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# An ultrathin Al<sub>2</sub>O<sub>3</sub> bridging layer between CdS and ZnO boosts photocatalytic hydrogen production†

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The rapid separation and transfer of photogenerated charge carriers are the essential criteria for high photocatalytic efficiency. Heterostructures (HSs) composed of several semiconductors can take advantage of energy level difference to promote the separation and transfer of photogenerated electrons and holes, where the interface between semiconductors is the most critical position which directly determines the transfer mechanism of photogenerated charge carriers, and ultimately determines the photocatalytic efficiency. Herein, an ultrathin Al<sub>2</sub>O<sub>3</sub> bridging layer is intentionally introduced into the interface between CdS and ZnO by using an atomic layer deposition (ALD) method, and the resultant CdS@Al<sub>2</sub>O<sub>3</sub>@ZnO catalyst exhibits a significantly enhanced H<sub>2</sub> evolution rate (1190 μmol h<sup>-1</sup> m<sup>-2</sup>) and a high apparent quantum efficiency (AQE, 31.14% at a wavelength of 420 nm). It is revealed that the ultrathin Al<sub>2</sub>O<sub>3</sub> bridging layer converts the original Z-scheme mechanism of photogenerated charge carriers in the CdS@ZnO catalyst into a modified type II mechanism in the CdS@Al<sub>2</sub>O<sub>3</sub>@ZnO catalyst, where the photogenerated electrons are tunnelled from the CdS core to the ZnO shell and then are depleted by H<sub>2</sub>O effectively in time to produce H<sub>2</sub>. Meanwhile, the photogenerated holes are consumed efficiently by the fixed negative charges of Al<sub>2</sub>O<sub>3</sub>. Detailed characterization reveals that the lattice mismatch between wurtzite ZnO and wurtzite CdS is effectively passivated by the Al<sub>2</sub>O<sub>3</sub> interlayer, facilitating the uniform deposition of the ZnO layer, which promotes the close contact between interfaces, suppressing the recombination of charge carriers. This work provides a new strategy for controlling the dominant transfer mechanisms of photogenerated electrons and holes in HSs through interface regulation.

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## 1. Introduction

Heterostructures (HSs) consisting of several semiconductors have the combined merits of different semiconductors that complement each other, which is important to overcome the inherent shortcomings of single semiconductors, such as the limited light absorption and the restricted separation efficiency of photogenerated electrons and holes.<sup>1–4</sup> Therefore, HSs have been considered as a promising strategy to construct new photocatalysts with high activity for hydrogen production from the water splitting process.<sup>5–7</sup> To date, studies on HSs composed of binary (CdS/ZnS, TiO<sub>2</sub>@gC<sub>3</sub>N<sub>4</sub>),<sup>7,8</sup> ternary (graphene–CdS–MoS<sub>2</sub>, In<sub>2</sub>O<sub>3</sub>/InCd<sub>2</sub>S<sub>4</sub>/In<sub>2</sub>S<sub>3</sub>, ZnS/CdS/ZnO)<sup>9–11</sup> and quaternary components (ZnO/ZnS/CdS/CuInS<sub>2</sub>, TiO<sub>2</sub>/CdS/Co–Pi)<sup>12,13</sup> have

been widely carried out, and all HSs exhibit higher activity than single semiconductors because photogenerated electrons and holes in HSs can migrate to different semiconductors, efficiently suppressing the recombination.<sup>14–18</sup> It has been widely demonstrated that the transfer of photogenerated electrons and holes in HSs follows several mechanisms, such as type I (nested energy band distribution),<sup>19,20</sup> type II (staggered band distribution),<sup>21,22</sup> Z-scheme,<sup>23,24</sup> *etc.* For a HS, the transfer mechanisms of photogenerated electrons and holes in it are not constant, but vary with the environment of the HS, especially the surface and interface of the HS. In our previous work, we introduced Au nanoparticles (NPs) into the interface of C-doped TiO<sub>2</sub> and g-C<sub>3</sub>N<sub>4</sub>, and found that the transfer mechanism of charge carriers changed from type II into Z-scheme.<sup>23,25</sup> Recently, we successfully prepared an Au decorated hollow ZnO@ZnS HS, where two kinds of Au NPs located on the surface and at the interface of ZnO@ZnS were disclosed, which played different roles and resulted in different transfer mechanisms of charge carriers, respectively.<sup>26</sup> Therefore, we believe that the dominant transfer mechanisms of photogenerated electrons and holes in HSs can be effectively adjusted and controlled according to the ultimate goal.

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In an early study,<sup>27</sup> we successfully constructed a CdS@ZnO HS with a core@shell structure by using an atomic layer deposition (ALD) method. The H<sub>2</sub> evolution rate of the obtained CdS@ZnO was 4.1 times higher than the highest reported value (17.40 mmol g<sup>-1</sup> h<sup>-1</sup>)<sup>28</sup> at that time among the CdS–ZnO catalyst systems, and a direct Z-scheme mechanism was confirmed to dominate the transfer of photogenerated charge carriers. However, there is still much room for improving its photocatalytic H<sub>2</sub> evolution when the optimization of the HS is taken into consideration: (1) owing to the Z-scheme mechanism, the photogenerated electrons are accumulated on the conduction band of the CdS core, which cannot be consumed effectively in time by H<sub>2</sub>O to produce H<sub>2</sub> due to the barrier of the ZnO shell; (2) the ZnO shell is adhered to the CdS core as nanoparticles rather than a layer due to the lattice mismatch between wurtzite ZnO and wurtzite CdS, which is not conducive to the formation of a close contacted high quality interface between the two semiconductors, resulting in inefficient transfer of photogenerated charge carriers. If the photogenerated electrons can be transferred from the CdS core to the ZnO shell, and the lattice mismatch between wurtzite ZnO and wurtzite CdS can be overcome, the resultant CdS@ZnO HS is expected to exhibit significantly enhanced photocatalytic H<sub>2</sub> evolution.

Based on the above considerations, in this work, we selectively introduced an ultrathin Al<sub>2</sub>O<sub>3</sub> interlayer between CdS and ZnO by using the ALD method. It is found that the introduction of the Al<sub>2</sub>O<sub>3</sub> interlayer is efficient in addressing the two above-mentioned problems simultaneously. In the resultant CdS@Al<sub>2</sub>O<sub>3</sub>@ZnO HS, the photogenerated electrons were transferred from the CdS core to the ZnO shell due to a tunneling effect induced by the Al<sub>2</sub>O<sub>3</sub> interlayer, and meanwhile the photogenerated holes were consumed efficiently by the negative fixed charges of the Al<sub>2</sub>O<sub>3</sub> interlayer, resulting in a modified type II transfer mechanism. Furthermore, the lattice mismatch between wurtzite ZnO and wurtzite CdS was effectively passivated by the Al<sub>2</sub>O<sub>3</sub> interlayer, facilitating the uniform deposition of the ZnO layer, which promoted the close contact between interfaces, effectively suppressing the recombination of charge carriers. The photocatalytic performance of the CdS@Al<sub>2</sub>O<sub>3</sub>@ZnO for H<sub>2</sub> evolution was further improved, benefiting from the dual roles of the Al<sub>2</sub>O<sub>3</sub> interlayer.

## 2. Experimental section

### 2.1 Synthesis of CdS nanoparticles

CdS nanoparticles (NPs) were fabricated by using a hydrothermal method according to our previous report.<sup>27</sup> To be specific, 100 mL of sodium sulfide (Na<sub>2</sub>S, 0.14 M) solution was slowly added to 125 mL of cadmium acetate dihydrate (Cd(CH<sub>3</sub>COO)<sub>2</sub>·2H<sub>2</sub>O, 0.14 M) solution and stirred for 24 h. The obtained mixture was stewed for another 24 h, and then washed with deionized water three times. After that, the yellow pulp was dispersed in 80 mL of deionized water, sealed into a 100 mL stainless steel Teflon-lined autoclave and heated at 200 °C for 72 h. After cooling down to room temperature, the obtained product was washed with deionized water and ethanol three times and dried at 80 °C, resulting in CdS NPs.

