# Facet-dependent active sites of a single Cu<sub>2</sub>O particle photocatalyst for CO<sub>2</sub> reduction to methanol

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Atomic-level understanding of the active sites and transformation mechanisms under realistic working conditions is a prerequisite for rational design of high-performance photocatalysts. Here, by using correlated scanning fluorescence X-ray microscopy and environmental transmission electron microscopy at atmospheric pressure, in operando, we directly observe that the (110) facet of a single Cu<sub>2</sub>O photocatalyst particle is photocatalytically active for CO<sub>2</sub> reduction to methanol while the (100) facet is inert. The oxidation state of the active sites changes from Cu(1) towards Cu(1) due to CO<sub>2</sub> and H<sub>2</sub>O co-adsorption and changes back to Cu(1) after CO<sub>2</sub> conversion under visible light illumination. The Cu<sub>2</sub>O photocatalyst oxidizes water as it reduces CO<sub>2</sub>. Concomitantly, the crystal lattice expands due to CO<sub>2</sub> adsorption then reverts after CO<sub>2</sub> conversion. The internal quantum yield for unassisted wireless photocatalytic reduction of CO<sub>2</sub> to methanol using Cu<sub>2</sub>O crystals is ~72%.

fficient catalytic reduction of  $CO_2$  into liquid fuels can play an integral role in potentially mitigating and resolving the challenge of climate change. Previous studies have focused on  $CO_2$ reduction by hydrogenation at elevated temperatures<sup>1-4</sup> although use of hydrogen raises cost and safety issues. Recently, electroreduction of  $CO_2$  has aroused considerable interest<sup>5-10</sup>. However, this process requires a large overpotential<sup>8,11</sup> and may cause decomposition of generated chemicals<sup>12</sup>. Additionally, use of electricity as a secondary form of energy is inefficient due to losses associated with conversion from primary sources as well as storage and transmission.

Photoreduction of  $CO_2$  using sunlight offers an efficient way to directly store solar energy in the form of chemical fuels. Of particular interest is its conversion directly into liquid fuels such as methanol<sup>13</sup>. In theory, combined photovoltaic-electrolysis approaches have lower 'solar-to-fuel' efficiency than photocatalysis due to the energy losses associated with additional conversion steps from solar to electricity and the subsequent conversion to fuels. Photocatalysis is a promising method of converting solar energy to fuel without secondary energy losses and it can potentially be scaled up for largescale applications<sup>14</sup>. Unfortunately, despite intensive efforts since 1978 (ref. <sup>15</sup>), the photocatalysts developed for  $CO_2$  reduction suffer from poor stability<sup>16,17</sup>, low efficiency in the visible part of the solar spectrum<sup>18</sup>, high cost as a result of the use of low-abundance elements<sup>19</sup> and low multi-electron transfer necessary for chemical fuel production<sup>20</sup>.

Cu compounds are promising photocatalysts with good multielectron transfer properties<sup>21</sup> because of their loosely bonded delectrons. Although in their relative infancy<sup>22</sup>, such catalysts show great potential for facilitating CO<sub>2</sub> activation and conversion to CO (refs. <sup>7,8,23,24</sup>) or ethylene<sup>25–27</sup> by electroreduction, as well as for water splitting<sup>28,29</sup>, solar cells<sup>30</sup> and Li ion batteries<sup>31</sup>. Furthermore, Cu<sub>2</sub>O is an inexpensive material based on relatively abundant elements. Cu<sub>2</sub>O is a direct-gap semiconductor with a band gap of 2.1 eV enabling the absorption of visible light that constitutes most of the solar spectrum<sup>28</sup>. Importantly, it shows intrinsic p-type conductivity due to the presence of negatively charged Cu vacancies and possesses one of the lowest electron affinities of the elements. This makes Cu<sub>2</sub>O a good candidate for the reduction of one of the most resilient compounds, namely CO<sub>2</sub> (refs. <sup>32-35</sup>). Previous studies have focused on traditional statistical ex situ ensemble bulk measurements. Solar-to-fuel efficiencies reported for unassisted wireless photocatalytic CO<sub>2</sub> reduction vary across laboratories; however, in general, they are substantially smaller compared to electrical to chemical fuel efficiency in photoelectrochemical cell devices<sup>36,37</sup>. To design highly efficient one-step CO<sub>2</sub> reduction using unassisted wireless photocatalysis in neutral pH, we need to design a photocatalyst with selective active sites. However, identification of the selective active sites by traditional ex situ ensemble methods is difficult or even impossible because of the inhomogeneity of the active sites and their sensitivity to their local environment.

Here, we report the facet specific active sites in  $Cu_2O$  capable of binding and reducing  $CO_2$ . This is unravelled using operando multimodal correlated characterization of a single particle, ex situ ensemble measurements and density functional theory (DFT) calculations. Through this approach, we exploit our understanding of the facet selective active sites to design nanoparticles with improved activity. We report a photocatalytic conversion of  $CO_2$  with liquid fuel methanol as the sole product of photoreduction where the internal quantum yield is ~72% and the solar-to-fuel efficiency is 10% with the major oxidation products being  $O_2$  and  $H_2O_2$ .

