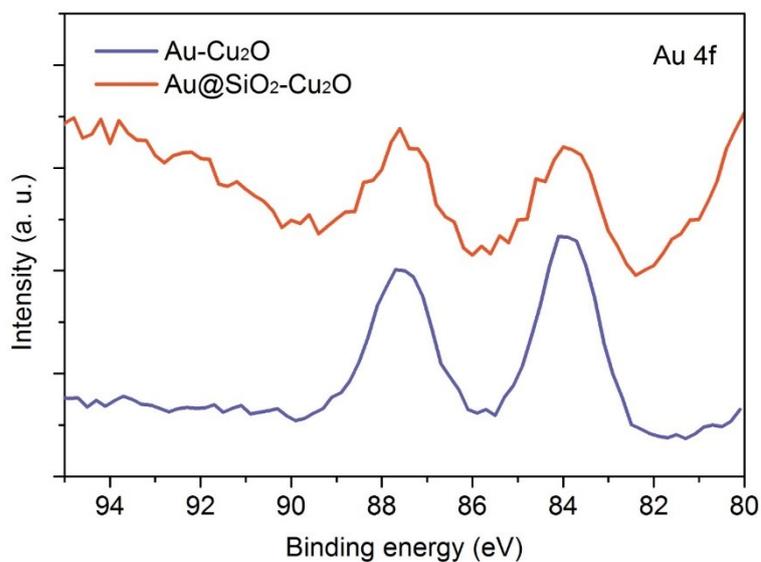


Fig. S14. XPS results of Au 4f of Au-Cu₂O and Au@SiO₂-Cu₂O hierarchical catalysts (56-57).

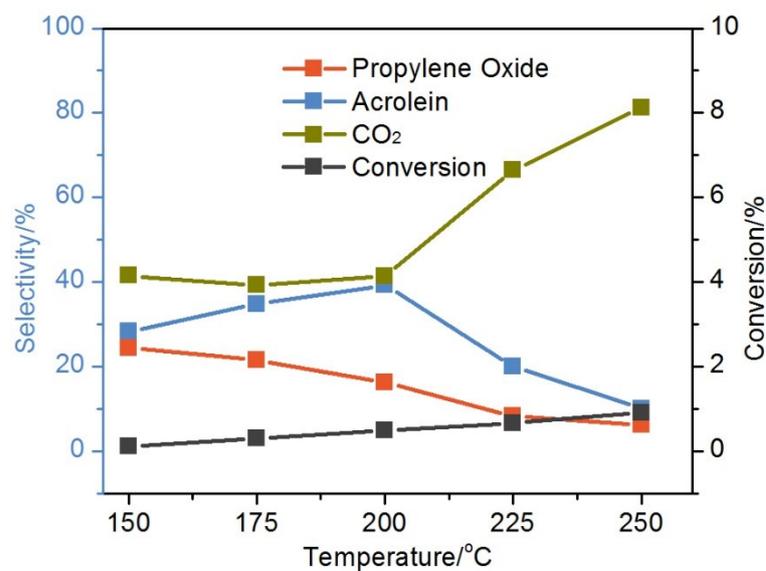


Fig. S15. Catalytic performance of surfactant free C-Cu₂O. The main product is CO₂.

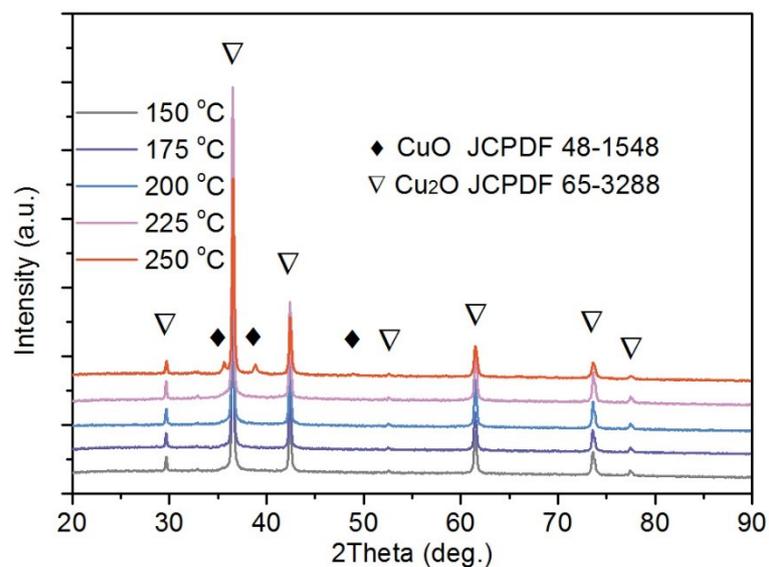


Fig. S16. XRD patterns and of C-Cu₂O after catalysis at different temperature: (a) 150 °C, (b) 175 °C, (c) 200 °C, (d) 225 °C, (e) 250 °C. No obvious CuO peak was observed until the operated temperature was higher than 250 °C.

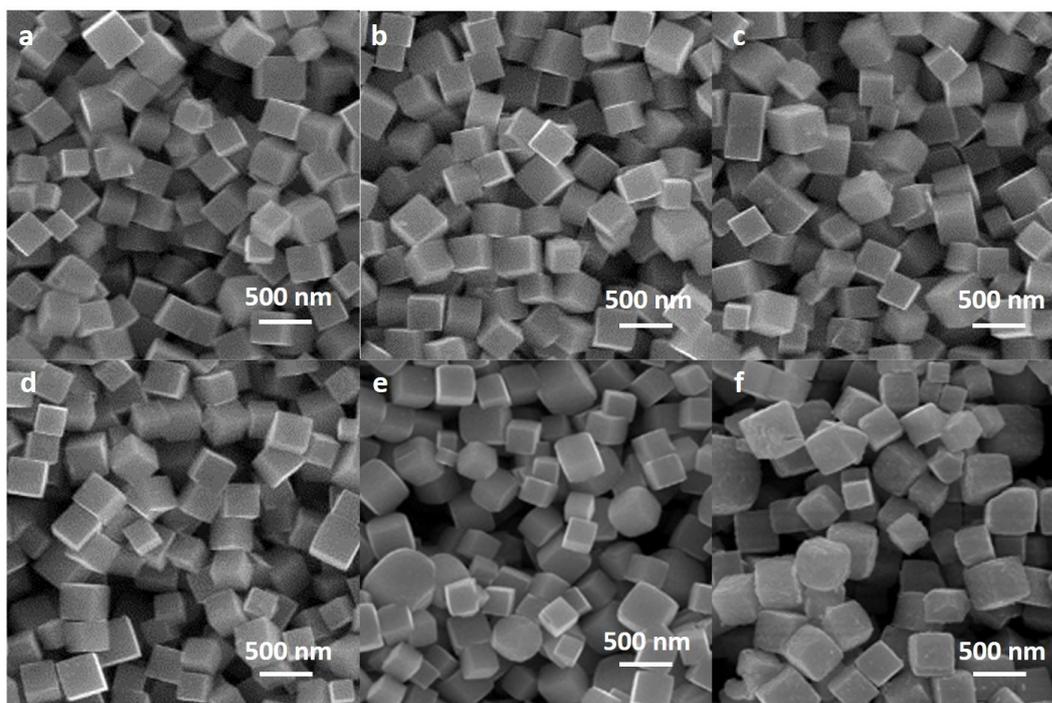


Fig. S17. SEM images of C-Cu₂O after catalysis at different temperature: (a) 125 °C, (b) 150 °C, (c) 175 °C, (d) 200 °C, (e) 225 °C, (f) 250 °C.

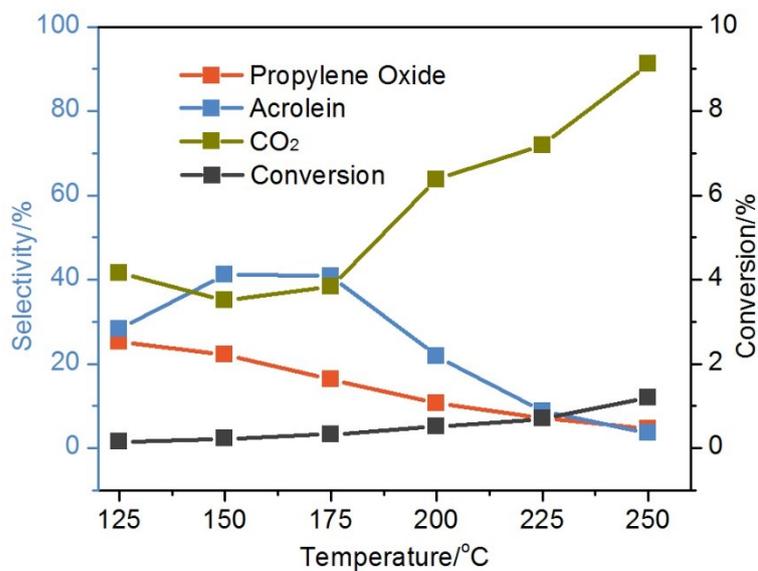


Fig. S18. Catalytic performance of surfactant free C-Cu₂O under illumination at different temperature. The main product is CO₂.