### 2.2 Atomic layer deposition of Al<sub>2</sub>O<sub>3</sub> and ZnO

The CdS@Al<sub>2</sub>O<sub>3</sub>@ZnO multi-shell heterostructure was prepared by using a consecutive atomic layer deposition method. The substrate sample was obtained by dripping 40 mg of CdS NPs (dispersed in ethanol) on an 8 inch silicon wafer which was then dried up at 60 °C. The substrate sample was then placed in a vacuum ALD reaction chamber (ALD 004, SENTECH, Germany). Trimethylaluminum (TMA) and diethyl zinc (Zn(C<sub>2</sub>H<sub>5</sub>)<sub>2</sub>) were utilized as Al and Zn precursors for the deposition of Al<sub>2</sub>O<sub>3</sub> and ZnO, respectively, and water was used as another precursor to provide the O source. Nitrogen served as a carrier gas. The temperature of the reaction chamber was set at 250 °C to ensure the evaporation of precursors during the reaction, and the deposition amount of the shell layer was controlled by the number of deposition cycles. One standard deposition cycle consisted of four steps: injection of the Zn or Al precursor, nitrogen flow, injection of the oxygen precursor, and nitrogen flow. The number of deposition cycles for Al<sub>2</sub>O<sub>3</sub> was 1, 2, 6 and 10, while the number of deposition cycles for ZnO was 100. The obtained CdS@Al<sub>2</sub>O<sub>3</sub>@ZnO core-shell structures were denoted as Cd@Al<sub>x</sub>@Zn, where *x* is the number of deposition cycles of Al<sub>2</sub>O<sub>3</sub>.

In addition, a control sample, CdS@ZnO@Al<sub>2</sub>O<sub>3</sub> (Cd@Zn@Al), was fabricated by switching the deposition order of Al<sub>2</sub>O<sub>3</sub> and ZnO. The number of deposition cycles of ZnO and Al<sub>2</sub>O<sub>3</sub> was 100 and 2, respectively, and the obtained sample was denoted as Cd@Zn@Al<sub>2</sub>.

In order to get pure Al<sub>2</sub>O<sub>3</sub> and minimize the influence of the synthesis method on the properties of the product, we employed the same ALD method in the fabrication of Al<sub>2</sub>O<sub>3</sub> by using an 8 inch silicon plate as the substrate. 200 deposition cycles were applied to get enough amount of Al<sub>2</sub>O<sub>3</sub> for the test. After deposition, the white sample was collected by scraping carefully with a clean blade.

### 2.3 Photocatalytic hydrogen generation

The photocatalytic water splitting reaction was carried out in a top-irradiation quartz cell connected to a closed glass gas circulation system (CEL-SPH2N, Beijing), and a 300 W xenon arc lamp was used as the simulated solar light source. 10 mg of the as-synthesized photocatalyst was dispersed in 25 mL of deionized water containing 0.35 M sodium sulfide and 0.25 M sodium sulfite anhydrous, and 1 wt% Pt was added as the co-catalyst. The amount of H<sub>2</sub> evolved under irradiation was monitored via a TCD gas chromatograph (GC-9720). Magnetic stirring was used during the water splitting experiment to ensure homogeneity of the suspension.

Besides the mass of catalysts, the light absorption ability and specific surface area are two of the most crucial factors in the photocatalytic reaction. Here, the multi-factor-normalized H<sub>2</sub> evolution rate was adopted to evaluate the photocatalytic performance of the catalysts.<sup>29,30</sup> CdS was used as the reference sample in the normalization treatment; for instance, the light absorption ability of CdS is defined as 1, and the light absorption ability of one catalyst *versus* CdS, named relative light absorption coefficient (*r*<sub>Abs.</sub>), is determined from the ratio of

total absorption intensity of the studied catalyst and that of CdS. The multi-factor-normalized H<sub>2</sub> evolution rate was determined from the following equation:

$$R(\text{H}_2) = \frac{N(\text{H}_2)}{t \times S_{\text{BET}} \times r_{\text{Abs.}}}$$

where,  $R(\text{H}_2)$  is the relative H<sub>2</sub> evolution rate and  $N(\text{H}_2)$  is the amount of hydrogen generated from the water splitting reaction during a certain period of time.  $t$ ,  $S_{\text{BET}}$  and  $r_{\text{Abs.}}$  are the reaction time, specific surface area and relative light absorption coefficient of the catalysts, respectively.

In the long-term stability test, every cycle was maintained for 5 hours. After a cycle, the pressure of the gas circulation system was balanced by connecting directly to environmental conditions, while the TCD gas chromatograph detection system was pumped with N<sub>2</sub> to remove the residual H<sub>2</sub>. At the same time, the sacrificial agent with Na<sub>2</sub>SO<sub>3</sub> and Na<sub>2</sub>S was refreshed for the next cycle. All the reaction conditions in the long-term stability test were the same as for the above-mentioned photocatalytic H<sub>2</sub> evolution.

## 2.4 Characterization

X-ray diffraction (XRD) patterns were recorded on a Bruker AXS D2 Phaser instrument using CuK<sub>α</sub> radiation in step mode between 20° and 60°. Transmission electron microscopy (TEM) and energy dispersive spectroscopy (EDS) mapping images were obtained on a JEOL JEM-2100F microscope operating at a voltage of 200 kV. Scanning electron microscopy (SEM) images were obtained on a field emission scanning electron microscope (Gemini SEM 500). X-ray photoelectron spectroscopy (XPS) analyses were performed on a Thermo Fisher ESCALAB Xi+ analyzer with aluminum K<sub>α</sub> radiation. UV-vis absorption spectra were recorded with a UV-Vis-NIR spectrometer (JASCO, V-670). Photoluminescence (PL) spectra and fluorescence lifetimes were measured using a fluorescence spectrometer (Edinburgh Instruments FLS 980), and the excitation wavelength was 370 nm. Transient photocurrent response measurement and electrochemical impedance spectroscopy (EIS) were carried out in a standard three-compartment cell, where a saturated calomel electrode (SCE) was used as the reference electrode. To avoid the effect of dissolved Pt<sup>2+</sup> on the measurement,<sup>31,32</sup> carbon paper (1 cm × 1 cm) was utilized as the counter electrode, and the working electrode was made by loading the catalyst on fluorine-doped tin oxide (FTO, 7 Ω per square) glasses. 0.1 mol L<sup>-1</sup> Na<sub>2</sub>SO<sub>4</sub> solution was used as the electrolyte, and a 300 W Xe lamp was utilized as the light source. Light on and off operation was controlled by a shutter every 20 seconds and the response of transient photocurrent was recorded on a CHI 660 E electrochemical workstation (Shanghai, China).

## 2.5 Detection of hydroxyl radicals

Hydroxyl radicals ( $\cdot\text{OH}$ ) were detected by the PL method using terephthalic acid (TA) as a probe molecule. In a typical procedure, 10 mg of the sample was dispersed in a 20 mL aqueous solution of  $5 \times 10^{-4}$  M TA and  $2 \times 10^{-3}$  M NaOH. The

suspension was irradiated under a 300 W Xe lamp for 10 min with stirring, after which, a solution was obtained through centrifugation to remove catalyst particles. Subsequently, the solution was analyzed under the excitation of 315 nm wavelength by using a FLS980 fluorescence spectrometer.

## 2.6 Computational details

All calculations were performed using periodic density functional theory (DFT) through the Vienna Ab Initio Simulation Package (VASP).<sup>33–35</sup> The interaction between the core and valence electrons was described using the projector-augmented wave (PAW) method.<sup>36,37</sup> The PBE sol functional was applied for geometry optimization.<sup>38</sup> However, the hybrid Heyd–Scuseria–Ernzerhof functional (HSE06) including 25% screened Hartree–Fock (HF) exchange,<sup>9,40</sup> hybrid PBE0 functional using 25% HF exchange,<sup>41,42</sup> and HSE with 37.5% HF exchange (HSE375) were employed for the electronic property calculations of wurtzite CdS,  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> and wurtzite ZnO, respectively, because these parameters were found to reproduce their band gaps.<sup>43–45</sup> A plane wave cutoff energy of 560 eV was applied for all electronic calculations, with  $\Gamma$ -centred  $k$ -point meshes of  $6 \times 6 \times 4$ ,  $8 \times 8 \times 4$  and  $8 \times 8 \times 4$  for unit cells of CdS, Al<sub>2</sub>O<sub>3</sub> and ZnO, respectively, and was determined to be sufficient to converge the total energy of each system to within 10<sup>-5</sup> eV per atom.