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**Fig. 1** | **Structure**, **characterization and photocatalytic performance of the Cu<sub>2</sub>O photocatalyst particles. a**, Atomic model of Cu terminated by the (100) surface of Cu<sub>2</sub>O used for DFT calculations and a TEM image of a single cubic Cu<sub>2</sub>O nanocrystal with well-defined facets. **b**, NBD from the edge of the nanocrystal in **a**, indicated by the blue arrow. **c**, Atomic model of Cu terminated by the (110) surface of Cu<sub>2</sub>O used for DFT calculations and a TEM image of a cubic Cu<sub>2</sub>O nanocrystal with well-defined facets. **b**, NBD from the edge of the nanocrystal in **a**, indicated by the blue arrow. **c**, Atomic model of Cu terminated by the (110) surface of Cu<sub>2</sub>O used for DFT calculations and a TEM image of another particle with a hexagonal projected shape and well-defined facets. **d**, NBD on the edge of **c**, indicated by the blue arrow. **e**, SEM image of the cubic/octahedral sample measured by GC during the photocatalytic reaction as a function of retention time (black lines), the standard methanol measured by GC (red lines) and the control experiment with CO<sub>2</sub> and H<sub>2</sub>O but without Cu<sub>2</sub>O by GC (blue lines). The dot-dashed line indicates the same retention times of the black-line and red-line peaks. a.u., arbitary units. **g**, Illumination time-dependent conversion of methanol produced by the photocatalytic reaction using CO<sub>2</sub> and liquid H<sub>2</sub>O (red squares) and using CO<sub>2</sub> and H<sub>2</sub>O vapour (blue triangles). The error bars are 10% of the average value of five measurements. Inset: mass spectra of methanol obtained by GC-MS analysis of the photocatalytic reaction products. **h**, O<sub>2</sub> evolution during CO<sub>2</sub> reduction photocatalysis.

#### Characterization of photocatalytic materials

We prepared Cu<sub>2</sub>O nanocrystal samples with well-defined (110) and (100) facets via colloidal synthesis (see Methods and Supplementary Fig. 1). Two types of Cu<sub>2</sub>O samples were synthesized (Fig. 1 and Supplementary Fig. 2): one containing a mixture of cubic and octahedral nanocrystals and the other primarily composed of cubic nanocrystals. Characterization of the first sample, which consisted of cubic/octahedral particles, by transmission electron microscopy (TEM) showed various projections of Cu<sub>2</sub>O crystals with welldefined facets and noticeable truncations at the corners (Fig. 1). Nanobeam electron diffraction (NBD) showed (110) and (200) facets with a [002] zone axis (Fig. 1b). This indicates that the particle shown in Fig. 1a is a cube enclosed by (100) facets with (110) facets at the truncated corners. Particles with a hexagonal projected shape were also observed (Fig. 1c) with a  $[1\overline{1}\overline{1}]$  zone axis, (110) edges and (211) corners seen by NBD (Fig. 1d). All Cu<sub>2</sub>O crystals were cubic phase with space group  $Pn\bar{3}m$ , which were confirmed both by TEM and X-ray powder diffraction measurements. Scanning electron microscopy (SEM) showed that the truncated cubic/octahedral crystals were 500 nm to 1 µm in size (Fig. 1e).

The illumination of faceted Cu<sub>2</sub>O nanoparticles in the presence of CO<sub>2</sub> and water showed the formation of a new molecular product. The products of this photocatalytic reaction (Fig. 1f, black line, see Methods) showed the same retention time as the methanol analytical standard (Fig. 1f, red line) by gas chromatography (GC) analysis, while no products (Fig. 1f, blue line) were observed in the absence of the purified Cu<sub>2</sub>O photocatalyst. The only product was methanol, as confirmed by gas chromatography mass spectrometry (GC–MS) that showed peaks with mass-to-charge ratios of 32, 31, 30, 29, 28 and 15 (Fig. 1g, inset), corresponding to [CH<sub>3</sub>OH•]+, [CH<sub>3</sub>O]+, [CH<sub>2</sub>O]+, [CHO]+, [CO]+ and [CH<sub>3</sub>]+. The same species were obtained in a standard methanol mass spectrum (Supplementary Fig. 3).

The amount of methanol produced was determined by comparing the integrated area of the GC peaks with that of known amounts of standard pure methanol. The methanol production and its dependence on irradiation time (Fig. 1g) showed an internal quantum yield of 72% (see Methods and Supplementary Fig. 4), which is comparable to the recent reported efficiencies of 50% and 63% obtained using electroreduction<sup>10,38</sup>. The solar-to-fuel efficiency was measured to be 10% (see Methods), which was smaller than the theoretically predicted maximum of 35.2% (ref. <sup>39</sup>). The plateau in the time conversion curve (Fig. 1g) indicated that the methanol produced by photoexcited electrons reached a steady state with the same formation and consumption rate beyond 6 min of irradiation time.

Concomitantly, the system produced oxygen in synchrony with periodic illumination (Fig. 1h and see Methods), suggesting that our photocatalysts can couple H<sub>2</sub>O oxidation with CO<sub>2</sub> reduction, functioning as an artificial leaf<sup>9</sup>, albeit in a substoichiometric ratio (see Methods and Supplementary Fig. 5). Substoichiometric oxygen production was previously observed; its formation rate is typically several orders of magnitude smaller than those of the main catalytic products<sup>40</sup>. Some of the photogenerated holes may be consumed in the photo-induced disproportionation of Cu<sub>2</sub>O nanoparticles, the formation of H<sub>2</sub>O<sub>2</sub> (see Methods) and other reactive oxygen species (ROS) and/or in re-oxidation of methanol, contributing to the steady state in the methanol evolution. The oxygen isotope tracer experiment using oxygen-labelled water (H<sub>2</sub><sup>18</sup>O; for details, see Methods and Supplementary Fig. 6) showed that the majority of photogenerated molecular oxygen is isotopically mixed <sup>16</sup>O<sup>18</sup>O with small amounts of  ${}^{18}O_2$  and  ${}^{16}O_2$ . This indicates that H<sub>2</sub>O is the major source for the oxygen evolution reaction (although the initial water is pure H<sub>2</sub><sup>18</sup>O, one molecule of water H<sub>2</sub><sup>16</sup>O is produced from each C<sup>16</sup>O<sub>2</sub> converted into methanol to contribute to <sup>16</sup>O content), with a possible smaller contribution from Cu216O isotopic exchange. Lattice oxygen in the highly active oxide was shown to contribute to the generation of oxygen due to the hybridization of the oxygen orbital with the Fermi level of the photocatalysts<sup>41,42</sup>. The total amount of oxygen produced in the reaction after 2 min of illumination was 1.45 mmol (~6% efficiency). The formation of  $H_2O_2$  was also determined to be 5.8 mmol using potassium permanganate titration (~12% efficiency). Singlet oxygen, adsorbed oxoand peroxo- species may also be produced but their amount was hard to quantify.