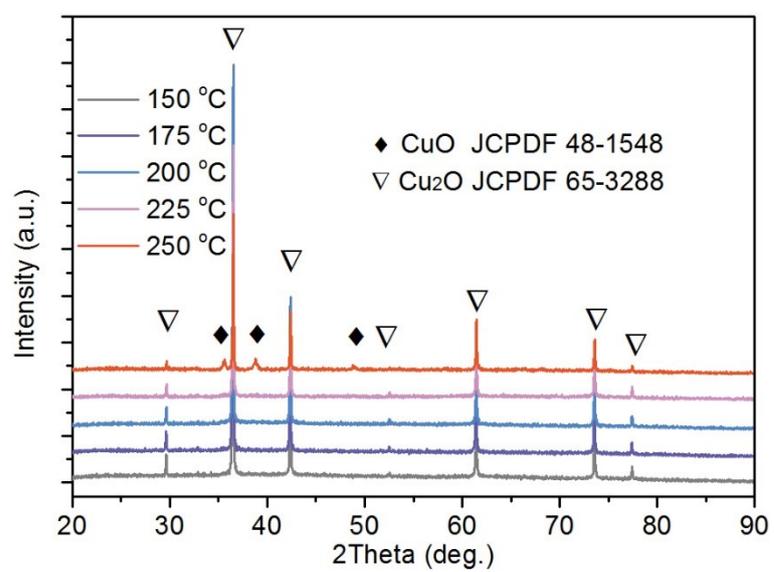


Fig. S19. XRD patterns and of C-Cu₂O after catalysis under illumination at different temperature: (a) 150 °C, (b) 175 °C, (c) 200 °C, (d) 225 °C, (e) 250 °C. No obvious CuO peak was observed until the operated temperature was higher than 250 °C.

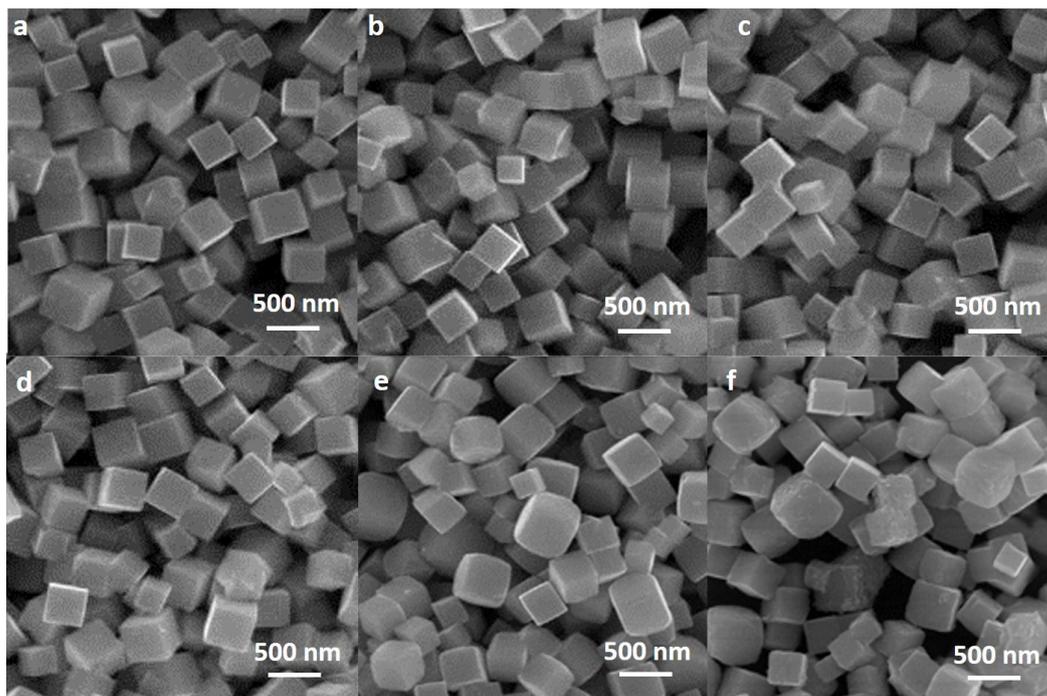


Fig. S20. SEM images of C-Cu₂O after catalysis under illumination at different temperature: (a) 125 °C, (b) 150 °C, (c) 175 °C, (d) 200 °C, (e) 225 °C, (f) 250 °C.

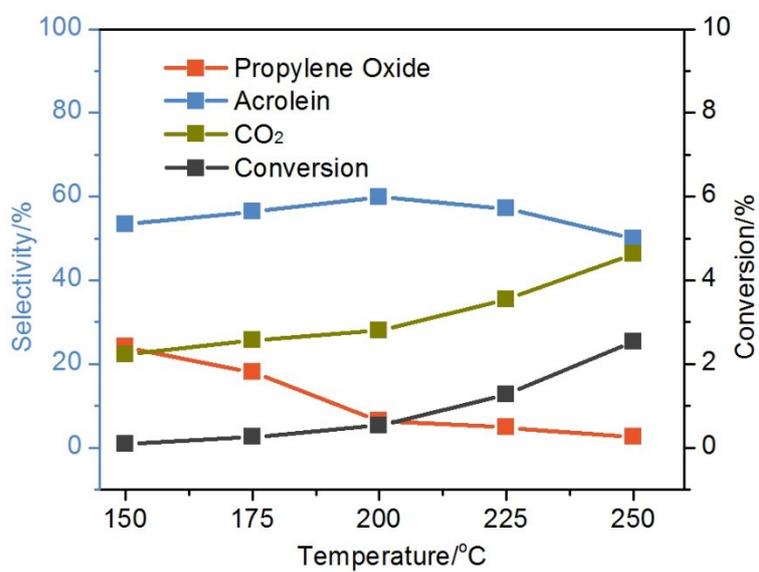


Fig. S21. Catalytic performance of Au-Cu₂O at different temperature. The main product is acrolein.

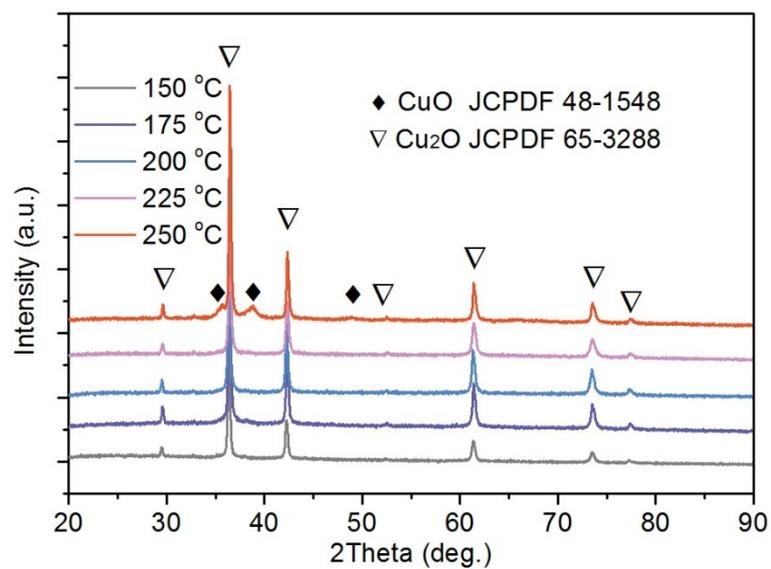


Fig. S22. XRD patterns and of Au-Cu₂O after catalysis at different temperature: (a) 150 °C, (b) 175 °C, (c) 200 °C, (d) 225 °C, (e) 250 °C. No obvious CuO peak was observed until the operated temperature was higher than 250 °C.

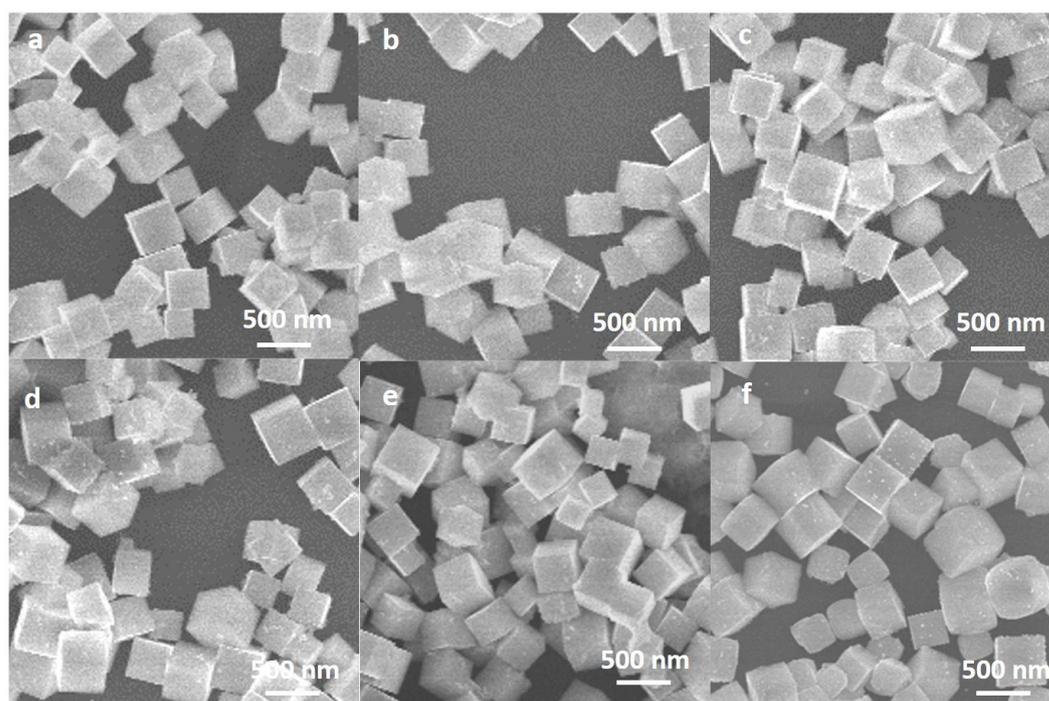


Fig. S23. SEM images of Au-Cu₂O after catalysis at different temperature: (a) 125 °C, (b) 150 °C, (c) 175 °C, (d) 200 °C, (e) 225 °C, (f) 250 °C.

