Solar cells Efficient, stable solar cells by using inherent bandgap of α -phase formamidinium lead iodide

Hanul Min, Maengsuk Kim, Seung-Un Lee, Hyeonwoo Kim, Gwisu Kim, Keunsu Choi, Jun Hee Lee, Sang II Seok*

In general, mixed cations and anions containing formamidinium (FA), methylammonium (MA), caesium, iodine, and bromine ions are used to stabilize the black α -phase of the FA-based lead triiodide (FAPbl₃) in perovskite solar cells. However, additives such as MA, caesium, and bromine widen its bandgap and reduce the thermal stability. We stabilized the α -FAPbl₃ phase by doping with methylenediammonium dichloride (MDACl₂) and achieved a certified short-circuit current density of between 26.1 and 26.7 milliamperes per square centimeter. With certified power conversion efficiencies (PCEs) of 23.7%, more than 90% of the initial efficiency was maintained after 600 hours of operation with maximum power point tracking under full sunlight illumination in ambient conditions including ultraviolet light. Unencapsulated devices retained more than 90% of their initial PCE even after annealing for 20 hours at 150°C in air and exhibited superior thermal and humidity stability over a control device in which FAPbl₃ was stabilized by MAPbBr₃.

erovskite solar cells (PSCs) have increased power conversion efficiencies (PCEs) with broader solar-light absorption through narrower bandgaps. Among lead halide perovskites (LHPs), the narrowest bandgap is afforded by the formamidinium (FA)-based lead triiodide (FAPbI₃) [1.45 to 1.51 eV (1-3) in thin films], which also affords improved thermal stability (4) relative to MAPbI₃ (where MA is methylammonium) because of its elevated decomposition temperature. However, FAPbI₃ readily transforms from the desired trigonal black α -phase into the undesired wide-bandgap δ -phase with hexagonal symmetry under ambient conditions at room temperature (5). We have shown that modification of FAPbI₃ with MAPbBr₃ stabilized the α-FAPbI₃ phase and subsequently improved the device efficiency to >18% in (FAPbI₃)_{0.85}(MAPbI₃)_{0.15} (6).

Efforts to improve the stability of FAPbI₃ have focused on mixed cation-anion hybrid LHPs that incorporate several cations, anions, or both, such as the FA_xMA_{1-x} double cation or the $FA_{1-x-y}MA_xCs_y$ triple cation. This incorporation is achieved by the partial replacement of FA^+ , MA^+ , Rb^+ , Cs^+ , Br^- , and I^- (7–11). Nevertheless, when MA and Br are alloyed in FAPbI₃, various problems such as low thermal stability (4) caused by the addition of MA. phase separation (12) caused by the presence of mixed halides, and reduced photon absorption arise, thereby resulting in low current density owing to an undesirable increase in the bandgap. Although the α -FAPbI₃ phase without MA can be stabilized by incorporating phenylethylammonium lead iodide (13) through surface functionalization (14) or by using both Rb and Cs (15), the resulting PCE is still low when compared with that obtained with the use of the FA_xMA_{1-x} double cation.