Noticeably, the same amount of methanol was produced if H<sub>2</sub>O vapour was used instead of liquid H<sub>2</sub>O (Fig. 1g and Supplementary Fig. 7), which indicates that H<sub>2</sub>O vapour oxidation follows a similar reaction mechanism to that of liquid H<sub>2</sub>O. The <sup>13</sup>CO<sub>2</sub> isotope tracer experiment (see Methods and Supplementary Fig. 8) showed the formation of <sup>13</sup>CH<sub>3</sub>OH as the only product, with GC-MS analysis showing distinct mass-to-charge ratios of 16 and 33 corresponding to  $[{}^{13}CH_3]^+$  and  $[{}^{13}CH_3OH\bullet]^+$  when compared with the product from the unlabelled <sup>12</sup>CO<sub>2</sub> reaction. The remaining peaks at 45 and 44 correspond to the unreacted <sup>13</sup>CO<sub>2</sub> and <sup>12</sup>CO<sub>2</sub> in the sealed vial, respectively (Supplementary Fig. 8). This gives strong evidence that the carbon in methanol originates from CO<sub>2</sub>. We did not observe methanol oxidation and formation of H<sub>2</sub> as a by-product from H<sub>2</sub>O reduction. This is probably due to the fast desorption of methanol from the surface of the catalyst. Methanol oxidation is relatively easy to observe using electron paramagnetic resonance (EPR) due to the doublet lines at the 130 G separation of the methanol radical and the formyl radicals<sup>43</sup>. However, we did not observe any similar radical during illumination of Cu<sub>2</sub>O particles with CO<sub>2</sub> and H<sub>2</sub>O in the EPR cavity at 4K, as shown in the section 'Ensemble characterization at a statistic level'. Nevertheless, illumination of the same amount of cubic Cu<sub>2</sub>O sample (Supplementary Fig. 2b) in the presence of CO<sub>2</sub> and H<sub>2</sub>O produced only a fraction (~10%) (Supplementary Fig. 2c,d) of the methanol observed with the cubic/octahedral sample (Supplementary Fig. 2a), suggesting that the (100) facets were nearly inactive whereas the sample with higher-index facets was more active photocatalytically.

### Operando multimodal characterization of single particles

To understand the role of the active sites and local crystal structure on the photoreduction of  $CO_2$  under real working conditions, we studied single Cu<sub>2</sub>O particles by both environmental transmission electron microscopy (ETEM) and scanning fluorescence X-ray microscopy (SFXM) in a nanoreactor<sup>44,45</sup> while exposed to a mixture of gaseous CO<sub>2</sub>, H<sub>2</sub>O and light (see Methods, Fig. 2a and Supplementary Fig. 9). For example, we first identified particle I (Fig. 2d), with a truncated hexagonal shape (red circle in Fig. 2c), by TEM. A corresponding single-particle electron diffraction (SPED) measurement indicated that the edges of particle I were {220} planes in the TEM view (Supplementary Fig. 10). Then, we transferred the nanoreactor containing the Cu<sub>2</sub>O particles from the ETEM to the SFXM platform.

Low-resolution scans (Fig. 2e) using the Cu and Ni Ka fluorescence signals from the Cu<sub>2</sub>O particles and grid bars, respectively, enabled us to locate the same particle I (white circle in Fig. 2e) for imaging and nanospectroscopic study by SFXM. In the high-resolution Cu Ka fluorescence image in Fig. 2f we observe particle I as a truncated square rather than a truncated hexagon observed in TEM. This was due to the relative projected views of the particle using the ETEM and SFXM platforms, which differed by nearly 60°, and can be understood by considering particle I to be a cube resting on one of its truncated corners, a (111) facet (Fig. 2b). With the TEM view directions as indicated by the red arrow, the particle appeared as a corner-truncated hexagon (Fig. 2b,d), while it presented as a corner-truncated square when viewed with X-ray imaging (Fig. 2b,f). Hence, the corner of particle I was a (110) facet while edges are exposed as a (100) facet along the X-ray view direction (Fig. 2b,f). We targeted specific facets of particle I by directing the nanofocused X-ray beam parallel to desired facets whilst scanning the incident X-ray energy (Supplementary Fig. 9), enabling us to obtain facet-dependent spectroscopic information of the Cu active sites. The spectroscopic signal, which resulted from a columnar projection of the X-ray beam through the plane of the facet, probably reflected an admixture of states resulting from both the surface and interior chemistry.