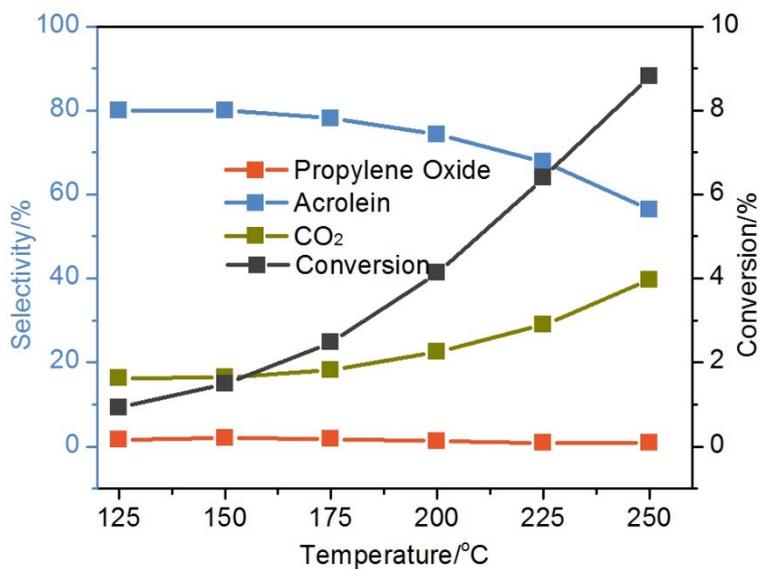


Fig. S24. Catalytic performance of Au-Cu₂O under illumination at different temperature. The main product is acrolein.

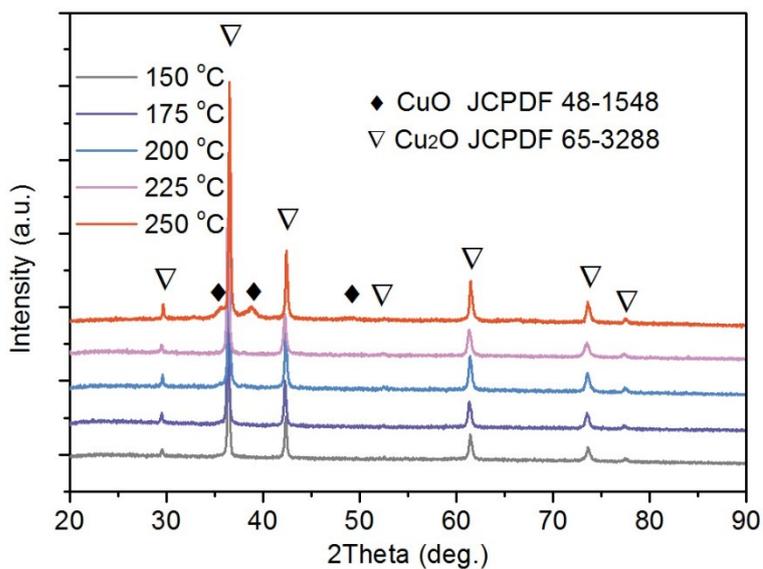


Fig. S25. XRD patterns and of Au-Cu₂O after catalysis under illumination at different temperature: (a) 150 °C, (b) 175 °C, (c) 200 °C, (d) 225 °C, (e) 250 °C. No obvious CuO peak was observed until the operated temperature was higher than 250 °C.

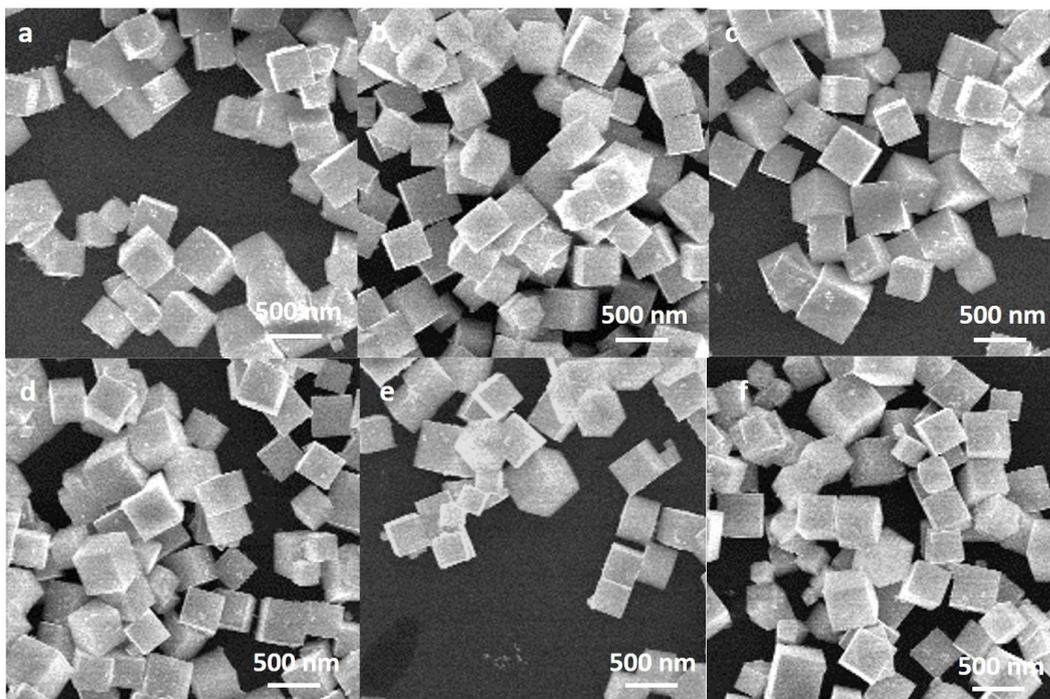


Fig. S26. SEM images of Au-Cu₂O after catalysis under illumination at different temperature: (a) 125 °C, (b) 150 °C, (c) 175 °C, (d) 200 °C, (e) 225 °C, (f) 250 °C.

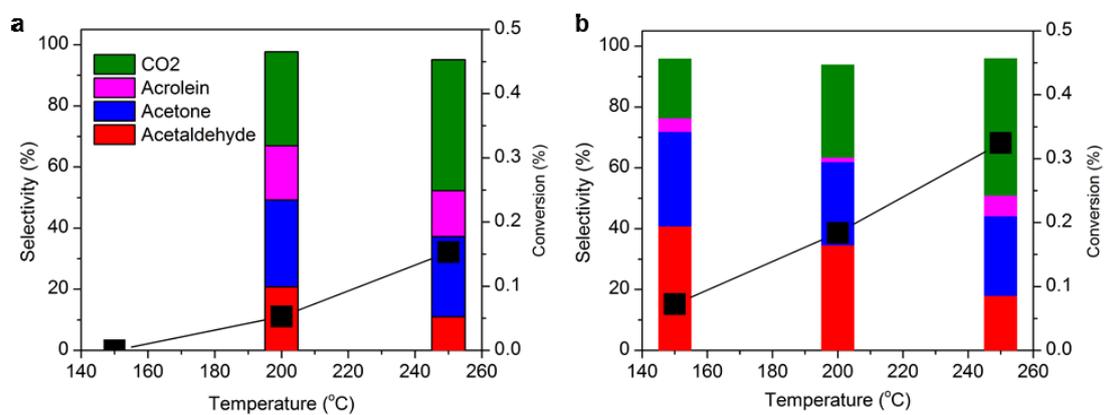


Fig. S27. Catalytic performance of Au-TiO₂ without (a) and with (b) illumination of visible light.

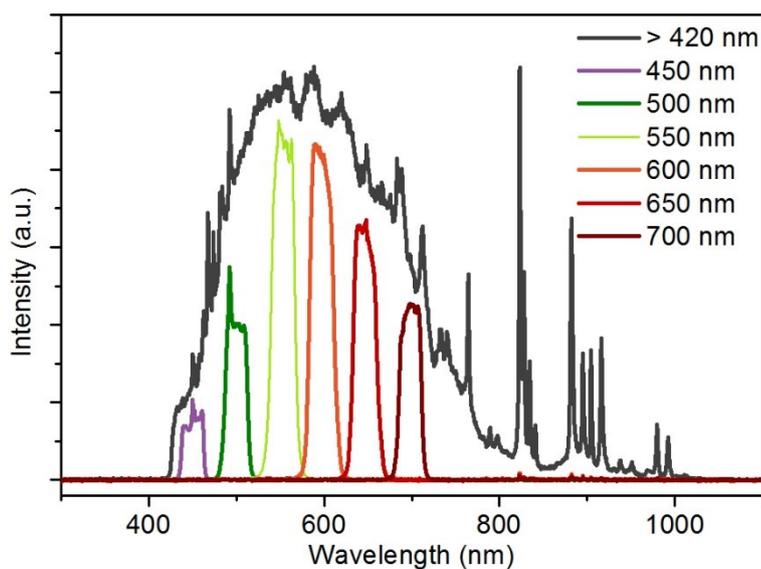


Fig. S28. Spectrum of the light source with 420 nm long-pass filter, and 450 nm, 500 nm, 550 nm, 600 nm, 650 nm, 700 nm band-pass filters (58).