To further increase the PCE by enhancing the photocurrent density resulting from increased light harvesting, a new composition is needed that stabilizes the α -phase while maintaining the inherent bandgap of FAPbI₃. The ionic radii and molecular configuration of methylenediammonium [+H₃N-CH₂-NH₃+ (MDA), 262 picometers (pm), calculated] and FA (⁺H₂N=CH-NH₂, 256 pm) are comparable except for the valence-state difference. However, MDA has more hydrogen atoms than FA and even MA, which means that it can form a greater number of H bonds with I⁻. Thus, MDA could structurally stabilize the α -FAPbI₃ in even smaller amounts than does MA. Some studies have reported the incorporation of divalent cations into monovalent A-site cations (16-21), but most of the divalent diammonium cations reported have been used to realize twodimensional (2D) or 2D/3D LHPs. Safdari et al. (16) reported on the 2D diammonium perovskites [NH₃(CH₂)₄NH₃]PbI₄, [NH₃(CH₂)₆NH₃] PbI₄, and [NH₃(CH₂)₈NH₃]PbI₄, which demonstrated PCEs of ~1%. The recent synthesis of $(NH_3C_mH_{2m}NH_3)(CH_3NH_3)_{n-1}Pb_nI_{3n+1}$ (m = 4-9/n = 1-4) that used complementary solidstate grinding and solution methods as a 2D/ 3D mixture has also been reported (21). The use of 2D or 2D/3D mixed-perovskite-containing diammonium cations increased the stability of perovskite materials by strengthening electrostatic interactions between the layers that rigidify the structure. However, larger organic cations in 2D structures can inhibit charge transport by creating insulating spacing layers between conductive inorganic slabs, so it would be desirable to fabricate 3D LHPs without the 2D perovskite layer.

We report highly efficient and stable PSCs fabricated with MA-, Cs-, and Br-free α -FAPbI₃ by incorporating a small amount of MDACl₂. The resulting FAPbI₃:3.8 mole % (mol %) MDACl₂ structure stabilized the α -phase with only a slight change in the bandgap. We achieved a short-circuit current density (J_{sc}) of between 26.1 and 26.7 mA cm⁻² (certified), which is the highest reported for PSCs fabricated from FA-based LHPs. The device exhibited a PCE of >24% (certified stabilized 23.7%) and long-term stability compared with the control [(FAPbI₃)_{0.95}(MAPbI₃)_{0.05}], which itself shows the highest efficiency among mesoporous (mp)-TiO₂-based PSCs.

We deposited a thin film of FAPbI₂ incorporating MDACl₂ with a process similar to that reported previously for state-of-art mixed perovskites (9-11) but adding MDACl₂ instead of MAPbBr3 (see the experimental section in the supplementary materials). The ultraviolet-visible (UV-vis) absorption spectra of $FAPbI_3: xMDACl_2 (x = 0, 1.9, 3.8, and 5.7 mol \%)$ slightly blue-shifted with increasing amounts of MDACl₂ (Fig. 1A) and were nearly identical to that of a thin film without the addition of MACl (fig. S1) as a mediator (22-26) for high-crystallinity perovskite. This result is consistent with the corresponding shifts of photoluminescence (PL) emission peaks (Fig. 1A and fig. S2) from 826 nm to 824, 822, and 820 nm, and 816 nm for the control. This result suggests that MACl almost completely disappeared during the annealing process, as previously reported (27, 28). The presence of MDA in the perovskite was confirmed by Fourier-transform infrared spectroscopy (fig. S3) and nuclear magnetic resonance imaging (fig. S4). Thus, MDACl₂ may have been incorporated into the perovskite lattice within the experimental x-value range.

The bandgap changes in FAPbI₃ caused by the introduction of MDACl₂ were calculated using density functional theory and compared for two cases as per Eqs. 1 and 2:

$$\begin{aligned} & \text{FA vacancy }:(\text{FA}_{1-2x}^{+1}\text{MDA}_{x}^{+2})\text{Pb}^{2+}\text{I}_{3}^{-1}\text{Cl}_{2x}^{-1} \rightarrow \\ & [(\text{FA}_{1-2x}^{+1}\text{MDA}_{x}^{+2})\text{Pb}^{2+}\text{I}_{3}^{-1}] + V_{x\text{FA}} + 2x\text{FACl} \end{aligned} \tag{1} \\ & \text{Cl}_{i} \text{ insertion }:(\text{FA}_{1-x}^{+1}\text{MDA}_{x}^{+2})\text{Pb}^{2+}\text{I}_{3}^{-1}\text{Cl}_{2x}^{-1} \rightarrow \end{aligned}$$

$$[(\mathbf{FA}_{1-x}^{+1}\mathbf{MDA}_{x}^{+2})\mathbf{Pb}^{2+}\mathbf{I}_{3}^{-1}\mathbf{Cl}_{x}^{-1}] + \mathbf{Cl}_{xi} + x\mathbf{FACI}$$
(2)