X-ray fluorescence spectra (Fig. 3a) taken near the Cu K edge on the (110) facet of Cu<sub>2</sub>O particle I in its initial pristine state revealed a white-line peak at 8,981.0 eV, indicating a Cu(1) oxidation state associated with this facet, in agreement with reference X-ray absorption spectra taken from pure Cu<sub>2</sub>O (Supplementary Fig. 11). By contrast, the white-line peak observed on the (100) facet was at 8,981.5 eV, indicating the presence of both Cu(1) and Cu(11) oxidation states (Fig. 3a). When measuring spectra on the (110) facet before and after flushing CO<sub>2</sub>/H<sub>2</sub>O through the nanoreactor for a few minutes, the peak shifted of between 1.0 and 1.5 eV towards higher energy (Fig. 3b and Supplementary Fig. 12), indicating that the oxidation state of the Cu changed from oxidation state 1<sup>+</sup> towards 2<sup>+</sup> due to the co-adsorption of CO<sub>2</sub>/H<sub>2</sub>O. In addition, consecutive spectra measured at intervals of 10 min at another position on the (110) facet of particle I showed a gradual shift of the peak by ~1 eV until saturation was reached in about 1 h (Supplementary Fig. 12). These results suggest that  $CO_2/H_2O$  co-adsorbed on the (110) facet reduced the electron density of the exposed Cu atoms, causing the oxidation state to shift from Cu(I) towards Cu(II). Furthermore, when the particle was excited with a green laser (532 nm) while maintaining a constant CO<sub>2</sub>/H<sub>2</sub>O gas flow, the peak reverted back by ~1 eV towards lower energy (Fig. 3b), indicating that the majority of the Cu oxidation states shifted back to Cu(I). In comparison, the oxidation states of the Cu atoms on the (100) facet did not change on  $CO_2/H_2O$  adsorption in either the dark or the illuminated states (Fig. 3c), indicating the (100) facet was photocatalytically inactive. In successive measurements on several other, similar Cu<sub>2</sub>O particles, we reproducibly observed similar spectroscopic shifts on the (110) facets followed by reversion on illumination, and statistically insignificant shifts on the (100) facets. We note that the X-ray measurements could not have triggered the photocatalytic reaction because the chemical state



**Fig. 2 | Operando multimodal imaging of a single photocatalyst Cu<sub>2</sub>O particle. a**, Schematic illustration of the setup showing the gas-flow nanoreactor. **b**, Schematic of the electron beam and X-ray directions for TEM and SFXM imaging on particle I, a truncated cube. **c**, TEM overview of Cu<sub>2</sub>O particles inside the nanoreactor. The red circle indicates particle I. **d**, TEM image showing particle I with the electron beam parallel to the [111] direction of the Cu<sub>2</sub>O cube. **e**, Low-resolution SFXM image of the particles inside the nanoreactor from Cu K $\alpha$  and Ni K $\alpha$  emission with a white circle indicating particle I. The different aspect ratios between **c** and **e** are due to the 60° difference between the X-ray and TEM imaging directions. **f**, High-resolution SFXM image of particle I from Cu K $\alpha$  emission with the incident X-ray beam parallel to the [001] direction. The colour scales in **e** and **f** indicate the intensity of the Cu fluorescence signal. The units in the colour bars are arbitrary.

of the Cu in the pristine  $Cu_2O$  did not change over the course of many measurements without the introduction of gas, and it did not change from Cu(II) to Cu(I) until subsequent optical illumination (Supplementary Fig. 12).

Particle I was checked again by TEM and SPED after the X-ray measurements, which showed that the location and structure of particle I were preserved (Supplementary Fig. 13). This finding rules out the possibility that the particle moved or was significantly damaged structurally by the X-ray beam. While we characterized the same Cu<sub>2</sub>O particles by both ETEM and SFXM, it should be noted that in TEM the particle could not be tilted to the same view orientation as for the SFXM measurements. To verify these results we used ETEM to examine another Cu<sub>2</sub>O particle in the nanoreactor, the truncated cube particle II, with the same orientation as particle I that we studied with X-ray imaging (we could not tilt particle I due to geometric limitations). SPED measurements confirmed the presence of {110} facets on the truncated corners and {100} facets on the faces of particle II (Supplementary Fig. 10), as with particle I in the X-ray measurements. Electron energy loss spectra (EELS) recorded on the (110) facet before and after CO<sub>2</sub>/H<sub>2</sub>O co-adsorption in the dark state showed that the Cu  $L_{2,3}$  edge (2p to 3d transition) shifted ~2 eV towards lower energy on  $CO_2/H_2O$  adsorption (Fig. 3d). The EELS spectrum from the pristine particle (Fig. 3d, black line) is consistent with the peak position and  $L_3/L_2$  peak intensity ratio of  $Cu_2O$  (ref. <sup>46</sup>). The broadening of the spectrum recorded with gas flow (Fig. 3d, red line) and the decrease in the  $L_3/L_2$  peak intensity ratio indicate the formation of Cu<sup>0</sup>. The energy shift between the spectral peaks of the particles taken before and after gas adsorption (Fig. 3d, black and red lines) shows that, of the superposition of copper phases present, Cu(II) dominates because the  $L_3$  peak energies of Cu<sup>0</sup> and Cu(I) are the same. This is further quantified by EPR with much higher sensitivity described in section 'Ensemble characterisation at a statistic level'. Similar energy shifts were observed by EELS on the particle III and particle IV (110) facets (Supplementary Fig. 12). These shifts are consistent with the oxidation state change from Cu(I) to Cu(II) seen with X-ray imaging<sup>46</sup>. Again, consistent with the X-ray measurements, the EELS data did not show a change in the Cu oxidation state on the (100) facet.