Three two-phase heterojunctions were built using  $1 \times 2$  CdS (110)/ $1 \times 1$  Al<sub>2</sub>O<sub>3</sub> ( $-111$ ),  $1 \times 1$  Al<sub>2</sub>O<sub>3</sub> ( $-111$ )/ $2 \times 2\sqrt{3}$  ZnO (001) and  $1 \times 1$  CdS (110)/ $2 \times \sqrt{3}$  ZnO (001). These combinations were necessary to balance the total number of atoms and the strain caused by the mismatch of two phases. The electrostatic potentials of the heterojunctions were calculated using the hybrid HSE06 including 25% screened Hartree–Fock exchange.<sup>39,40</sup> A 20 Å vacuum was employed for all surface slabs to vanish wave function overlaps across the vacuum region. The ionization potential was calculated using the core-level alignment approach of Wei and Zunger,<sup>46</sup> and the non-polar CdS (110), Al<sub>2</sub>O<sub>3</sub> ( $-111$ ) and ZnO (100) surfaces were chosen.

# 3. Results and discussion

## 3.1 The microstructure and photocatalytic performance of Cd@Al<sub>2</sub>@Zn

Fig. 1a shows the XRD patterns of the four selected samples, CdS, Cd@Zn, Cd@Al<sub>2</sub> and Cd@Al<sub>2</sub>@Zn. It can be observed that all samples display the dominant characteristic peaks of wurtzite CdS (JCPDS 41-1049).<sup>47–49</sup> After 100 deposition cycles of ZnO, several weak peaks at 30–35° can be found from the XRD of Cd@Zn, and their low intensity can be attributed to the very small amount as well as the high dispersivity of ZnO.<sup>48–51</sup> No XRD signal of Al<sub>2</sub>O<sub>3</sub> was detected, which is mainly due to the amorphous nature and very small amount of Al<sub>2</sub>O<sub>3</sub>.<sup>16,43</sup> The XRD spectra of the Si substrate before and after Al<sub>2</sub>O<sub>3</sub> deposition are presented in Fig. S1.† The wide diffraction peaks between 15 and 35 degrees come from the SiO<sub>2</sub> on the surface of the Si substrate, and no new peak appeared even after 200 cycles of Al<sub>2</sub>O<sub>3</sub> deposition, indicating the amorphous nature of Al<sub>2</sub>O<sub>3</sub>.<sup>43,52,53</sup> It should be noticed that with the same number of

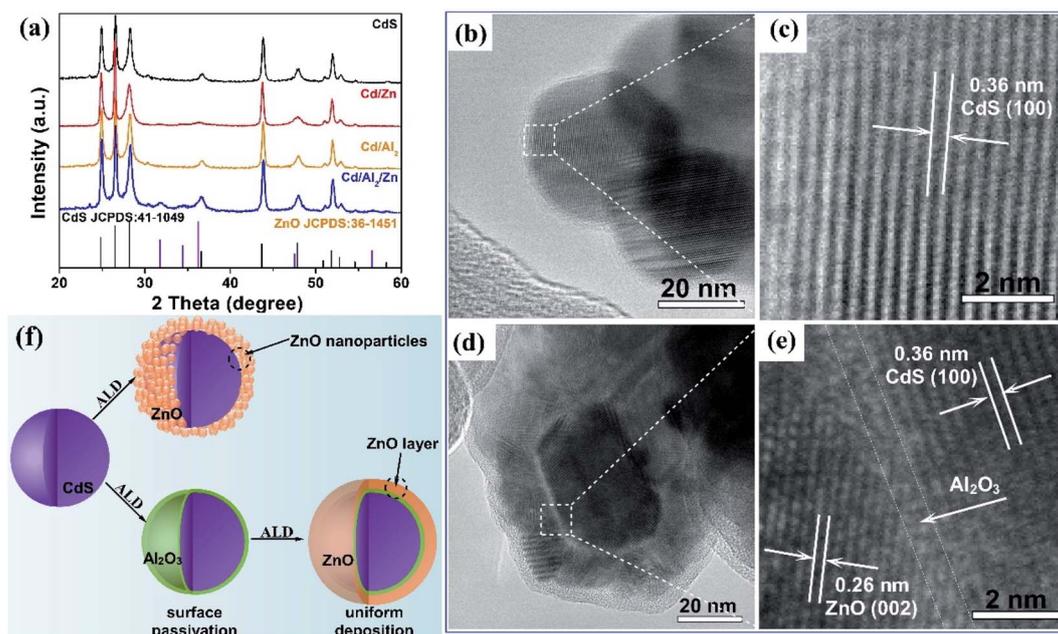


Fig. 1 (a) XRD spectra of CdS, Cd@Zn, Cd@Al<sub>2</sub> and Cd@Al<sub>2</sub>@Zn (the peaks of CdS and ZnO are marked by black and purple graphic patterns, respectively). (b and d) TEM and (c and e) HRTEM images of (b and c) CdS and (d and e) Cd@Al<sub>10</sub>@Zn. (f) Schematic view of the interface difference between Cd@Zn and Cd@Al@Zn fabricated by the ALD method.

deposition cycles of the ZnO layer, the ZnO peaks in the XRD pattern of Cd@Al<sub>2</sub>@Zn were obviously stronger than those of Cd@Zn, and thus we suppose that the introduction of Al<sub>2</sub>O<sub>3</sub> facilitates the deposition of the ZnO layer.

The above assumption can be confirmed by TEM characterization. The pristine CdS presented a uniform spherical morphology with an average diameter of  $35 \pm 5$  nm and a smooth surface (Fig. 1b). The characteristic (100) lattice fringes (0.36 nm) of wurtzite CdS can be clearly observed from Fig. 1c, indicating the good crystallinity of CdS NPs.<sup>48</sup> The TEM images of Cd@Zn show the deposited ZnO as nanoparticles attached on the surface of CdS (Fig. S2a†). The distinguished lattice spacing of about 0.26 nm (Fig. S2b†) matched well with the (002) interplanar spacing of wurtzite ZnO (JCPDS 36-1451).<sup>54–56</sup> For easy observation, the TEM of Cd@Al<sub>10</sub>@Zn was carried out. As shown in Fig. 1d and e, the surface of CdS was covered by the Al<sub>2</sub>O<sub>3</sub> layer and ZnO layer, and a light-colored Al<sub>2</sub>O<sub>3</sub> layer appeared between the CdS core and the ZnO shell (Fig. 1d). The relative position of CdS, Al<sub>2</sub>O<sub>3</sub> and ZnO can be further identified by EDX mapping. As shown in Fig. S3,† Cd and S signals are located at the center of the particles while the outer position is dominated by O, Zn and Al elements. In addition, the overlap images of S&Cd&Zn and Cd&Al&Zn are in accord with the HADDF image of the sample, demonstrating the formation of the core-shell heterostructure. Two kinds of lattice parameters, corresponding to CdS (0.36 nm) and ZnO (0.26 nm), can be observed from the HRTEM image (Fig. 1e), and the light-colored region (labeled with a white dotted line in Fig. 1e) shows the amorphous characteristics of the Al<sub>2</sub>O<sub>3</sub> layer with an average thickness of about 1 nm. It should be noted that the ZnO deposited directly on the CdS core is in the form of