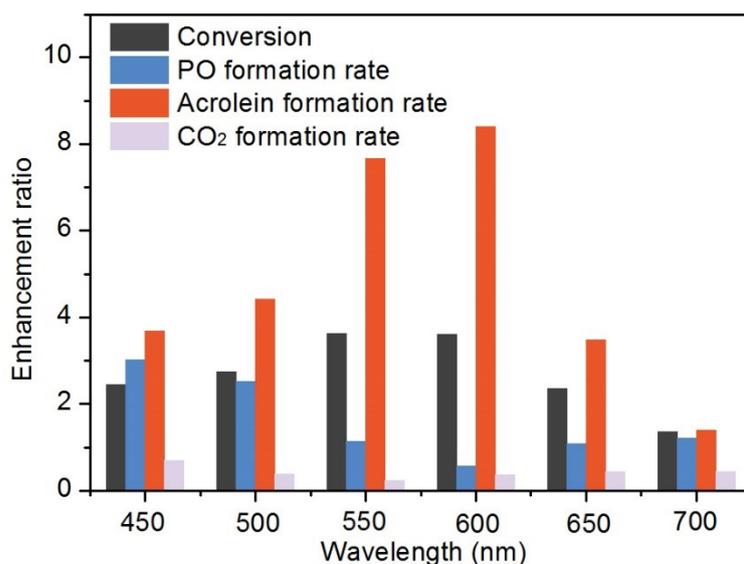


Fig. S29. The enhancements of the propylene conversion and the product formation rate for Au-Cu₂O hierarchical structure at 150 °C as a function of incident light wavelengths. They are calculated by dividing the conversion or formation rate with illumination by that without illumination.

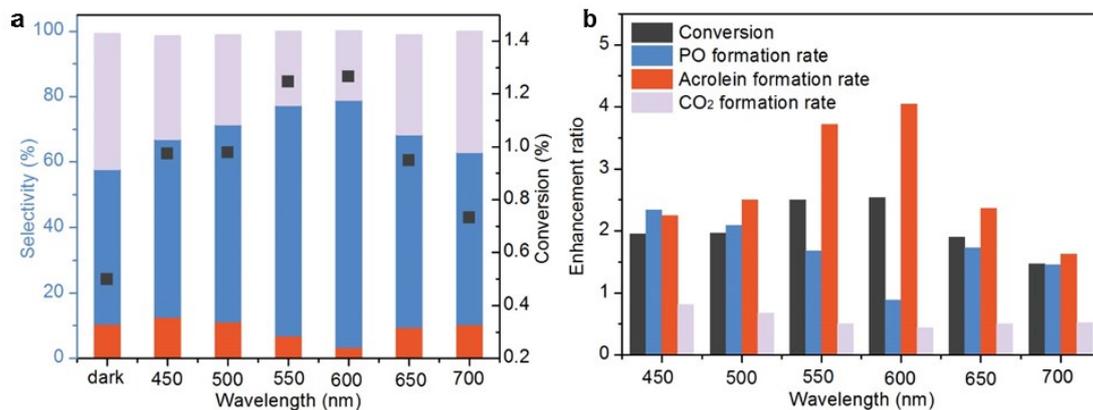


Fig. S30. The catalytic performance (conversion and selectivity) and enhancement ratios for the Au-Cu₂O hierarchical structure at 200 °C as a function of incident light intensity. (a) The catalytic performance (conversion and selectivity) for the Au-Cu₂O hierarchical structure as a function of incident light wavelength. (b) The enhancements of the propylene conversion and the product formation rate for Au-Cu₂O hierarchical structure at 200 °C as a function of incident light wavelengths.

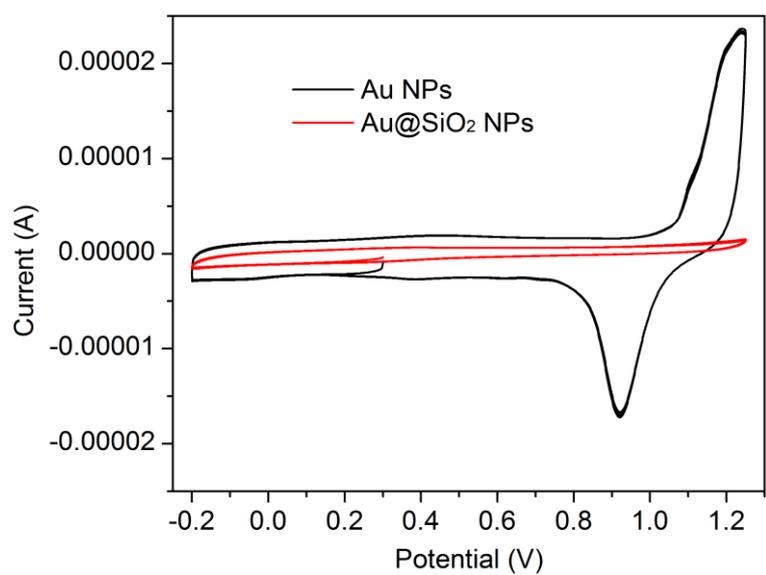


Fig. S31. CV curves of Au NPs or Au@SiO₂ NPs on glassy carbon electrode in 0.1 M H₂SO₄ solution.

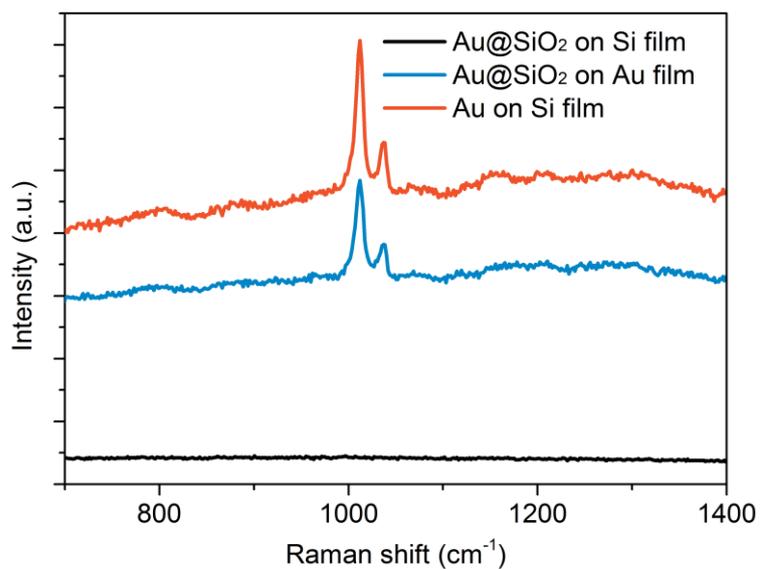


Fig. S32. Raman spectra of pyridine on gold for Au@SiO₂ on Si film, Au@SiO₂ on Au film and Au on Si film systems.

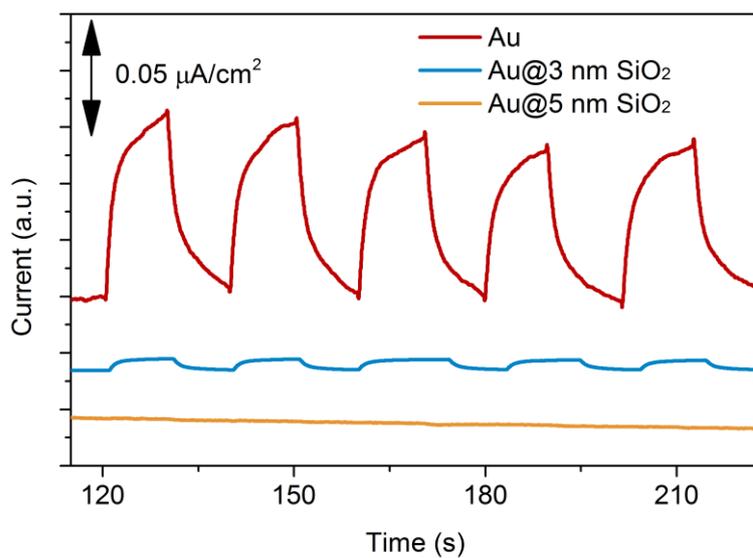


Fig. S33. The photocurrent of Au, Au@SiO₂ 3 nm and Au@SiO₂ 5nm on glassy carbon electrode in 0.1 M Na₂SO₄ solution at 0.6 V vs SCE.