We also carried out DFT calculations to understand the mechanism of CO<sub>2</sub> and H<sub>2</sub>O adsorption on Cu<sub>2</sub>O on different facets (Fig. 4a,b). The calculations showed that the co-adsorption of CO<sub>2</sub> and H<sub>2</sub>O are thermodynamically favourable on both the (110) and (100) facets (Cu termination was considered for both facets), although the co-adsorption on (110) is much more favourable (energy of adsorption,  $E_{ad} \approx -2.8 \text{ eV}$ ) than that on (100) ( $E_{ad} \approx -1.87 \text{ eV}$ ). It is notable that the co-adsorption of CO<sub>2</sub> and H<sub>2</sub>O



**Fig. 3 | Operando multimodal nanospectroscopy of a single photocatalyst Cu<sub>2</sub>O particle under various working conditions. a**, Scanning X-ray nanospectroscopy of various facets of particle I (Fig. 2f) in the pristine state.  $\Delta$  denotes the shift in the white line peak in the Cu K-edge X-ray spectrum. Inset: the locations for spectra recorded from particle I, with the corresponding colours indicated. **b**, Oxidation state changes at the (110) facet of particle I in the pristine state, with CO<sub>2</sub>/H<sub>2</sub>O gas exposure under dark and under illuminated conditions. **c**, Oxidation states at the (100) facet of particle I under the same conditions. **d**, In situ EELS on the (110) facet of particle II. Inset: the TEM images of particle II. The lower signal-to-noise ratio of the spectrum taken with gas is due to greater electron scattering in comparison to in vacuum. a.u., arbitrary units.

on Cu-terminated (110) led to a dramatic change of the arrangement of Cu atoms on the surface leading to the formation of Cu clusters (Fig. 4a,c,d). This was confirmed by the EPR measurements. This cluster formation, however, did not occur on either the Cu/Oterminated (the termination have both Cu and O atoms) (110) or the Cu-terminated (100) on CO2 and H2O adsorption, based on our calculations (Fig. 4b,e,f). This suggests that the (110) facets of the synthesized Cu<sub>2</sub>O particles are likely to be Cu terminated and therefore have greater photocatalytic activity for co-adsorption of CO<sub>2</sub> and H<sub>2</sub>O. This also indicates that the surface termination of Cu<sub>2</sub>O has a strong effect on the photocatalytic reaction activity and may affect the final product from the reaction. We carried out further calculations and found that, among all the Cu<sub>2</sub>O surfaces we considered, only CO<sub>2</sub> co-adsorption with  $H_2O$  on the Cu-terminated Cu<sub>2</sub>O (110) leads to bent  $CO_2$  ( $CO_2^{-}$ ) (Supplementary Fig. 14), and the conversion from a physisorbed  $CO_2$  to a bent  $CO_2$  on this surface is energetically favourable ( $\Delta E = -6.3 \text{ eV}$ ). Such a highly exergonic reaction indicates that the first energy barrier for the reduction could be low. For the Cu/O-terminated Cu<sub>2</sub>O (110) and all the Cu<sub>2</sub>O (100) terminations, the geometry optimizations show that the bent CO<sub>2</sub> is not stable on these surfaces. Our calculations also show that, without H<sub>2</sub>O co-adsorption, the bent CO<sub>2</sub> is not stable on any of the considered surfaces.

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#### Ensemble characterization at a statistical level

We characterized the properties of the ensembles of Cu<sub>2</sub>O particles by high-resolution X-ray powder diffraction (HRXRD) and EPR to obtain complementary information about the photocatalytically active sites. Operando high-energy X-ray diffraction measurements (Supplementary Fig. 15) show the absence of noticeable structural changes (no line broadening, disappearance/appearance of additional lines or decrease of line intensity) in the Cu<sub>2</sub>O particles over 8h during the photocatalytic reaction. The stability is also confirmed by the ex situ X-ray photoemission spectroscopy. Formation of strong Cu 2p satellite peaks, characteristic of oxidized CuO samples, were not observed in both pristine and reacted Cu2O catalyst (Supplementary Fig. 16). Fine structural changes in the Cu<sub>2</sub>O particles were probed ex situ using HRXRD (Fig. 5a). In the HRXRD data in Fig. 5b,c, we see that CO<sub>2</sub>/H<sub>2</sub>O co-adsorption under dark conditions increased the crystal lattice constant, which is consistent with our spectroscopic observation (Fig. 3b) that the CO<sub>2</sub> withdrew the electron density from the Cu active sites, leading to an oxidation state change from Cu(1) to Cu(11). Consistently, using DFT calculations we found substantial charge transfer at the co-adsorption sites of CO<sub>2</sub> and H<sub>2</sub>O at the Cu-binding site of the (110) surface (Fig. 5f). Relative to bulk crystalline  $Cu_2O$  with Cu–O bond lengths of 1.86 Å, the Cu-O bonds in the vicinity of the CO<sub>2</sub> and H<sub>2</sub>O co-adsorption