nanoparticles rather than a layer, which has been confirmed in our previous work.<sup>48</sup> The reason can be ascribed to the mismatched lattice between wurtzite ZnO and wurtzite CdS. After the wurtzite CdS core is wrapped by an amorphous Al<sub>2</sub>O<sub>3</sub> layer, the wurtzite ZnO deposited on Cd@Al is a layer rather than nanoparticles, suggesting that the mismatched lattice between wurtzite ZnO and wurtzite CdS is passivated by the amorphous Al<sub>2</sub>O<sub>3</sub> layer. As illustrated in Fig. 1f, when Al<sub>2</sub>O<sub>3</sub> is deposited on the surface of CdS, a uniform deposition layer and an intimate interface can be obtained due to the good surface compatibility of Al<sub>2</sub>O<sub>3</sub>. After that, diethyl zinc and water are alternatively pumped into the reaction chamber for the deposition of ZnO. The uniformly deposited Al<sub>2</sub>O<sub>3</sub> provides an ideal substrate for the nucleation and growth of ZnO, thus playing the role of a seed layer. In addition, in the formation of the heterostructure, the good surface compatibility of Al<sub>2</sub>O<sub>3</sub> helps shrink the mismatch and lattice distortion between CdS and ZnO and facilitates the formation of an intimately contacted interface, thus playing the role of a buffer layer. The SEM images (Fig. S4†) further verified the role of the Al<sub>2</sub>O<sub>3</sub> interlayer in the deposition of the ZnO shell. A fluffy surface of the sample can be clearly observed from the SEM images of Cd@Zn (Fig. S4a and b†), which mainly comes from the uneven deposition of ZnO NPs. However, a relatively smooth surface was observed over the Cd@Al<sub>10</sub> and Cd@Al<sub>10</sub>@Zn samples (Fig. S4c–f†). Since the same deposition procedure of ZnO was utilized in all samples, the different morphologies of ZnO in Cd@Zn and Cd@Al<sub>10</sub>@Zn can be ascribed to the Al<sub>2</sub>O<sub>3</sub> interlayer.

UV-vis absorption spectroscopy and BET measurements were carried out to investigate the effect of the Al<sub>2</sub>O<sub>3</sub> interlayer on the light absorption ability and specific surface area of the catalysts.

As shown in Fig. S5a,<sup>†</sup> no obvious change of light absorption ability was observed for the samples, this is because of the translucent properties of ZnO and Al<sub>2</sub>O<sub>3</sub> as well as the very low loading amount of Al<sub>2</sub>O<sub>3</sub>. The deposition of ZnO and Al<sub>2</sub>O<sub>3</sub> resulted in obvious change in the specific surface area of the samples, and as can be observed from Fig. S5c,<sup>†</sup> Cd@Zn displayed a much higher surface area than pure CdS due to the uneven deposition of ZnO, which can be confirmed by the results of SEM and TEM observation. After the introduction of the Al<sub>2</sub>O<sub>3</sub> interlayer, the specific surface area of Cd@Al<sub>x</sub>@Zn gradually decreased to a stable level, the same as that of CdS. Based on previous experiments and discussion, the decrease of surface area can be attributed to the introduction of a highly-surface-compatible-Al<sub>2</sub>O<sub>3</sub> layer, which facilitated the uniform and dense deposition of the ZnO layer.

The photocatalytic performances of the samples were investigated by monitoring the hydrogen evolution rate. Chloroplatinic acid was added as the precursor to provide a Pt cocatalyst, and we detected the components of the solution by using ICP-MS and found there were no Pt ions in the solution after reaction, indicating that Pt ions have been reduced into Pt nanoparticles and deposited on the surface of the photocatalyst. Multi-factor-normalized H<sub>2</sub> evolution rate was adopted to give a more accurate determination of the photocatalytic activity of the samples, and the effect of light absorption ability was eliminated by the introduction of the relative light absorption coefficient (*r*<sub>Abs.</sub>). As shown in Fig. 2a, CdS displayed a relatively low H<sub>2</sub> evolution rate of 45 μmol h<sup>-1</sup> m<sup>-2</sup> due to the serious charge recombination.<sup>55,56</sup> After depositing Al<sub>2</sub>O<sub>3</sub>, a slightly increased hydrogen generation ability was observed over the

Cd@Al<sub>x</sub> samples, which is much lower than that over Cd@Zn samples (548 μmol h<sup>-1</sup> m<sup>-2</sup>) due to the as-reported Z-scheme transfer mechanism between ZnO and CdS.<sup>48,57,58</sup> After a ZnO layer was deposited on the Cd@Al<sub>2</sub> sample, the photocatalytic activity of Cd@Al<sub>2</sub>@Zn was significantly enhanced. As observed in Fig. 2b, the H<sub>2</sub> evolution rate climbed from 807 μmol h<sup>-1</sup> m<sup>-2</sup> for Cd@Al<sub>1</sub>@Zn to 1190 μmol h<sup>-1</sup> m<sup>-2</sup> for Cd@Al<sub>2</sub>@Zn, and then reduced to 1084 μmol h<sup>-1</sup> m<sup>-2</sup> for Cd@Al<sub>6</sub>@Zn, suggesting that the Al<sub>2</sub>O<sub>3</sub> interlayer with a suitable thickness contributes to the improvement of photocatalytic H<sub>2</sub> evolution performance of Cd@Al@Zn. As an important control sample, we fabricated Cd@Zn@Al<sub>2</sub> by switching the deposition order of ZnO and Al<sub>2</sub>O<sub>3</sub>, and the H<sub>2</sub> evolution ability of Cd@Zn@Al<sub>2</sub> (692 μmol h<sup>-1</sup> m<sup>-2</sup>) was much lower than that of Cd@Al<sub>2</sub>@Zn (1190 μmol h<sup>-1</sup> m<sup>-2</sup>) (the H<sub>2</sub> evolution rates before normalization are shown in Fig. S6<sup>†</sup>). We carefully studied the hydrogen evolution rates of different CdS-ZnO based H<sub>2</sub> evolution improvement strategies reported in the literature (Table S2<sup>†</sup>), and found that the hydrogen evolution rate of 101.3 mmol h<sup>-1</sup> g<sup>-1</sup> (1190 μmol h<sup>-1</sup> m<sup>-2</sup>) over Cd@Al<sub>2</sub>@Zn is the highest level. The result demonstrates that the introduction of the Al<sub>2</sub>O<sub>3</sub> passivation interlayer is promising for improving the photocatalytic H<sub>2</sub> evolution ability. The apparent quantum efficiency (AQE) of Cd@Al<sub>2</sub>@Zn at different wavelengths was further tested. As shown in Fig. 2c, the highest AQE value of about 31.14% was obtained at a wavelength of 420 nm. A long-term cycle test exhibited that no obvious difference in photocatalytic ability (Fig. 2d) and crystal structure (Fig. S7<sup>†</sup>) appeared during the 20 hour reaction (4 cycles), indicating the good photocatalytic stability of the as-fabricated Cd@Al<sub>2</sub>@Zn.

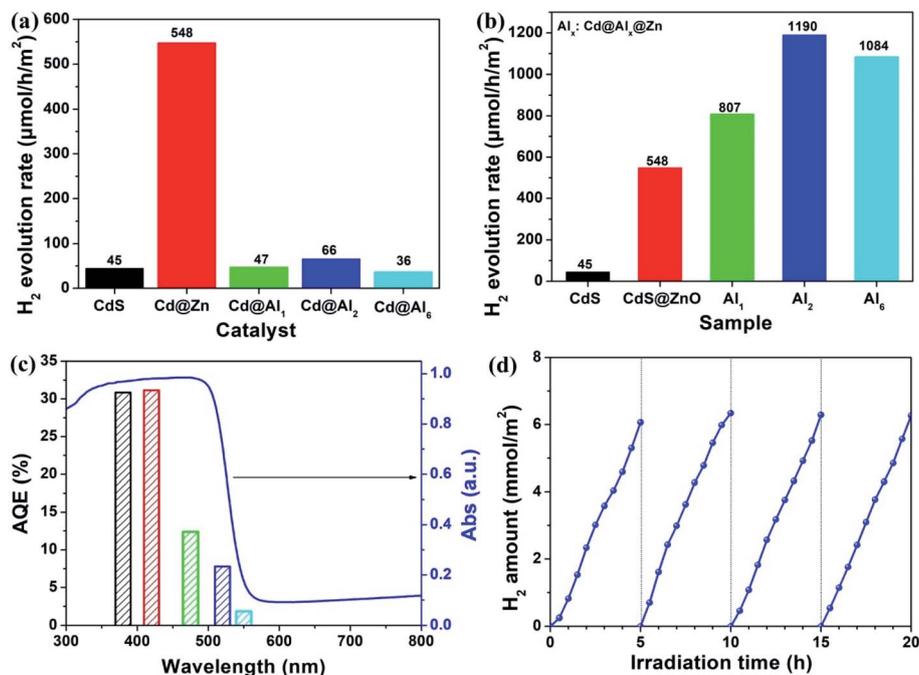


Fig. 2 (a and b) The comparison of photocatalytic H<sub>2</sub> generation rates over the as-fabricated samples based on normalized light absorption ability and specific surface area, the amount of catalysts is 10 mg. (c) The measured AQE of Cd@Al<sub>2</sub>@Zn at different wavelengths. (d) Long-term cycling reaction test of Cd@Al<sub>2</sub>@Zn for 20 hours.