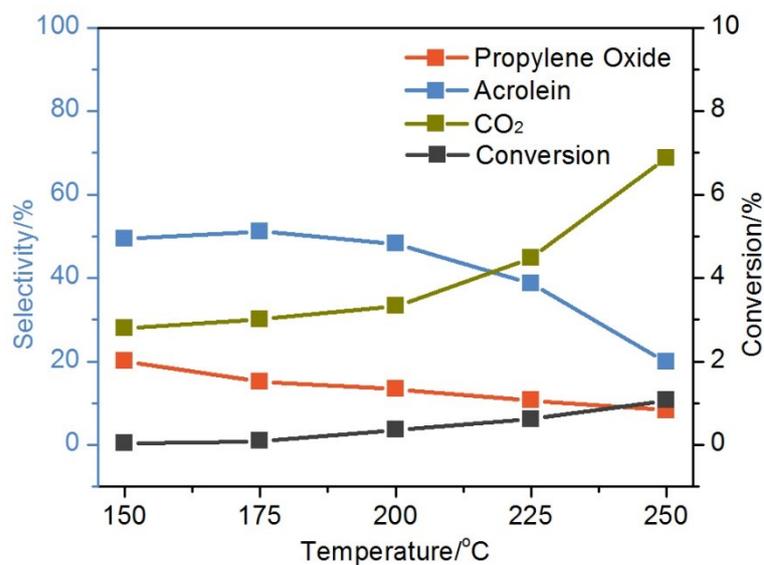


Fig. S34. Catalytic performance of Au@SiO₂-Cu₂O at different temperature. The main product is acrolein.

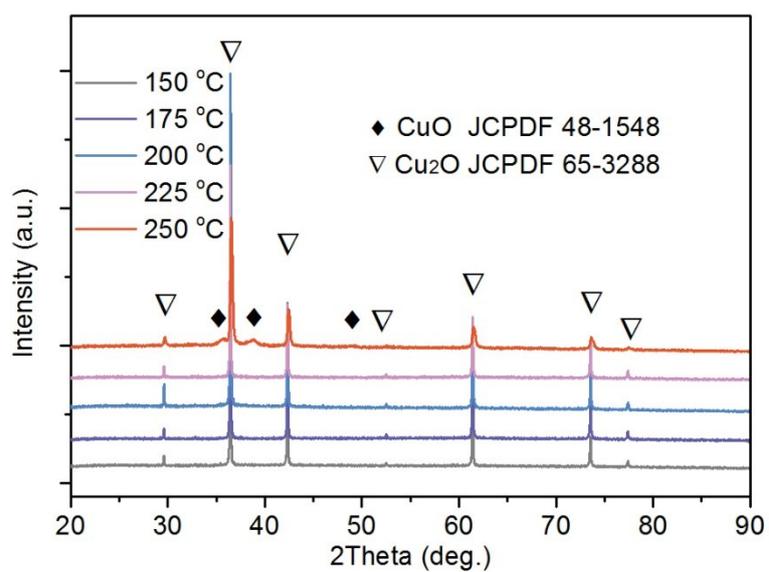


Fig. S35. XRD patterns and of Au@SiO₂-Cu₂O after catalysis at different temperature: (a) 150 °C, (b) 175 °C, (c) 200 °C, (d) 225 °C, (e) 250 °C.

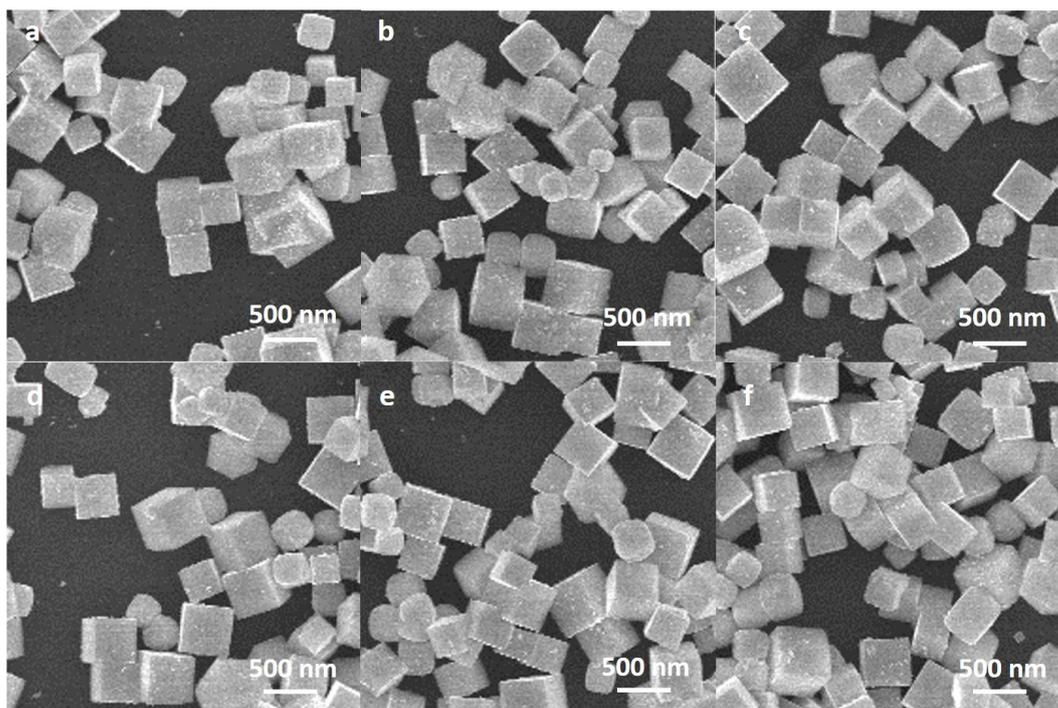


Fig. S36. SEM images of Au@SiO₂-Cu₂O after catalysis at different temperature: (a) 125 °C, (b) 150 °C, (c) 175 °C, (d) 200 °C, (e) 225 °C, (f) 250 °C.

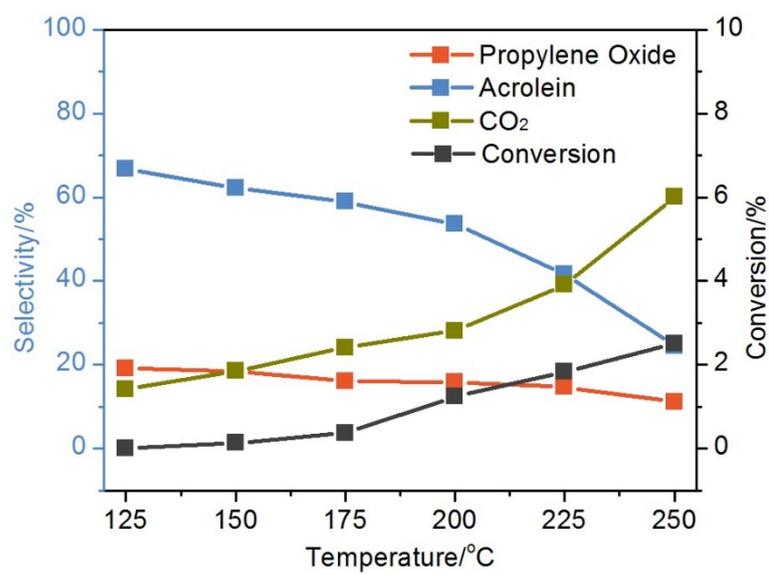


Fig. S37. Catalytic performance of Au@SiO₂-Cu₂O under illumination at different temperature. The main product is acrolein.

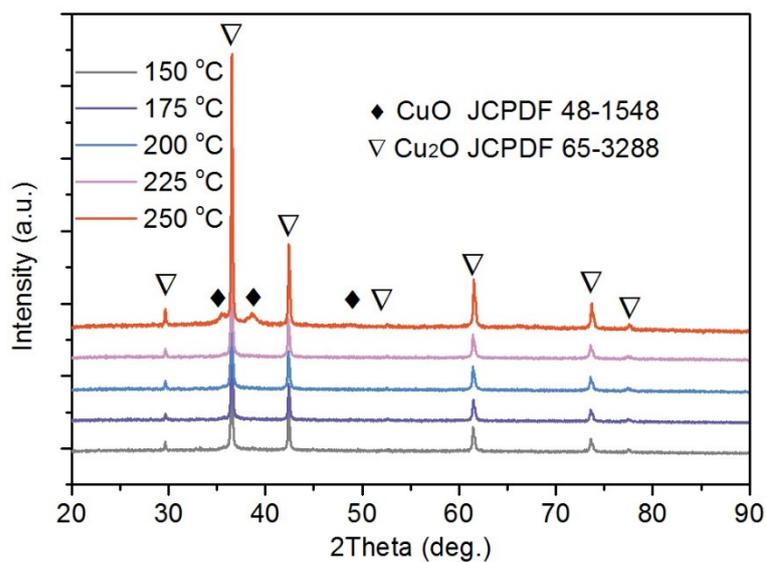


Fig. S38. XRD patterns and of Au@SiO₂-Cu₂O after catalysis under illumination at different temperature: (a) 150 °C, (b) 175 °C, (c) 200 °C, (d) 225 °C, (e) 250 °C.