Equations 1 and 2 represent compositions that assume the FA cation vacancies and the insertion of Cl⁻ with a small ionic radius, respectively, upon the addition of MDACl₂. The bandgap (1.47 eV) with the FA vacancy ($V_{\rm FA}$) composition as per Eq. 1 is slightly greater than that of pristine FAPbI₃ (1.45 eV) and the Cl interstitial composition (Cl_i) based on Eq. 2 yielded an increased bandgap of 1.69 eV (Fig. 1B). We

Department of Energy Engineering, School of Energy and Chemical Engineering, Ulsan National Institute of Science and Technology, 50 UNIST-gil, Eonyang-eup, Ulju-gun, Ulsan 44919, Korea.

^{*}Corresponding author. Email: seoksi@unist.ac.kr

expected that the addition of MDACl₂ to FAPbI₃ would mainly result in FA defects and the possible induction of PL quenching. However, adding 3.8 mol % MDACl₂ to FAPbI₃ enhanced the PL quantum yield as measured with an integrated sphere, but was greatly reduced with further MDACl₂ addition. This result implies that the FA defects did not act as deep electron traps. As reported previously (4, 29), pristine FAPbI₂ thin films annealed at high temperatures exhibited a black α -phase that absorbed long-wavelength light. However, the α -phase transitions to the yellow δ -phase within 10 days at room temperature and within 1 day under high-humidity conditions (5, 30) because the metastable polymorphs stabilized at high temperature are preferably converted back to the thermodynamically stable phase below 120°C.

The structural stabilization of the α -phase in FAPbI₃ with added cations can be explained by several factors. First, the α -phase stabilization by smaller cations such as Cs⁺ can be understood from the Goldschmidt tolerance factor *t* approaching 0.9, which is similar to that of MAPbI₃, by mixing FAPbI₃ (*t* ~ 1) and CsPbI₃ (*t* ~ 0.8) (*31*). Second, cation mixing in the FA sites affords entropic stabilization through the resulting entropic gain and small internal energy input to form their solid solution (*32*). Third, the stabilization of α -FAPbI₃ with MA cations can be explained by the presence of strong H bonds between the I⁻ and H–N groups (*33–35*). The α -FAPbI₃ phase can be more easily stabilized when MA, Cs, or both is used with Br⁻ of small ionic radii, and highefficiency PSCs are normally fabricated using α -FAPbI₃ stabilized by mixed cations and anions.

We expected that MDA could stabilize α -FAPbI₃ through the partial replacement of FA sites because MDA has more H groups with an ionic radius similar to that of FA and a stronger ionic interaction of its divalent state. The x-ray diffraction (XRD) patterns of FAPbI₃:*x*MDACl₂ (*x* = 0, 1.9, 3.8, and 5.7 mol %) and the control layers exposed to 80% humidity for 24 hours after annealing of the coatings of precursor solution at 150°C for 10 min (Fig. 1C) show two peaks characteristic of the α -FAPbI₃ phase at 14.3° and 28.6° assigned to the (001) and (002) crystal planes, respectively, and one peak at 11.6° corresponding to the δ -phase. The as-annealed films (fig. S5) exhibited a pure α -phase for all compositions, but after exposure to high humidity for 24 hours, pure $FAPbI_3$ completely converted to the δ -phase and FAPbI3 with 1.9 mol % MDACl2 exhibited a strong phase transition, with the control also exhibiting a certain amount of the δ -phase.