**Fig. 4 | DFT calculations of binding energies. a**, Confirmation of  $CO_2/H_2O$  co-adsorption on  $Cu_2O$  (110) Cu termination. The binding energy is -2.8 eV. **b**, Confirmation of  $CO_2/H_2O$  co-adsorption on  $Cu_2O$  (100) Cu termination. The binding energy is -1.87 eV. **c**, Individual binding confirmation of  $H_2O$  on  $Cu_2O$  (110) Cu termination. The binding energy is -1.04 eV. **d**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (110) Cu termination. The binding energy is -0.04 eV. **d**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (110) Cu termination. The binding energy is -0.21 eV. **f**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (110) Cu/O termination. The binding energy is -0.10 eV. **g**, Individual binding confirmation of  $H_2O$  on  $Cu_2O$  (100) Cu termination. The binding energy is -0.75 eV. **h**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (100) Cu termination. The binding energy is -0.75 eV. **h**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (100) Cu termination. The binding energy is -0.75 eV. **h**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (100) Cu termination. The binding energy is -0.75 eV. **h**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (100) Cu termination. The binding energy is -0.75 eV. **h**, Individual binding confirmation of  $CO_2$  on  $Cu_2O$  (100) Cu termination. The binding energy is -0.79 eV. The Cu, O, C and H atoms are shown in blue, red, brown and white, respectively.

sites were found to be about 1.92-1.95 Å by DFT calculations, which is very similar to the DFT computation of 1.95 Å in bulk crystalline CuO. After optical illumination (Fig. 5b,c, blue lines), the lattice constant of pristine Cu<sub>2</sub>O was recovered (Fig. 5b,c, black lines).

The above results suggest that, on co-adsorption of  $CO_2/H_2O$ on  $Cu_2O$  particles, intermolecular electron transfer occurs between the Cu active sites at the (110) surfaces and  $CO_2$  molecules. This redistribution of electron density results in partial reduction of adsorbed  $CO_2/H_2O$  molecules, lowering the energy barrier for their complete reduction. On the other hand, the electron redistribution reduces the electron density of the (110) surface. This in turn facilitates localization of photogenerated electrons on these sites, enhancing the probability of their interaction with adsorbed molecules. Illumination of  $CO_2/H_2O$  co-adsorbed on the (110) facet, therefore, leads to efficient electron transfer and removal of the bonding between  $CO_2$  molecules and active sites, restoring the electron density of the pristine  $Cu_2O$ . Co-adsorption of  $CO_2/H_2O$  by the particles is consistent with the lattice constant changes seen by HRXRD (Fig. 5b,c and Supplementary Fig. 17), which suggest volumetric expansion and possible  $CO_2$  induced lattice strain. Similar behaviour has been observed in Au nanoparticles undergoing CO or thiol gas adsorption<sup>47,48</sup>.

The complementary role of the adsorbed water to that of  $CO_2$  reduction gives strong corroboration of the photocatalysis process. Besides providing protons for the photocatalytic reaction, HRXRD studies of water adsorption on  $Cu_2O$  particles indicate that the  $H_2O$  molecule donates electrons and pushes the electron density towards the  $Cu_2O$  surface (Supplementary Fig. 18). However, the reverse Cu oxidation state change on co-adsorption of  $CO_2$  and  $H_2O$  (Fig. 3b) indicates the dominating effect of  $CO_2$  over  $H_2O$ .

To investigate binding and electron transfer between Cu active sites and CO<sub>2</sub> molecules, we made EPR measurements on pristine Cu<sub>2</sub>O particles, exposed to CO<sub>2</sub>/H<sub>2</sub>O under dark and illuminated conditions (Fig. 5d). The pristine particles (Fig. 5d, black lines) showed the existence of Cu(II) defects with unpaired electrons in 3d orbitals consistent with the p-doping of Cu<sub>2</sub>O nanoparticles with g tensor characteristic of magnetic moment and anisotropy of copper ions. The spectral shape suggests that the loosely packed pristine Cu<sub>2</sub>O powder is aligned in the magnetic field<sup>49</sup>. The hyperfine pattern of typical Cu complexes disappears in favour of a new aligned single-crystal signal at intermediate g with large and almost unresolved hyperfine splitting<sup>50</sup>. On CO<sub>2</sub> adsorption in the dark, however, the anisotropy of the Cu<sub>2</sub>O powder spectrum and a hyperfine splitting (A) pattern reappears (Fig. 5d, red lines). The full anisotropic spectrum of saturated coordinated Cu(II) is, however, restored only by CO<sub>2</sub>/H<sub>2</sub>O co-adsorption (Fig. 5d, blue lines). In this case, the restored g and A tensors ( $A_{\parallel}$  for Cu of 137 G and  $g_{\parallel}$  of 2.3420) provide evidence for a distorted tetrahedral complex rather than octahedral or square planar coordination, which is typically found for Cu complexes<sup>51</sup>. This indicates the proximity effect of H<sub>2</sub>O and CO<sub>2</sub> co-adsorbed on the same Cu(II) defect and the structural changes of surface defects (Cu(II), unpaired spin) under adsorption. The overall intensity of the spectrum was increased, suggesting the number of spins, and consequently the number of Cu(II) ions, increased on CO<sub>2</sub> and H<sub>2</sub>O co-adsorption in the dark. This suggests that the CO<sub>2</sub> molecules are withdrawing electron density from Cu(I) active sites, in agreement with our X-ray spectroscopic and structural observations (Figs. 3b and 5a-c). The CO<sub>2</sub> and H<sub>2</sub>O co-adsorption increases spin density (Cu(II)) isotopically (Fig. 5d, blue lines), compared to CO<sub>2</sub> adsorption alone (Fig. 5d, red lines) in the dark. This indicates H<sub>2</sub>O co-adsorption can facilitate CO<sub>2</sub> adsorption.