Considering that cocatalysts play an important role in the photocatalytic reaction, we further investigated the effect of the Pt cocatalyst by monitoring the H<sub>2</sub> evolution over the Pt-free samples. The H<sub>2</sub> evolution ability of CdS, Cd@Zn and Zd@Al<sub>2</sub>@Zn in the absence of the Pt cocatalyst is shown in Fig. S8,† and it can be observed that the H<sub>2</sub> evolution rate over Cd@Al<sub>2</sub>@Zn is much higher than that of Cd@Zn, which is in accord with the performance in the presence of the cocatalyst, indicating that the introduction of the Al<sub>2</sub>O<sub>3</sub> interlayer contributes to the H<sub>2</sub> evolution. In the absence of the Pt cocatalyst, the highest H<sub>2</sub> evolution rate over Cd@Al<sub>2</sub>@Zn was 201.3 μmol h<sup>-1</sup>, which is much lower than that in the presence of the Pt cocatalyst (1013.4 μmol h<sup>-1</sup>), demonstrating that the existence of the Pt cocatalyst is important in facilitating the generation of H<sub>2</sub>.

### 3.2 The role of the ultrathin Al<sub>2</sub>O<sub>3</sub> interlayer

The ultrathin Al<sub>2</sub>O<sub>3</sub> interlayer between CdS and ZnO significantly boosted the photocatalytic hydrogen production. In order to get better understanding of the structure–property relationship in Cd@Al@Zn, we carried out a series of experimental and theoretical studies to disclose the role of the ultrathin Al<sub>2</sub>O<sub>3</sub> interlayer. Fig. S9a† displays the XPS spectra of Cd@Zn and Cd@Al<sub>2</sub>@Zn, which confirmed the presence of Cd, S, O, and Zn elements, while the signal of Al was detected in Cd@Al<sub>2</sub>@Zn. The peak of C 1s centered at about 284.5 eV (Fig. S9b†) came from the unavoidably introduced C-based impurity in the carbon tape.<sup>59–61</sup> It should be noticed that the XPS peaks of Cd 3d in Cd@Al<sub>2</sub>@Zn slightly shifted to lower electronic energies than those in Cd@Zn (Fig. S9c†), but the Zn signal in the two samples was detected at the same position (Fig. S9d†). Since Al<sub>2</sub>O<sub>3</sub> was directly covered on CdS, the lower-shift of the Cd signal in Cd@Al<sub>2</sub>@Zn can be attributed to the effect of Al<sub>2</sub>O<sub>3</sub> on the chemical environment of CdS, demonstrating the interaction between deposited Al<sub>2</sub>O<sub>3</sub> and the CdS core.<sup>61–63</sup> A weak signal at about 74.1 eV was detected in the Al 2p XPS of Cd@Al<sub>2</sub>@Zn (Fig. S9e†), which corresponds to the vibration of the Al–O bond,<sup>64–66</sup> confirming the presence of Al<sub>2</sub>O<sub>3</sub>. At the same time, the deposition data of ZnO in Cd@Zn and Cd@Al<sub>2</sub>@Zn determined by XPS are shown in Table S1.† It can be observed that the percentage of Zn increased from 10.11%

for Cd@Zn to 19.90% for Cd@Al<sub>2</sub>@Zn, demonstrating that the existence of Al<sub>2</sub>O<sub>3</sub> contributes to the adhesion of the ZnO layer, which is in accord with the results of XRD, TEM and SEM measurements.

**3.2.1 The passivation effect.** Fig. 3a shows the photoluminescence (PL) spectra of pristine CdS, Al<sub>2</sub>O<sub>3</sub> and Cd@Al<sub>2</sub> samples. The pristine CdS exhibited a wide region band gap emission with the center at 550 nm wavelength under the excitation of a 370 nm Xe lamp. By contrast, the pure Al<sub>2</sub>O<sub>3</sub> displayed barely an emission signal in the detected region, which may be because the Al<sub>2</sub>O<sub>3</sub> fabricated by the ALD method only responds weakly to ultraviolet light with a wavelength shorter than 350 nm (Fig. S10†). When 2 cycles of Al<sub>2</sub>O<sub>3</sub> was deposited on the surface of CdS, the obtained Cd@Al<sub>2</sub> depicted a similar PL spectrum to that of CdS but with much higher emission intensity. Considering that the pure Al<sub>2</sub>O<sub>3</sub> barely has emission in this region, the enhanced PL intensity can be attributed to the suppressed non-radiative relaxation of CdS.<sup>67–69</sup> The photocurrent response of the samples is shown in Fig. S11.† Compared with pure CdS, the Cd@Al<sub>2</sub> exhibited an observable enhancement in photocurrent density even with only 2 cycles of Al<sub>2</sub>O<sub>3</sub> deposition. The PL spectra and photocurrent response curves provide direct evidence for the surface passivation effect of the Al<sub>2</sub>O<sub>3</sub> layer. The unavoidable surface defects, dangling bonds and the adsorbates on the surface of CdS NPs inhibit the transfer of charge carriers and consume the excited energy through a non-radiative recombination process, thus quenching the band gap emission of the semiconductor materials.<sup>58,70,71</sup> The introduced Al<sub>2</sub>O<sub>3</sub> layer covers the surface of CdS, repairs the surface defects of CdS and suppresses the formation of dangling bonds and adsorbates, which result in an enhanced band gap emission.<sup>71,72</sup> In addition, Cd@Al<sub>2</sub> exhibited a longer PL decay than pure CdS (Fig. 3b), which can be ascribed to the field passivation effect of Al<sub>2</sub>O<sub>3</sub>. It was proven that the insulating amorphous Al<sub>2</sub>O<sub>3</sub> has a high fixed negative charge density (10<sup>12</sup> to 10<sup>13</sup> cm<sup>-2</sup>) due to the existence of intrinsic (O interstitials and Al vacancies) and extrinsic (interstitial H) defects,<sup>49,60,73,74</sup> which can consume the photogenerated holes and thus efficiently inhibit the recombination of holes and electrons, prolonging the lifetime of electrons.<sup>72–76</sup>

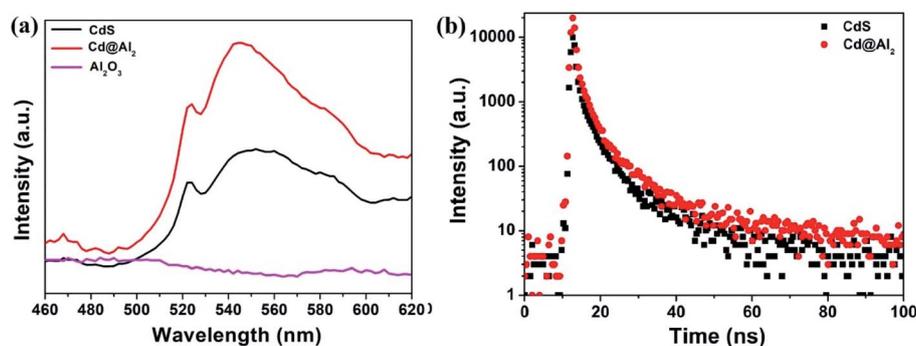


Fig. 3 (a) Photoluminescence spectra of CdS, Al<sub>2</sub>O<sub>3</sub> and Cd@Al<sub>2</sub> under the excitation of a 370 nm Xe lamp, normalized by using the UV-vis absorption ability. (b) Time-resolved photoluminescence spectra at 550 nm of CdS and Cd@Al<sub>2</sub> catalysts.