By contrast, FAPbI₃ incorporating 3.8 and 5.7 mol % of MDACl₂ retained the α -phase. Furthermore, as shown in fig. S5, the overall



Fig. 1. Preparation of FAPbl₃:xMDACl₂ (x = 0, 1.9, 3.8, and 5.7 mol %). (A) UV-vis absorption and PL spectra of perovskite layers with different *x* values. The perovskite layer based on 0.95FAPbl₃/ 0.05MAPbBr₃ as a control was also investigated. (B) Electronic band structures of (a) FAPbl₃, (b) representative FA_{0.963}(V_{FA})_{0.037}MDA_{0.037}Pbl₃, and (c) representative FA_{0.963}MDA_{0.037}Pbl₃(Cl_i)_{0.037}. The doping amount of 0.037 in the calculation was given by 1/27 (= 0.03704) based on the 3 × 3 × 3 supercell for the composition close to the actual experiment. (C) XRD patterns of perovskite prepared with different *x* values and the control layer exposed to 80% humidity for 24 hours after annealing of the coatings of the precursor solution at 150°C for 10 min. (D) GIWAXS 2D images at full depth of x-ray incident angle for perovskite layers for (a) x = 3.8 mol % and (b) control.

signal intensity increased for MDACl₂ addition up to 3.8 mol % without any impurity peaks. The incorporation of >3.8 mol % MDACl₂ may result in increased crystallinity because of the reduced lattice strain at FA defects (fig. S6). As can be inferred from Eqs. 1 and 2, the perovskite structure may contain defective FA⁺ or interstitial Cl⁻ to satisfy the electrical neutrality rule with the release of FACl as a by-product when MDA²⁺ ions are substituted into the FA sites in FAPbI₃.

The addition of MDACl₂ likely formed FACl or substituted residual MDACl₂. However, the final annealed perovskite films did not exhibit any impurity peaks such as those of FACl and MDACl₂. This result implies that MDACl₂ was successfully incorporated into the FAPbI₃ perovskite lattices, and the resulting FACl was eliminated by annealing at 150°C for 10 min, because FACl easily volatilized during heat treatment. The structural phase of FAPbI₃ with and without the substitution of the representative 3.8 mol % MDA and the presence of other phases such as FACl or MDACl2 was further characterized using grazing-incidence wide-angle x-ray scattering (GIWAXS). In Fig. 1D, we observed diffraction rings assigned to α-FAPbI₃-(100)_c, α-FAPbI₃-(200)_c, α-FAPbI₃-(210)_c, δ-FAPbI₃-(100)_h, and PbI₂-(001)_t for the two representative samples (36). In addition, different crystal orientations of α -FAPbI₃, [100]_c and [200]_c, were observed as preferential GIWAXS peaks for both samples. The PbI₂ components remaining in both perovskite layers showed similar out-of-plane orientations. The GIWAXS ring patterns of α -FAPbI₃:xMDACl₂ (x = 0 and 3.8 mol %) did not exhibit appreciable differences, and the fitted azimuthal circular average GIWAXS 1D full spectra were also nearly identical, without showing any peaks related to FACl or MDACl₂. This result indicated that MDA is substituted into the FAPbI3 lattice, as estimated from the conventional XRD analysis.

Point defects corresponding to $V_{\rm FA}$ or $Cl_{\rm i}$ expressed by Eqs. 1 or 2, respectively, could decrease the open-circuit voltage (V_{OC}) and the fill factor (FF) rather than the current density $J_{\rm SC}$, although $V_{\rm FA}$ defects form shallow traps near the conduction band (37). Figure 2A compares the PCE distributions of the PSCs fabricated with FAPbI₃:xMDACl₂ (x = 0, 1.9, 3.8, and 5.7 mol %) and the control. The average PCE values of the PSCs fabricated with no MDACl₂ improved from $22.012 \pm 0.51\%$ to 22.46 ± 0.34% for 3.8 mol % MDACl₂, mainly from an increase in J_{SC} , while exhibiting similar or greater $V_{\rm OC}$ and FF values (fig. S7). When more MDACl₂ (5.7 mol %) was added, the degradation of crystallinity and PL caused the PCE to slightly decline from a decrease in both J_{SC} and $V_{\rm OC}$. The initial PCE of the PSC fabricated without MDACl₂ was fairly high (29) but decreased substantially over time because of the