A new weak peak with g = 2.0023 was observed with illumination of Cu<sub>2</sub>O at 4 K (Fig. 5e and Supplementary Fig. 19). This signal originates from the conduction electrons in the metallic Cu clusters formed on illumination<sup>52</sup>. Formation of Cu clusters may play an important role in electron-hole pair separation and catalytic conversion of CO<sub>2</sub> and H<sub>2</sub>O to methanol by lowering the energy barrier and providing a reaction pathway for methanol formation<sup>1,53</sup> as single-atom catalysts. Small Cu<sup>0</sup> clusters can also have strong adsorption strength to the adsorbate species due to the undercoordinated bonds. The Cu clusters can also act as plasmonic centres to enhance the electrical field around the adsorbates while having minor thermal effects on the particle and its surroundings, due to weak electron photon coupling and smaller absorption cross-section compared to nanoparticles. The synergetic interplay between Cu(I) and Cu clusters may be the reason for this high selectivity and high efficiency of methanol formation.

The formation of Cu clusters may be a consequence of the Cu<sub>2</sub>O disproportionation reaction. While the X-ray nanospectroscopy and



**Fig. 5 | Exsitu ensemble-averaged measurement of Cu<sub>2</sub>O particles and charge transfer. a**, HRXRD measurements, showing that Cu<sub>2</sub>O crystals were cubic phase with space group  $Pn\bar{3}m$ . **b**, Close-up view of the (220) peak, showing a noticeable shift. **c**, Close-up view of the (200) peak. **d**, Dark EPR spectra (thicker solid lines) of Cu<sub>2</sub>O particles, showing adsorption induced electron transfer between Cu<sub>2</sub>O particles, CO<sub>2</sub> and H<sub>2</sub>O. Light-induced EPR spectra (thinner solid lines), show redistribution of the electron transfer between the Cu<sub>2</sub>O particles, CO<sub>2</sub> and H<sub>2</sub>O. Light-induced EPR spectra (thinner solid lines), show redistribution of the electron transfer between the Cu<sub>2</sub>O particles, CO<sub>2</sub> and H<sub>2</sub>O with simulated sunlight illumination. The g-factor represents the characteristic magnetic moment and anisotropy of the investigated paramagnetic species. **e**, Formation of Cu clusters under illumination. *g* = 2.0023 is a g-factor of free electrons and is characteristic for collective plasmon orbitals of small clusters. **f**, DFT calculation showing that a substantial charge transfer takes place between the adsorbed CO<sub>2</sub> and H<sub>2</sub>O on Cu-binding sites at the (110) facet that leads to the change of Cu-O bonds on the surfaces, as discussed in the text. The black lines represent the periodic boundaries of the simulation cell, the yellow and light blue shaded regions represent the electronic charge transfer of the system. The Cu, O, C and H atoms are shown in blue, red, brown and white, respectively.

EELS measurements probed the dominant formation of Cu(II), they were not sensitive enough to probe the formation of a small quantity of Cu clusters. Herein, the disproportionation reaction was quantified by EPR and found to be seven orders less than the yield of the main photocatalytic reaction and O<sub>2</sub> generation ( $4.98 \times 10^{-12}$  mol, see Methods). Hence, the observed O<sub>2</sub> evolution (Fig. 1h) does not only originate from Cu(I) disproportionation. Consistent with oxidation state (Fig. 3b) and structural measurements (Fig. 5b), the number of Cu(II) into non-paramagnetic Cu(I).

All the above evidence indicates that the adsorption effect is the dominant effect on  $CO_2$  photocatalytic reduction on  $Cu_2O$  facets. The photoreduction of other electron acceptors such as Au ions does not show obvious charge separation efficiency differences between the (110) and (100) facets (see Methods and Supplementary Fig. 20), in contrast to previously reported spatial charge separation efficiency heterogeneity for water splitting<sup>54</sup>. These results suggest that the coating of  $Cu_2O$  particles with other materials, such as graphene oxide<sup>34</sup>, could change the adsorption sites and active sites and

lead to different reaction pathways that would produce different products instead of methanol.

## Conclusions

Our material design strategy, which integrated X-ray and electron microscopy, bulk X-ray diffraction and EPR methods with DFT calculations, enabled us to correlate the electronic structure with the atomic structure, both at the single-particle and ensemble levels. It was found that the electron density of the Cu active sites on the (110) facet shifted from Cu(I) to Cu(II) on co-adsorption of  $CO_2/H_2O$ . Under 532 nm illumination, the Cu oxidation state changed from Cu(II) back to Cu(I), which indicates bond breakage between the  $CO_2$  and Cu active sites. These changes were accompanied by a volumetric response to photocatalytic stimulation.  $CO_2$  was reduced by accepting photogenerated electrons while the H<sub>2</sub>O was oxidized by accepting photogenerated holes. The disproportionation reaction of Cu<sub>2</sub>O might be a part of the photocatalytic process but the yield was seven orders of magnitude lower than the major photocatalytic reaction. The oxidation state of the (100)

facet remained unchanged on  $CO_2$  adsorption under both dark and illuminated conditions, indicating that the (100) facet is photocatalytically inert.

We anticipate that this facet-dependent single-particle approach can be extended to design other catalytic materials displaying sitespecific behaviour, such as two-dimensional (2D) nanoplatelets and thin films of Cu<sub>2</sub>O with large-area (110) surfaces. We further expect that Cu<sub>2</sub>O, with its site-specific photocatalytic reduction of CO<sub>2</sub>, small electron affinity and large reduction potential, and a band structure optimized for charge pair formation and adsorption of solar radiation, will inspire the development of new classes of materials offering a path towards CO<sub>2</sub> utilization as well as efficient conversion of solar energy into chemical liquid fuels.