The behavior of charge carriers was further investigated to get insight into the passivation effect of the  $\text{Al}_2\text{O}_3$  layer.<sup>72,77</sup> Fig. 4 shows the Nyquist plots of  $\text{Cd@Zn}$  and  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$  under different measurement circumstances, where a smaller radius reflects relatively low transfer resistance of charge carriers.<sup>75</sup> The EIS curves measured in a dark environment (without light irradiation) give the migration resistance of charge carriers at the electrode interface while those measured under light irradiation relate to the transfer resistance of photogenerated charge carriers.<sup>65,78</sup> In addition, the EIS circles in the lower frequency region show the charge recombination resistance while those in the higher frequency region are related to the charge transfer resistance.<sup>76–80</sup>  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$  exhibited the highest charge transfer ability in the dark environment (Fig. 4a) due to the smallest Nyquist radius. Under light irradiation, the EIS curves of  $\text{Cd@Zn}$  and  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$  were recorded under different bias voltages to investigate the electrode polarized action of the measuring system. As can be observed in Fig. S12,† the radius of Nyquist semicircles decreased with the voltage increasing from 0 V to 1.2 V, indicating a continual polarization process of the samples.<sup>77,81</sup> We investigated the Nyquist curves of the samples under 0 V and 0.8 V bias voltages. As shown in Fig. 4b, when the EIS was performed under a bias voltage of 0 V, there was co-existence of recombination and transfer resistance in the catalyst.  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$  exhibited a larger Nyquist radius in the low frequency region, indicating that an increased interface charge recombination resistance was caused by the  $\text{Al}_2\text{O}_3$  interlayer. When the EIS was carried out under a bias of 0.8 V, the built-in electric field and recombination were balanced out under the drive force of the polarization effect. EIS curves give the information about charge transfer resistance. It can be observed from Fig. 4c that the Nyquist radius of  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$  was smaller than that of  $\text{Cd@Zn}$ , indicating that the  $\text{Al}_2\text{O}_3$  interlayer is beneficial for the decrease of charge transfer resistance. Based on the analysis, it can be concluded from the EIS results that the introduced ultrathin  $\text{Al}_2\text{O}_3$  interlayer helps to suppress interfacial charge recombination and facilitate the separation of electron–hole pairs, which is very much desired for an efficient photocatalytic reaction.

In order to verify the effect of  $\text{Al}_2\text{O}_3$  on the capture of photogenerated holes, we further carried out the  $\text{H}_2$  evolution

reaction in the absence of an extra sacrificial agent, and the results are shown in Fig. S13.† It can be observed that the pure CdS showed barely photocatalytic activity in the absence of a sacrificial agent due to the serious charge carrier recombination. The  $\text{Cd@Al}_2\text{O}_3$  displayed a  $\text{H}_2$  evolution rate of  $3.09 \mu\text{mol h}^{-1}$ , and the improved  $\text{H}_2$  evolution rate indicates that the  $\text{Al}_2\text{O}_3$  ultrathin layer suppresses the recombination of photogenerated charge carriers in CdS. Assuming that the photogenerated holes can't be consumed by  $\text{Al}_2\text{O}_3$ ,  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$  would display a slightly enhanced  $\text{H}_2$  evolution rate in comparison with  $\text{Cd@Zn}$  due to the improved interface contact. But in fact, much higher  $\text{H}_2$  evolution ability was observed over  $\text{Cd@Al}_2\text{O}_3\text{@Zn}$ . The dramatically enhanced  $\text{H}_2$  evolution rate indicates that the  $\text{Al}_2\text{O}_3$  consumes the photogenerated holes and plays the role of a sacrificial agent. Fig. S13b† shows the comparison of the  $\text{H}_2$  evolution ability of the catalysts with and without an additional sacrificial agent. We noticed that the photocatalytic activity of the  $\text{Cd@Al@Zn}$  catalyst in aqueous solution was much lower than that with the addition of an external sacrificial agent even though  $\text{Al}_2\text{O}_3$  plays the role of a sacrificial agent. This is because the addition of  $\text{Na}_2\text{S}$  and  $\text{Na}_2\text{SO}_3$  not only provides hole acceptors but also results in an alkaline environment with a high pH value for the photocatalytic reaction, which contributes to the splitting of water.<sup>82–84</sup>

Based on the above results and discussion, the surface passivation effect of  $\text{Al}_2\text{O}_3$  can be understood to be both chemical passivation and field passivation, as follows: (1) amorphous  $\text{Al}_2\text{O}_3$  synthesized by the ALD strategy has good compatibility and covering capacity on the substrate materials, which help to repair the surface defects and stop the formation of harmful dangling bonds, thus efficiently improving the quality of substrate materials and inhibiting the charge recombination induced by defects through the chemical passivation effect;<sup>50,59,80</sup> (2) due to the existence of fixed negative charges,<sup>73,74</sup> amorphous  $\text{Al}_2\text{O}_3$  works as a hole trapper, and can effectively consume the photogenerated holes in semiconductors, thus suppressing the recombination of charge carriers by the means of a built-in electric field passivation effect.<sup>58–60,73,74</sup>

In order to clarify the contribution of the introduced  $\text{Al}_2\text{O}_3$  passivation layer to the improvement of  $\text{H}_2$  evolution ability, we

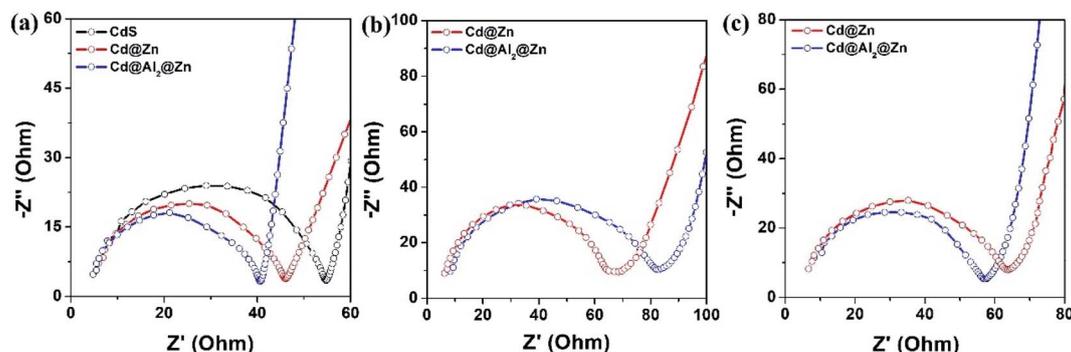


Fig. 4 EIS curves of the catalysts measured at different circumstances: (a) at 0 V bias voltage in the dark; (b and c) at 0 V and 0.8 V bias voltage under 300 W Xe lamp irradiation, respectively.

carefully studied past reports about the interface-passivated design and the performance enhancement of these systems, and the results are listed in Table S3.† Due to the special ability to shrink the interface defects and improve the performance of semiconductor based materials, the strategy of passivation layer has been widely utilized in various fields including photoelectrocatalytic water splitting, CO<sub>2</sub> reduction, solar cells, LED devices and so on in the past ten years. However, barely photocatalytic H<sub>2</sub> evolution has been reported. As was observed, all of the reported systems displayed improved performance, and the strengthening of interface contact and decreased interface defects contribute the most to the enhancement. In our work, the Cd@Al@Zn system exhibited 2.2 times higher H<sub>2</sub> evolution rate than Cd@Zn, and the enhancement coefficient is similar to that in the case of utilization of a passivation layer in other systems. The result demonstrates that we have carried out the first successful attempt to improve the photocatalytic H<sub>2</sub> evolution ability through passivating the interlayer of two semiconductors, which is meaningful in the exploration of high activity photocatalysts.