Fig. 2. Device performance. (**A**) PCE of perovskite solar cells fabricated with different *x* values of FAPbl₃:*x*MDACl₂ (x = 0, 1.9, 3.8, and 5.7 mol %) and the control. (**B** and **C**) Comparison of (B) *I–V* curves and (C) external quantum efficiency between target (x = 0.38 mol %) and control. (**D**) Comparison of surface morphologies and cross-sectional images of target and control acquired with scanning electron microscopy.

 α - to δ -phase transition (fig. S8). Therefore, in terms of efficiency and phase stability, we fixed 3.8 mol % in an appropriate amount and compared its characteristics with the control. We passivated the surface of the target and control layers by means of previously reported methods (9, 11, 38, 39), including ours (40). Figure 2B compares the current density-voltage (J-V) characteristics, reverse and forward bias sweep, for one of the "bestperforming" PSCs fabricated with 3.8 mol % MDACl₂ (denoted as the target) and the control. The J_{SC} , V_{OC} , and FF values (table S1) calculated from the J-V curves of the target were estimated as 26.50 mA cm^{-2} , 1.14 V, and 81.77%, respectively: these correspond to a PCE of 24.66% under standard air mass (AM) 1.5 conditions, mainly the result of the very high J_{SC} value. The control exhibited a PCE of 23.05%, with $J_{\rm SC}$ = 25.14 mA cm $^{-2},$ $V_{\rm OC}$ = 1.14 V, and FF = 80.55%. As expected, the surface passivation of perovskite layers improved both $V_{\rm OC}$ and FF, but J_{SC} remained almost unchanged, which resulted in a PCE of >24%. The external quantum efficiency comparison of the control and target in Fig. 2C showed that the efficiency improvement of the PSC with MDACl₂ was the result of the expansion of the range of absorption wavelengths.

The improvement in PCE with a suitable amount of MDA meant that the negative influence of $V_{\rm FA}$ defects was compensated by other beneficial factors. Because the PCE of



Fig. 3. Changes of defects and Cl⁻ ions resulting from the doping of MDACl₂. (A) Dark *I–V* curves of electron-only devices and (B) PL decay (on quartz substrate) of perovskite layers of FAPbl₃:xMDACl₂ (*x* = 0, 1.9, 3.8, and 5.7 mol %) and control. (C) X-ray photoelectron spectroscopy of Cl 2*p* and (D) time-of-flight secondary ion mass spectrometry of Cl anion results of FTO/Bl-TiO₂/mp-TiO₂/perovskite (a) target and (b) control layer (measurement area 50 µm by 50 µm). (E) XRD patterns of perovskite layers of FAPbl₃:xMDACl₂ (*x* = 0, 1.9, 3.8, and 5.7 mol %) and control magnified (001) orientation peaks.



Fig. 4. Long-term stability test. Comparison of **(A)** humidity (85% RH, 25°C) and **(B)** thermal (150°C at ~25% RH) stability performances of unencapsulated control and target. **(C)** Maximum power point tracking measured with the encapsulated target device under full solar illumination (AM 1.5 G, 100 mW/cm² in ambient conditions) without a UV filter.