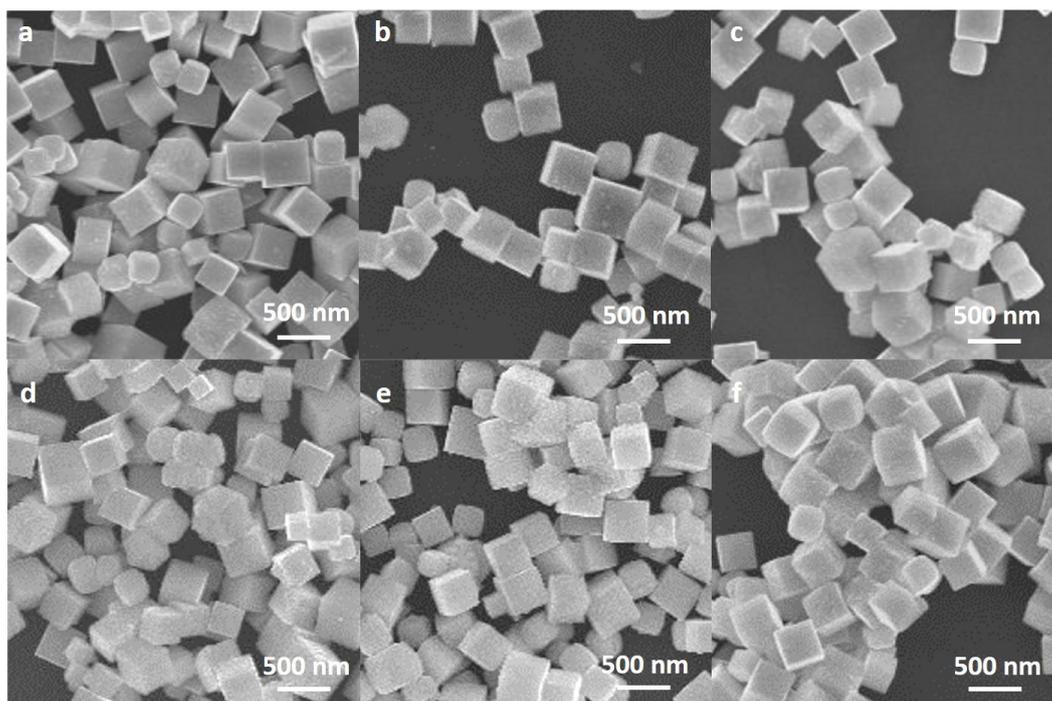


Fig. S39. SEM images of Au@SiO₂-Cu₂O after catalysis under illumination at different temperature: (a) 125 °C, (b) 150 °C, (c) 175 °C, (d) 200 °C, (e) 225 °C, (f) 250 °C.

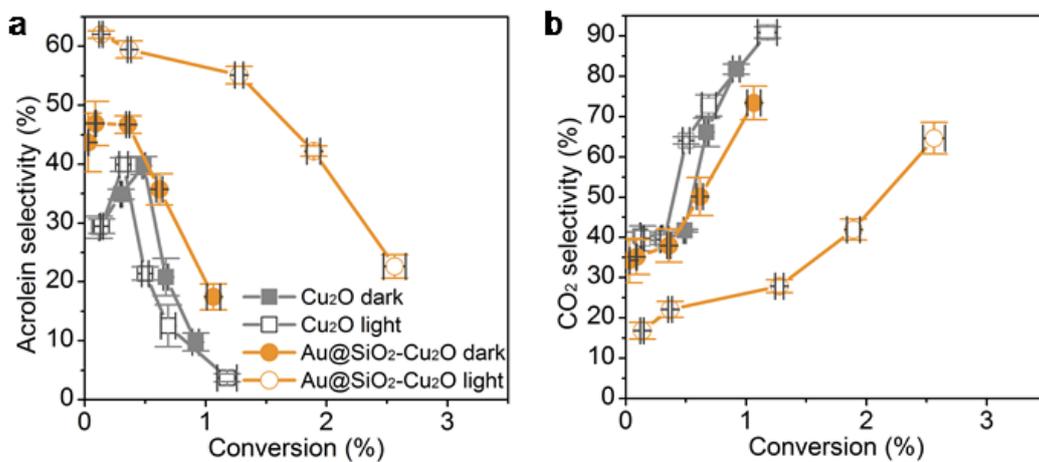


Fig. S40. Catalytic performance of Cu₂O and Au@SiO₂-Cu₂O. **a**, The selectivity of acrolein catalyzed by Cu₂O (blue) and Au@SiO₂-Cu₂O (red) with and without illumination as a function of propylene conversion. **b**, The selectivity of CO₂ for Cu₂O (blue) and Au@SiO₂-Cu₂O (red) with and without illumination as a function of propylene conversion.

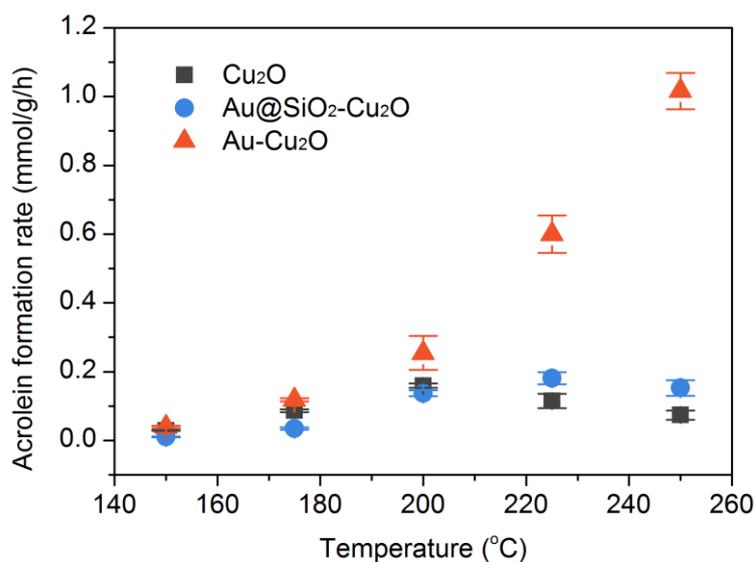


Fig. S41. Temperature dependent acrolein formation rate of Cu₂O, Au@SiO₂-Cu₂O and Au-Cu₂O.

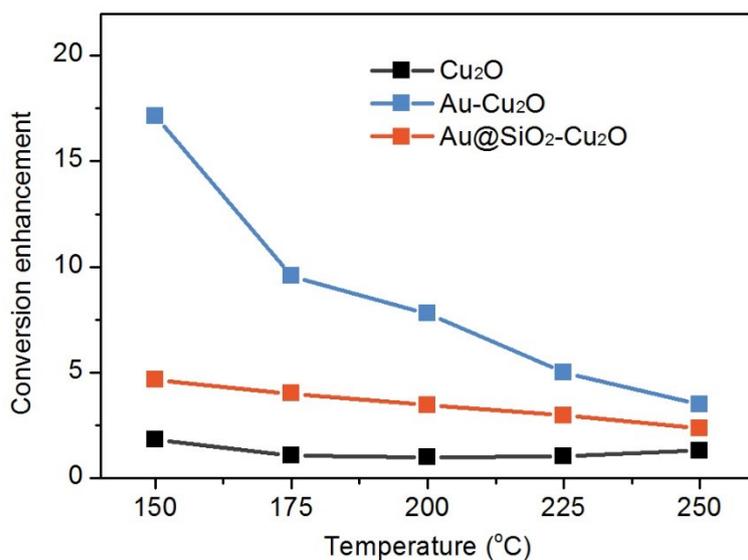


Fig. S42. The conversion enhancements induced by illumination for Cu₂O, Au-Cu₂O and Au@SiO₂-Cu₂O as a function of operating temperature.

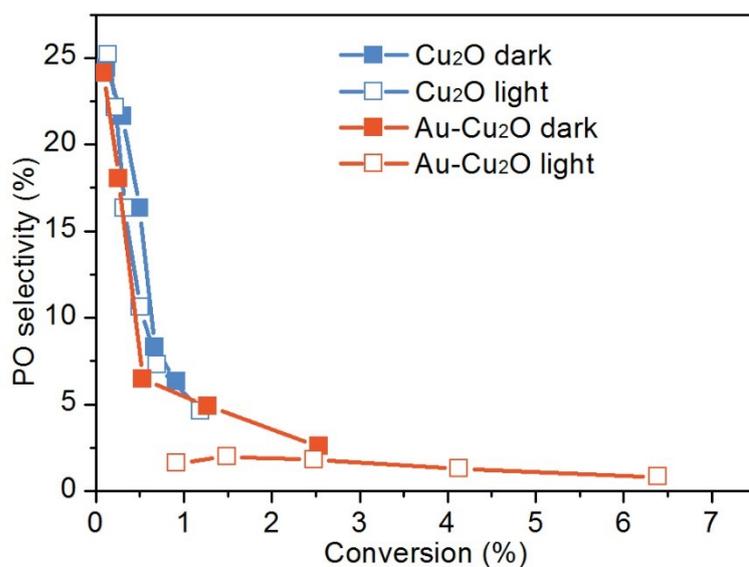


Fig. S43. The selectivity of propylene oxide catalyzed by Cu₂O and Au-Cu₂O with and without illumination as a function of propylene conversion.