**3.2.2 The tunneling effect.** The study on the thickness of the Al<sub>2</sub>O<sub>3</sub> interlayer is to reveal the inner charge transfer mechanism in Cd@Al@Zn. Fig. S14a and b† show the EIS and photocurrent response curves of the samples, respectively. In comparison to Cd@Zn, Cd@Al<sub>1</sub>@Zn and Cd@Al<sub>2</sub>@Zn displayed obviously decreased charge transfer resistance and increased photocurrent response ability, implying that the charge carriers can transfer through the Al<sub>2</sub>O<sub>3</sub> interlayer *via* the tunneling effect when the Al<sub>2</sub>O<sub>3</sub> layer is thin enough.<sup>60,71</sup> When the number of deposition cycles of Al<sub>2</sub>O<sub>3</sub> was increased to 6, increased transfer resistance and decreased photocurrent response were observed in Cd@Al<sub>6</sub>@Zn, suggesting that the tunneling effect is decreased due the excessive thickness of the Al<sub>2</sub>O<sub>3</sub> layer. According to the equation<sup>63,79</sup>  $k_{ET} = k_{ET}^0 \times e^{-\beta d}$  (where  $k_{ET}$ ,  $\beta$  and  $d$  are, respectively, the electron transfer rate, tunneling parameter, and the thickness of the Al<sub>2</sub>O<sub>3</sub> layer), when a thicker Al<sub>2</sub>O<sub>3</sub> interlayer is introduced, the carrier transfer through the tunneling effect will decline exponentially. Therefore, the over-deposition of the Al<sub>2</sub>O<sub>3</sub> interlayer is harmful to the efficient migration of photogenerated electrons.

Mott–Schottky (M–S) plots were collected to ascertain the band edge potentials of the samples. As shown in Fig. S15,† the observed positive slope confirmed the n-type characteristic of the samples.<sup>14,18</sup> It can be determined from the M–S curves that the CB of Cd@Al<sub>2</sub> shifted to a more negative region compared to that of pure CdS, which can be ascribed to the field passivation effect caused by Al<sub>2</sub>O<sub>3</sub>. Under irradiation, the photogenerated holes on the VB of CdS were transferred to the Al<sub>2</sub>O<sub>3</sub> layer because of the presence of the negative fixed charges on the surface of amorphous Al<sub>2</sub>O<sub>3</sub>, while the photogenerated electrons quickly accumulated on the CB of CdS, thus resulting in an upward band bending.<sup>50,66</sup>

The energy band structures of CdS, Al<sub>2</sub>O<sub>3</sub> and ZnO were further determined by DFT calculations to clarify the transfer behavior of charge carriers in the Cd@Al@Zn system. The Al<sub>2</sub>O<sub>3</sub> displayed a deep valence band (VB) position and a high conductive band (CB) position compared to both CdS and ZnO

(Fig. 5a), indicating its insulating nature. Due to the existence of negative fixed charges,<sup>57,63,65</sup> the holes on the VB of ZnO and CdS can be consumed by Al<sub>2</sub>O<sub>3</sub>, thus inhibiting the formation of the Z-scheme energy transfer mechanism. The work function of CdS is lower than that of ZnO, and thus the electrons on the CB of CdS will cross over Al<sub>2</sub>O<sub>3</sub> through the tunneling effect and finally jump to the CB of ZnO.<sup>59,64,66</sup> The transfer speed of electrons can be determined from the contact potential difference (CPD), which reflects the driving force of charge transfer at the interface between two semiconductors, and a higher CPD means a faster electron transfer speed from one material to another.<sup>65,66,85</sup> The CPD in CdS–ZnO (4.19 eV, Fig. 5b) is much lower than that in CdS–Al<sub>2</sub>O<sub>3</sub> (6.69 eV, Fig. 5c) and Al<sub>2</sub>O<sub>3</sub>–ZnO (7.83 eV, Fig. 5d), implying that the existence of Al<sub>2</sub>O<sub>3</sub> helps to trap electrons from CdS and then transfer to the surface of ZnO, realizing the transfer of electrons from the CdS core to the ZnO shell. As a result, the Al<sub>2</sub>O<sub>3</sub> interlayer traps the photogenerated electrons from the CdS core and then transfers them to the ZnO shell by the tunneling effect, and meanwhile the photogenerated holes produced in the CdS core and the ZnO shell are efficiently consumed by the negative fixed charges of the Al<sub>2</sub>O<sub>3</sub> interlayer, resulting in a modified type II transfer mechanism.

The tunneling effect of the ultrathin Al<sub>2</sub>O<sub>3</sub> interlayer was further determined by the detection of active intermediate hydroxyl radicals ( $\cdot\text{OH}$ ) over Cd@Zn and Cd@Al<sub>2</sub>@Zn. Ammonium carbonate (AC), a hole scavenger, was added to the suspension of Cd@Zn (denoted as Cd@Zn + AC for convenience) to distinguish the hole trapping ability of the Al<sub>2</sub>O<sub>3</sub> interlayer in Cd@Al<sub>2</sub>@Zn.<sup>86,87</sup> PL spectra using terephthalic acid (TA) as a probe molecule were measured to detect the formation of  $\cdot\text{OH}$  in the suspension, where a higher PL peak intensity reflects a larger amount of  $\cdot\text{OH}$ . In general, the  $\cdot\text{OH}$  can be generated from both the photogenerated electrons and holes when the CB and VB are in suitable positions.<sup>66,86,87</sup> (1) The photogenerated electrons react with the dissolved oxygen to produce superoxide radical anions ( $\cdot\text{O}_2^-$ ), which further react with hydrogen ions (H<sup>+</sup>) to produce  $\cdot\text{OH}$  radicals. (2) The photogenerated holes directly react with hydroxyl ions (OH<sup>-</sup>) or water molecules (H<sub>2</sub>O) to produce  $\cdot\text{OH}$  radicals. As shown in Fig. 6a, in Cd@Zn (sample I), the electrons on the CB of CdS and the holes on the VB of ZnO contribute to the generation of  $\cdot\text{OH}$ , and a relatively high PL intensity can be observed over Cd@Zn (Fig. 6b). In Cd@Zn + AC (sample II), the holes on the VB of CdS cannot produce  $\cdot\text{OH}$  radicals due to its less positive position than OH<sup>-</sup> and H<sub>2</sub>O, and the holes on the VB of ZnO can be trapped by AC, so the  $\cdot\text{OH}$  radicals are only generated from the electrons on the CB of CdS and ZnO, resulting in a much lower PL intensity in comparison with Cd@Zn. In Cd@Al<sub>2</sub>@Zn (sample III), the Al<sub>2</sub>O<sub>3</sub> interlayer plays a similar role to AC; the photogenerated holes on the VB of both CdS and ZnO are consumed by its negative fixed charges, which inhibits the generation of  $\cdot\text{OH}$  from the VB of the two semiconductors. Due to the tunneling effect caused by the ultrathin Al<sub>2</sub>O<sub>3</sub> interlayer, the photogenerated electrons can be transferred quickly from the CB of CdS to the CB of ZnO, which is adverse to the formation of  $\cdot\text{OH}$  because the electrons on the CB of ZnO have a lower reducing capability than the electrons on the CB of CdS