PSCs depends on the surface morphology of the perovskite layers, we compared the surface roughness and grain sizes of the target and the control with scanning electron microscopy. We found no notable differences in the cross-sectional thickness of the two representative layers, the control and the target (Fig. 2D), so the introduction of MDACl₂ into the precursor of FAPbI3 did not affect the features of the perovskite layers. The performance factors for the device shown in Fig. 2B, certified by an accredited laboratory (Newport, USA) by means of the quasi-steady-state (QSS) method, a newly established measurement standard for certification, were J_{SC} = 26.1 mA cm⁻², $V_{\rm OC}$ = 1.15 V, and FF = 79.0%, which corresponds to a stabilized PCE of 23.73%, the highest reported for devices using mp-TiO₂ as an electrode (fig. S9). Another certified device recorded a very high $J_{\rm SC}$ (26.70 mA cm⁻²) value, which is the highest reported in FAPbI3based PSCs, along with $V_{\rm OC}$ = 1.144 V and FF = 77.56%, corresponding to a stabilized PCE value of 23.69% (fig. S10). As per the certification process, the $J_{\rm SC}$ and $V_{\rm OC}$ values obtained from the QSS method were similar to the corresponding ones of the I-V measurements conducted in the reverse-bias mode; however, the FF was reduced.

The presence of $V_{\rm FA}$ defects may have resulted in an increase in the electron-carrier density in the conduction band in FAPbI₃,

which can be expressed using Eq. 3:

$$(FA_{1-x}^{+1}MDA_x^{+2})Pb^{2+}I_3^{-1}CI_x^{-1} \rightarrow [(FA_{1-x}^{+1}MDA_x^{+2})$$

 $Pb^{2+}I_3^{-1}]^{x+} + xCl^{-} \rightarrow [(FA_{1-x}^{+1}MDA_x^{+2})$
 $Pb^{2+}I_3^{-1}]^{x+} + xe_{CB}^{-} + \frac{x}{2}Cl_2$ (3)

Defects in semiconductors can capture and trap free charges, and thus the electrical properties of the material can change after all the defects are "filled" with charge in the trap-fill limit. To quantitatively assess defect density, we fabricated an electron- and hole-only device with the configuration of fluorine-doped tin oxide(FTO)/SnO₂/perovskite/phenyl-C61butyric acid methyl ester(PCBM)/Au and FTO/poly(3,4-ethylenedioxythiophene)-poly (styrenesulfonate)(PEDOT:PSS)/perovskite/ poly(triaryl amine)(PTAA)/Au, and characterized the evolution of the space-charge-limited current for different biases. The sharp rise in the J-V curve (Fig. 3A) correlated with the trapfilled limit. The defect density was calculated according to the following equation:

$$N_{\text{defects}} = 2\varepsilon\varepsilon_0 V_{\text{TFL}} / eL^2 \tag{4}$$

Where ε and ε_0 represent the dielectric constants of FAPbI₃ and the vacuum permittivity, respectively, *L* the thickness of the obtained perovskite film, and *e* the elementary charge. We estimated the corresponding electronic trap densities ($N_{\rm defects}$) to be 5.4 × 10¹⁵, 7.6 × 10¹⁵, 5.7 × 10¹⁵, and 8.0 × 10¹⁵ cm⁻³ for x = 0, 1.9, 3.8, and 5.7, respectively, in FAPbI₃:*x*MDACl₂ and 1.0 × 10¹⁶ cm⁻³ for the control. Thus, the addition of MDACl₂ had little effect on the electron trap densities, although the hole-trap density for the hole-only device decreased with the addition of MDACl₂ relative to the control (fig. S11).

Furthermore, we compared the chargecarrier lifetimes of the target and control deposited on a quartz substrate. Figure 3B shows the time-dependent photoluminescence recorded with a commercial time-correlated single-photon counting instrument. The lifetime for nonradiative recombination for the target was 1562 ns, longer than that of the control (715 ns). The recombination value was obtained with the use of the following biexponential equation: $Y = A_1 \exp(-t/\tau_1) +$ $A_2 \exp(-t/\tau_2)$, where τ_1 and τ_2 denote the fast and slow decay time constants, respectively, and are related to the radiative and trap-assisted nonradiative recombination processes.