Section S2. Theoretical Section

The photo thermal effect and temperature distribution is analyzed using the conventional macroscopic model, which is described by the usual heat transfer equation (52): $\rho c \frac{\partial T}{\partial t} = \nabla \cdot (k \nabla T) + Q_0$, where T is the local temperature increase, ρ is mass density, c is heat capacity and k is thermal conductivity. The local heat source density Q_0 originates from the non-radiative energy damping in gold nanoparticles, which writes as $Q_0 =$

$\frac{1}{2}\omega\varepsilon_0\varepsilon_{au}''|E_\omega|^2$, where ω is angular frequency of incident photon, ε_0 is the vacuum permittivity, ε_{au}'' is the imaginary part of the relative permittivity of gold adopted from Johnson and Christy's experimental data (53), and E_ω the local electric field derived from electromagnetic simulations. In the experimental condition, the illumination time is in tens of minutes, which is much longer than thermal relaxation times, therefore equation can be simplified as the Poisson's equation: $\nabla^2 T = -\frac{Q_0}{k}$. In homogenous medium, the stationary solutions of the equation involve the thermal Green's function and a Dirac source inside particles (46). The parameters used in the calculation is listed in Table S1.

Table S1. Parameters used in the calculations (59)

Materials	Refractive index, n	Thermal conductivity, k(W/m · k)
Gold	Ref. (53)	132.08
Cuprous (I) oxide	Ref. (59)	1.07
Helium	1	0.144

We begin with the analysis of single gold particle on cuprous (I) oxide substrate. As shown in figure S44, under normal incidence excitation, the electric field is strongly localized in the vicinity of particle-substrate gap. The confinement or field spatial resolution is about 20 nm, which is highly connected with the curvature of particle (here is the radius of gold sphere). Under asymmetric excitation condition, i.e. oblique incidence with a certain angle, the field will be further localized and enhanced. Figure S45 show a case of incident angles of 30° , yielding the confinement of field down to 6 nm.

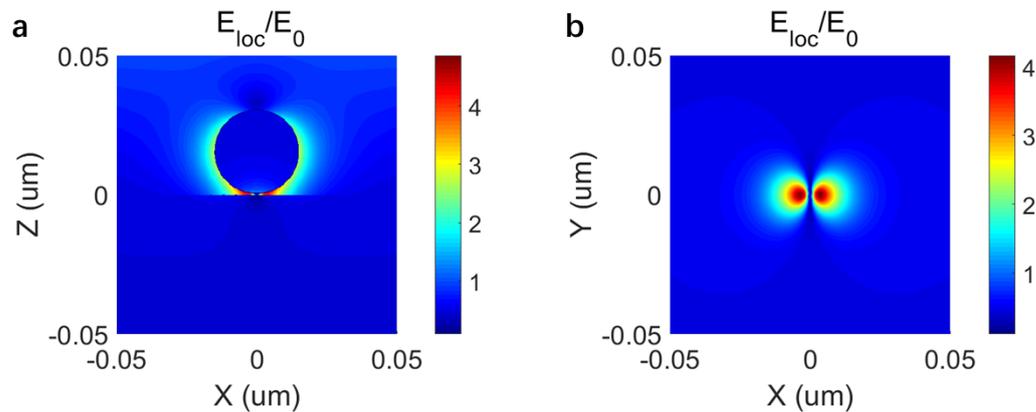


Fig. S44. The local electric field distribution, under normal incident from top to bottom. a) side view, the electric field is highly confined at the particle-substrate gap and forming electromagnetic hotspot, the air-substrate interface is at $z = 0$ plane. b) Electric

field in the surface of substrate, show a symmetrical dipole-like pattern due to the symmetrical excitation.

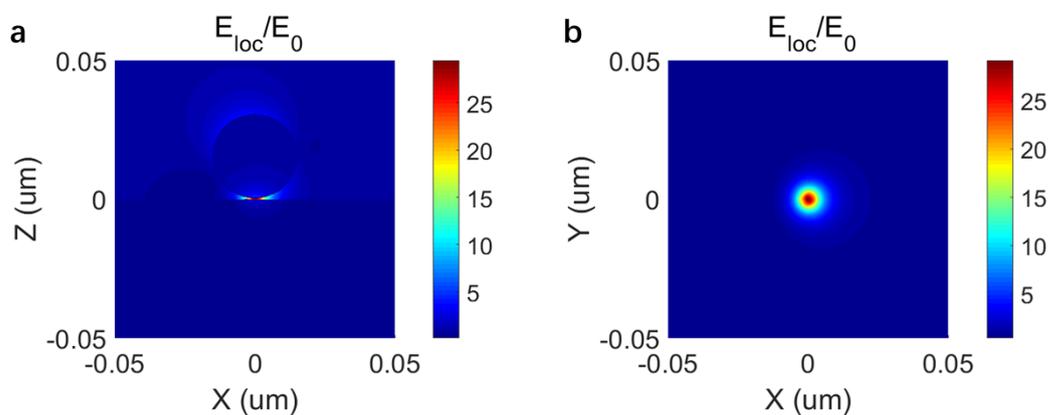


Fig. S45. The local field distribution under oblique incidence with an angle of 30° . a) side view, the field is localized in hotspots, with a maximal enhancement 25 times. b) top view, the electric field at the surface of substrate is further enhanced and confined due to the breaking of symmetry.

The temperature distribution can be derived via solving the heat transfer equation in steady-state. The results shown in figure S46 indicate a modest temperature increase of only 0.8 K, and the maximal temperature increase is located at the surface of nanoparticles. Since there are no other heat sources in the surrounding medium and the thermal conductivity of the particle is much higher than that of the medium, the temperature is almost uniform inside the particle and the thermal field also shows confinement in the vicinity of the nanoparticle surface and slowly decays in the medium. Figure S46c compares the distribution of temperature and electric field under normal and oblique incidence; the thermal field shows much weaker confinement than the electric case. The rate of most chemical reactions follows the Arrhenius law, hence, the reaction is much easier to trigger at the interface of the particle and substrate with both higher temperature and catalytic surface. We assume this active nano region is within 2 nm above the substrate and extract the temperature distribution at this cross-section, as shown in figure S46b and S46d. The results show a negligible temperature jump (~ 0.07 K) when approaching the particle surface. With such a modest temperature increase and local temperature jump, we will not expect any accelerated reactions.

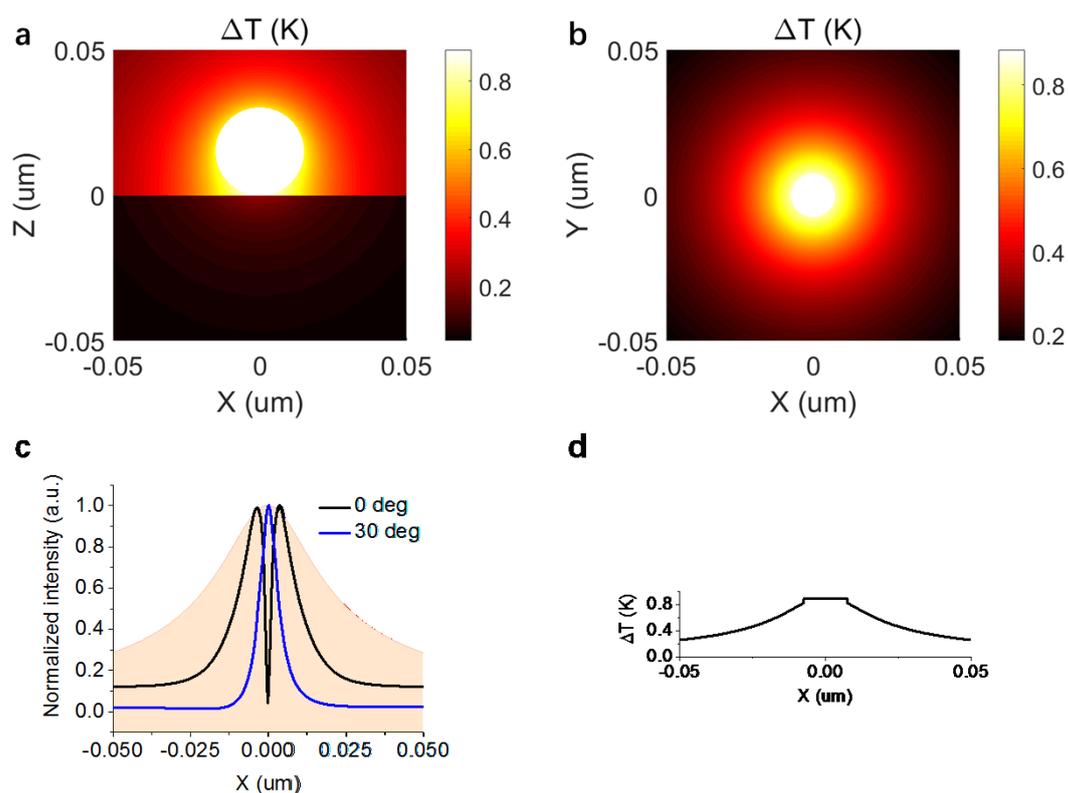


Fig. S46. Photo thermal effect induced by single gold particle. a) Side view of temperature distribution, the gas-substrate interface is at $z = 0$ plane. b) top view of temperature distribution, to show the temperature of nanoparticle surface, the cutting cross section is at $z = 2$ nm plane. c) Comparison of temperature (shown in light yellow), electric field under normal incidence (black solid line) and oblique incidence (blue solid line), at the surface of substrate. d) Temperature distribution at the surface of 2 nm above substrate, show a slight jump near particle region.