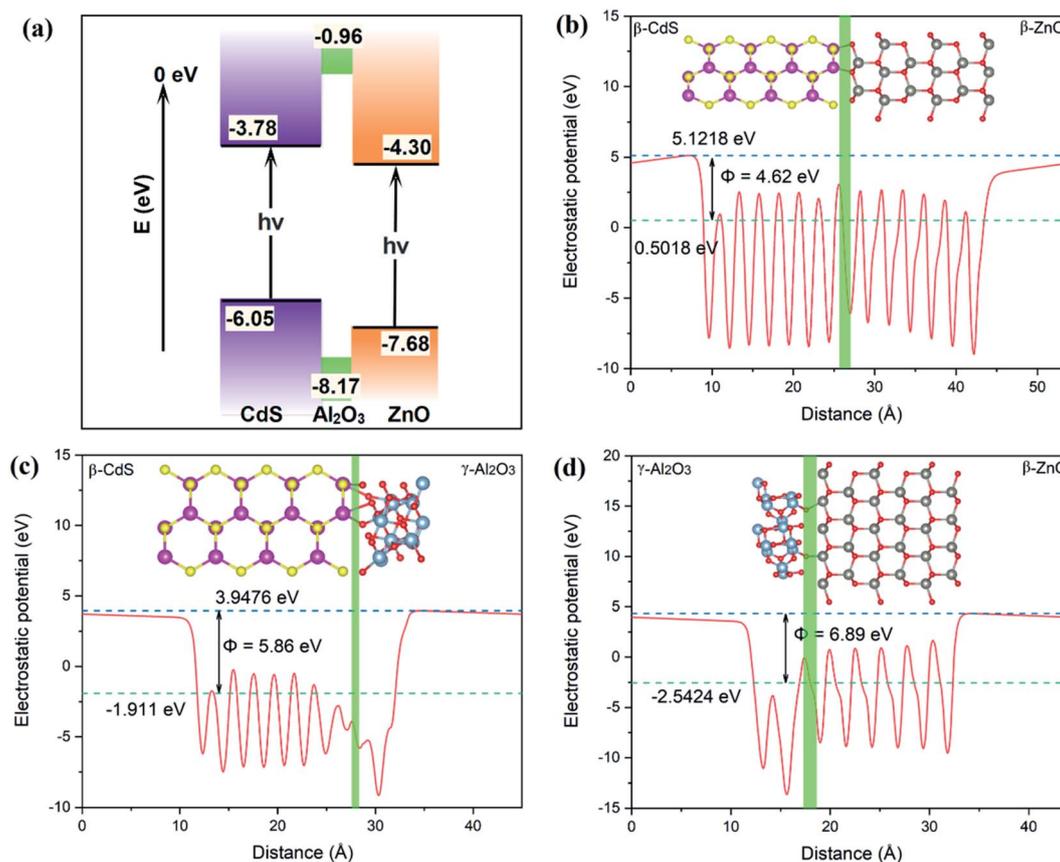


Fig. 5 (a) The calculated energy band position of CdS, Al<sub>2</sub>O<sub>3</sub> and ZnO; (b)–(d) the calculated slab module of the energy change along the *c*-direction of the slab cell for CdS–ZnO, CdS–Al<sub>2</sub>O<sub>3</sub> and Al<sub>2</sub>O<sub>3</sub>–ZnO interfaces, respectively,  $\phi$  is the calculated contact potential difference.

(the CB position of ZnO is more positive than that of CdS). Therefore, the PL intensity produced from Cd@Al<sub>2</sub>@Zn will be lower than that from Cd@Zn + AC, as confirmed in Fig. 6b.

### 3.3 The proposed transfer mechanism of charge carriers

Based on the results and discussion, we can draw a conclusion that the existence of the Al<sub>2</sub>O<sub>3</sub> interlayer in Cd@Al@Zn results in a new transfer mechanism of charge carriers, which is totally different from the Z-scheme mechanism in the Cd@Zn system (Scheme 1a). As illustrated in Scheme 1b, after the introduction of Al<sub>2</sub>O<sub>3</sub>, the holes on the VB of both CdS and ZnO are consumed

by the negative fixed charges of Al<sub>2</sub>O<sub>3</sub>, which inhibits the direct Z-scheme mechanism. Due to the different energy band positions of CdS and ZnO, the photogenerated electrons transfer from the CB of CdS (higher energy level) to the CB of ZnO (lower energy level), thus forming a modified type II transfer mechanism, which is also different from the traditional type II mechanism because the photogenerated holes do not transfer from the VB of ZnO to that of CdS. The photogenerated electrons transfer along the direction of CdS–Al<sub>2</sub>O<sub>3</sub>–ZnO, which efficiently suppresses the recombination of charge carriers. Moreover, the photogenerated electrons transferred from the CdS core to the ZnO shell can be consumed effectively in time by H<sub>2</sub>O, thus facilitating the

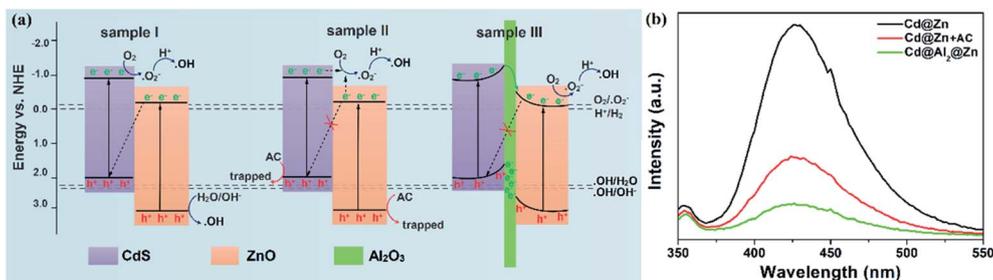
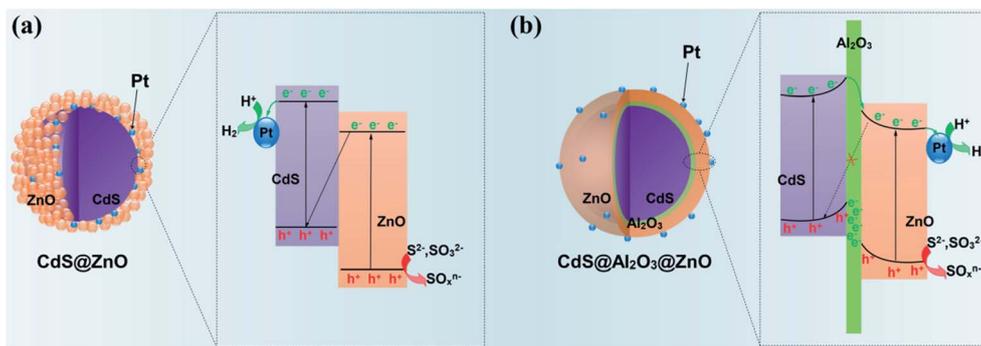


Fig. 6 (a) Pathways of  $\cdot\text{OH}$  produced from Cd@Zn (sample I), Cd@Zn + AC (sample II) and Cd@Al<sub>2</sub>@Zn (sample III); (b) the corresponding PL spectra of the samples.



Scheme 1 The possible energy transfer mechanism in (a) Cd@Zn and (b) Cd@Al@Zn.

generation of  $H_2$ . The formation of the modified type-II energy transfer mechanism can also be verified by the discordance between the electron behavior and  $H_2$  evolution ability. It is observed from Fig. S14<sup>†</sup> that Cd@Al<sub>6</sub>@Zn presents a smaller photocurrent and a larger electrochemical impedance than Cd@Zn, but the  $H_2$  evolution rate of Cd@Al<sub>6</sub>@Zn is higher than that of Cd@Zn. The increased  $H_2$  production performance can be ascribed to the exposure of the active reaction surface. In Cd@Zn, a direct Z-scheme energy transfer mechanism is proven, where the photogenerated electrons accumulate on the surface of the CdS core. However, the surface of CdS is partially covered by ZnO nanoparticles, and the contradiction in the exposure of the CdS surface and the formation of a heterojunction between CdS and ZnO is irreconcilable and thus results in a limited  $H_2$  evolution rate. As to Cd@Al@Zn, the photogenerated electrons transferred to the surface of the ZnO shell through the modified type-II mechanism, so the exposure of active reaction sites and the formation of the heterojunction are complementary, thus accelerating the photocatalytic reaction.

## 4. Conclusions

In summary, an ultrathin  $Al_2O_3$  interlayer was introduced into the interface between the CdS core and the ZnO shell using a facile ALD method. The obtained Cd@Al<sub>x</sub>@Zn displayed dramatically enhanced photocatalytic hydrogen evolution performance. The optimized Cd@Al<sub>2</sub>@Zn presented a photocatalytic hydrogen evolution rate of  $1190 \mu\text{mol h}^{-1} \text{m}^{-2}$ , which was 2.17 times higher than that of Cd@Zn. The enhanced photocatalytic performance can be attributed to the following reasons: (i) the tunneling effect of the  $Al_2O_3$  interlayer induces the transfer of photogenerated electrons from the CdS core to the ZnO shell, and then the electrons are consumed by  $H_2O$  to produce  $H_2$ ; meanwhile the photogenerated holes are depleted by the fixed negative charges of  $Al_2O_3$ , forming a modified type II transfer mechanism. (ii) The surface passivation effect of the ultrathin  $Al_2O_3$  interlayer overcomes the lattice mismatch between wurtzite ZnO and wurtzite CdS, which facilitates the uniform deposition of the ZnO layer and promotes the close contact between interfaces. Benefitting from this, the separation and transfer of charge carriers are efficiently promoted, thus boosting the photocatalytic  $H_2$  generation.

## Conflicts of interest

The authors declare no conflict of interest.

## Acknowledgements

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