The increase in the charge-carrier lifetime indicates the presence of another beneficial factor, and we focused on the coexistence of Cl^{-} in the FAPbI₃ lattice substituted for I^{-} sites, in interstitial sites, or both. A small amount of Cl⁻ substitution can reduce the lattice strain of FAPbI3 and suppress defect formation. Also, MDA²⁺ cations could ease the insertion of smaller Cl- ions into the interstitial spaces. The residual Cl content determined by x-ray photoelectron spectroscopy and time-of-flight secondary-ion mass spectrometry (Fig. 3, C and D) of FAPbI3 with 3.8 mol % MDACl₂ was higher than that of the control in the absence of MDACl₂. The Cl content was higher at the interface with the TiO₂ electrode, which is expected to increase the light stability of PSCs (41, 42). The small amount of Cl observed in the control is consistent with a previously reported result (43); smaller Cl⁻ ions substituted at the I sites of FAPbI₃ contribute to phase stabilization by reducing the lattice strain of the α -phase (44). However, the Cl concentration in FAPbI₃ remained nearly unchanged even when excess Cl[−] ions were added, so we attributed the increase in residual Cl observed in Fig. 3, C and D, to MDA²⁺-related interstitial sites rather than to substitution into I sites. Interstitial Cl⁻ ions would account for the slight increase in bandgap (Fig. 1B) and a shift in the XRD peak to a lower angle (Fig. 3E) from an expanded unit cell versus contraction from the increased number of hydrogen bonds (I--H-N) between PbI_2 and MDA and the increased V_{FA} to compensate for charge neutrality.

We compared the long-term humidity-, thermal-, and photostability performances of the unencapsulated control and target in

Fig. 4. We used copper phthalocyanine (CuPC) as the hole-transporting material (HTM) to prevent degradation by hygroscopic dopants and spiro-OMeTAD itself at 150°C. The initial device parameters and J-V curves are presented in table S2 and fig. S12, respectively. The target device exhibited higher humidity stability, retaining >90% of the initial PCE after 70 hours under high humidity [85% relative humidity (RH), 25°C], than the control PCE, which reduced to 40% of the initial value (Fig. 4A). The thermal stability monitored at 150°C and ~25% RH (Fig. 4B) indicates that the control device PCE degraded gradually, reaching <20% of the initial PCE after 17 hours due to MA evaporation. The target device maintained >90% of its initial PCE and exhibited greatly improved thermal stability. In addition, the long-term photostability (encapsulation and ambient condition) of the PSC, including spiro-OMeTAD as HTM, was tested with maximum power point tracking under full solar illumination without a UV filter (Fig. 4C). Despite the use of a TiO_2 photoelectrode with high photocatalytic effect, the target device exhibits very high photostability, maintaining ~90% of its initial PCE (>23.0%) over 600 hours of irradiation. This result can be attributed to both the high concentration of Cl ions in the interface between photoelectrode and perovskite (42) and the stabilization of the α -phase on FAPbI₃ by MDACl₂.

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SUPPLEMENTARY MATERIALS

science.sciencemag.org/content/366/6466/749/suppl/DC1 Materials and Methods Figs. S1 to S12 Tables S1 to S2 References (45–50)

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Efficient, stable solar cells by using inherent bandgap of α -phase formamidinium lead iodide

Hanul Min, Maengsuk Kim, Seung-Un Lee, Hyeonwoo Kim, Gwisu Kim, Keunsu Choi, Jun Hee Lee and Sang II Seok

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Maintaining the bandgap The bandgap of the black α -phase of formamidinium-based lead triiodide (FAPbI₃) is near optimal for creating high-efficiency perovskite solar cells. However, this phase is unstable, and the additives normally used to stabilize this phase at ambient temperature—such as methylamonium, caesium, and bromine—widen its bandgap. Min *et al.* show that doping of the α -FAPbI₃ phase with methylenediammonium dichloride enabled power conversion efficiencies of 23.7%, which were maintained after 600 hours of operation. Unencapsulated devices had high thermal stability and retained >90% efficiency even after annealing for 20 hours at 150°C in air. Science, this issue p. 749

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