The mismatch between theory and experiments is also discussed in many photo thermal systems in nanoscale, i.e. solar steam generation (60, 61). One explanation is that the macroscopic model neglect of the interfacial thermal resistance between the particle and the surrounding medium (46, 61, 62), which is particularly important at the nanoscale.

For example, the surface potential can induce a vibrational energy shift of molecules adsorbed on gold surface, and the mismatch of the vibrational energies of an inner layer of molecules that attach particles and a layer away from particle surface can prevent the energy transfer hence block the heat flow (61, 63). In our case, a thermal insulating layer can be introduced due to adsorption of reactive C_3H_6 molecules, and resulting in the heat localization.

We compare the temperature distribution with and without thermal barrier layer. The results are shown in figure S47. For nanoparticle without thermal barrier, the thermal field is diffused and the temperature increase is limited, as shown in figure S47a. While for particle with thermal barrier (figure S47b, S47c), the heat is highly localized around

particle surface due to the insulating of energy exchange, and in the vicinity of particle surface, an abrupt temperature jump is observed. One can infer that the reactive molecules enriched on the surface of particles as well as substrate, and yielding a thermal insulating layer that block the heat flow and forming a high temperature region that highly localized around the surface of nanoparticles.

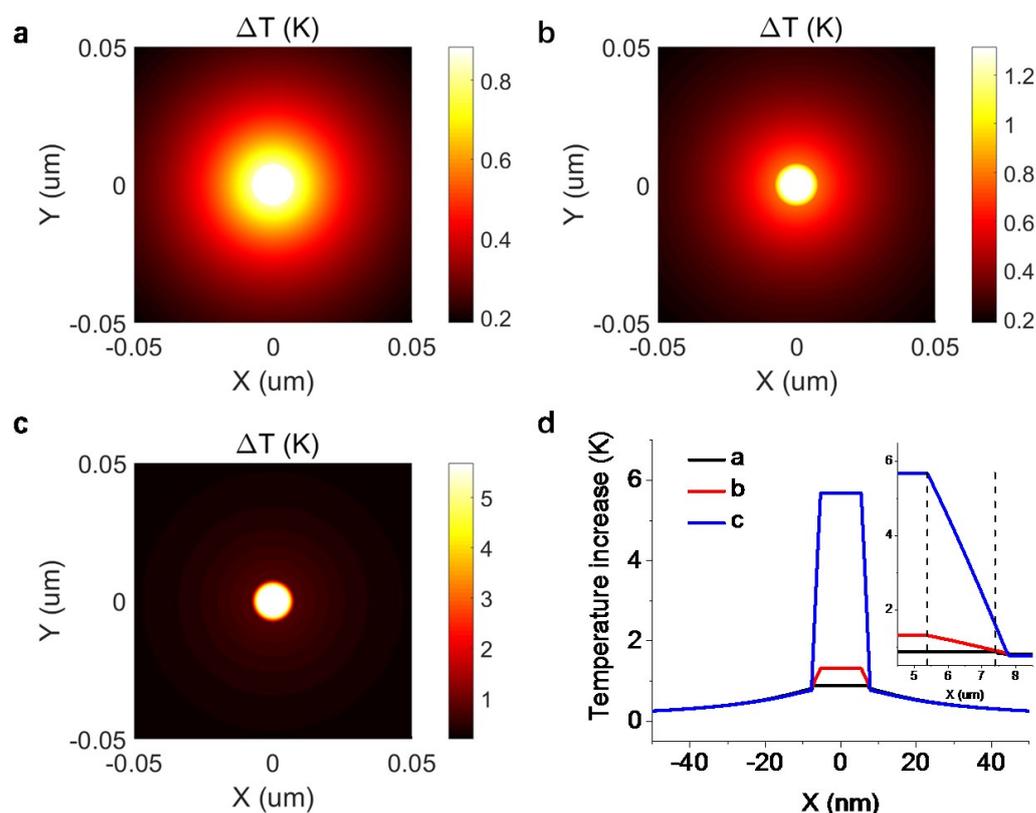


Fig. S47. The influence of surface thermal barrier. Temperature distribution at the surface of 2 nm above substrate, without surface thermal barrier layer (a), with 1 nm thermal barrier layer and the thermal conductivity equals to conductivity of C_3H_6 gas (b), with 1 nm thermal barrier layer and the thermal conductivity is one tenth of conductivity of C_3H_6 gas (c). (d) Temperature distribution along x direction of case in figure a (black solid line), b (red solid line) and c (blue solid line), respectively. The inset is a zoom in in figure d around the particle surface region, show an obvious temperature jump.

We have to emphasize that when a large amount of particles is illuminated, the temperature profile can be distinct from single particle situation, due to the thermal collective effect (64, 65). The temperature can become uniform throughout the particles, although the heat sources are localized. We analysis the collective heating effect via simulating 121 particles that assembled with varied periodicity, which reflect the surface density of particle numbers, as shown in figure S48. In a low particle density, as shown in figure S48a, the temperature is localized around particles and the temperature increase of surrounding medium is limited. While in a high particle density, as shown in figure S48c,

the temperature is no longer localized to the surface of particles, and we can observe a significant temperature increase of surrounding medium. In our experimental condition, the surface particle density can be approximately estimated from the SEM images, in which show a moderate number of $100 - 300/\mu\text{m}^2$. Therefore, we can safely claim that the localization of temperature still holds in our case, as shown in figure S48b.

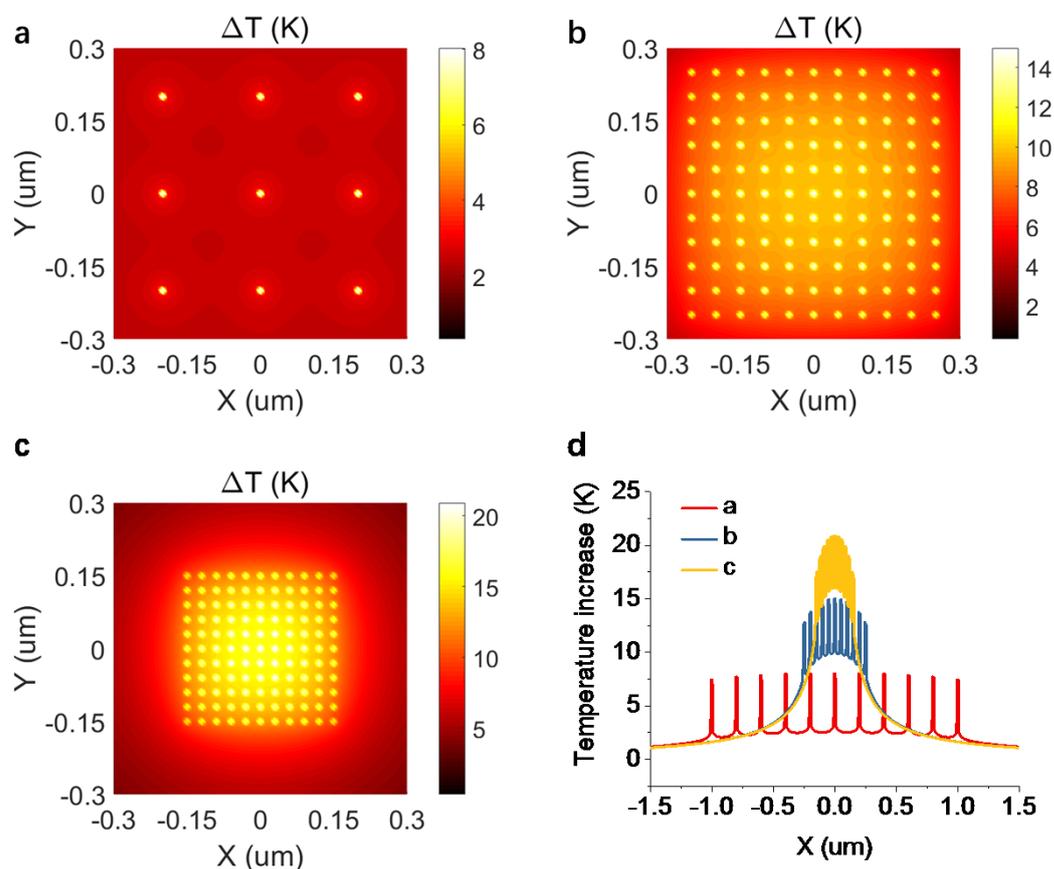


Fig. S48. Collective heating effect with different particle concentration. Here we used 11×11 particle array with varied periodicity to simulate the multi-particles on substrate surface. (a) The Temperature distribution with a low surface particle density of $25/\mu\text{m}^2$, the temperature field is localized in the vicinity of particle. (b) Temperature distribution with a moderate surface particle density of $300/\mu\text{m}^2$, the collective heating effect yielding an obvious temperature raise in surrounding medium. (c) Temperature distribution with a high surface particle density of $1300/\mu\text{m}^2$, the temperature is delocalized with a high temperature increase surrounding medium. (d) Temperature distribution along x direction, shown in figure a (red solid line) and b (blue solid line) and c (yellow solid line).

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