#### Methods

**Synthesis of Cu<sub>2</sub>O photocatalyst particles.** We synthesized two types of Cu<sub>2</sub>O samples, one containing a mix of cubic and octahedral nanocrystals (Supplementary Fig. 2a) and the other primarily composed of cubic nanocrystals (Supplementary Fig. 2b) using standard colloidal synthesis methods<sup>55</sup>.

To synthesize the cubic/octahedral Cu<sub>2</sub>O sample, deionized water (88.2 ml), Cu(CH<sub>3</sub>COO)<sub>2</sub> aqueous solution (5 ml, 0.1 M) and sodium dodecyl sulfate (0.87 g) were successively added into a three-neck flask. The three-neck flask was placed in a water bath at 60 °C with vigorous magnetic stirring. After complete dissolution of the sodium dodecyl sulfate powder in 1 h, NaOH solution (1.8 ml, 1.0 M) and D-(+)-glucose aqueous solution (5 ml, 0.1 M) were successively quickly injected into the flask. The total volume of the solution was 100 ml. The flask was kept in a water bath at 60 °C after the injection under vigorous magnetic stirring for 1 h. The solution colour gradually changed from blue to green, yellow, orange and finally brick-red (Supplementary Fig. 1). After the reaction, the precipitate was separated from the solution by centrifugation at 5,000 r.p.m. for 5 min and washed several times with deionized-water–ethanol solution.

Synthesis of the cubic Cu<sub>2</sub>O sample was performed by dissolving CuCl<sub>2</sub>•2H<sub>2</sub>O (0.171 g) into aqueous solution (100 ml). NaOH aqueous solution (10 ml, 2.0 M) was then added drop-wise into the above transparent light green solution. After vigorous magnetic stirring for 0.5 h, ascorbic acid solution (10 ml, 0.6 M) was added to the solution. The colour of the liquid changed gradually from black to turbid red. The mixture was aged for 3 h. All procedures were carried out under constant vigorous magnetic stirring and heated at 55 °C. The resulting precipitate was separated from the solution by centrifugation at 5,000 r.p.m. for 5 min and washed several times with deionized water and ethanol.

All the particles were washed and purified to remove ligands after the synthesis. Both samples were dried inside a  $N_2$  gas-flow glove box under room temperature for 48 h and then stored inside the  $N_2$  gas-flow glove box to avoid any contamination from organics for later characterization. All chemicals were purchased from Sigma Aldrich.

Photocatalytic activity characterization. We characterized the photocatalytic activity of both Cu<sub>2</sub>O particle syntheses by measuring their methanol production (Supplementary Fig 4c,d). In a typical reaction, 0.01 g of purified particles were dispersed in 5 ml of ultrapure deionized H<sub>2</sub>O in a 20-ml Agilent Headspace vial and sealed. The suspension was thoroughly degassed to remove air with CO2 gas for 10 min. CO2 gas (Airgas, Inc.) was bubbled through a deionized H2O bubbler glassware for the CO<sub>2</sub>/H<sub>2</sub>O gas mixture. The suspension was continuously flushed with the CO<sub>2</sub>/H<sub>2</sub>O gas mixture for another 30 min to saturate the suspension with CO2. Then the suspension was illuminated using a 300 W Xe lamp (Perkin-Elmer Optoelectronics) at a power of 204 W with continuous CO2/H2O gas flow for 0-60 min. The sealed suspension was removed from the gas line after illumination. The methanol in both the gaseous product and the liquid in aqueous solution was extracted out at a temperature above the boiling point of methanol and then injected for quantification using an Agilent Headspace sampler (model 7697 A) connected to an Agilent GC-MS analyser (model 5975C GC-MS, with quadrupole detector, thermal conductivity detector (TCD) detector, DB-5ms column and helium carrier gas). Standard methanol (LC-MS ultra CHROMASOLV, tested for UHPLC-MS) was purchased from Fluka. The formation rate of methanol is about 1.2 mol g<sup>-1</sup>h<sup>-1</sup>. For quantification of oxygen, CO<sub>2</sub> was saturated in suspension first. Then the suspension was illuminated using a 300 W Xe lamp (Perkin-Elmer Optoelectronics) at a power of 204 W for 2 min in the sealed headspacer without any gas flow. The amount of oxygen in gas phases was extracted for the quantification in GC (Supplementary Fig. 5). The dissolved oxygen in H<sub>2</sub>O was calculated using the solubility of 8.6 mg l<sup>-1</sup>.

In situ oxygen measurement was performed using an oxygen needle sensor from Unisense. A mass of 0.01 g of purified particles was dispersed in 20 ml of ultrapure deionized H<sub>2</sub>O in an Agilent Headspace vial and sealed. The oxygen needle sensor was inserted into the solution through the rubber stopper for the whole measurement. The suspension was thoroughly degassed to remove air with CO<sub>2</sub> gas for 30 min with vigorous magnetic stirring. CO<sub>2</sub> gas (Airgas, Inc.) was bubbled through a deionized  $H_2O$  bubbler glassware to prepare the  $CO_2/H_2O$  gas mixture. The initial drop of  $O_2$  was due to this purge of air out of the bottle. The suspension was illuminated using the Xe lamp at a power of 204 W with continuous  $CO_2/H_2O$  gas flow for the oxygen evolution measurements under illuminated conditions. Then the lamp was turned off for about 30 min for the oxygen evolution measurements under dark conditions. This light on/off cycle was repeated for several cycles.

**Carbon isotope tracer measurements.** The  ${}^{13}CO_2$  ( ${}^{13}C$  99%) was purchased from Cambridge Isotope Laboratories, Inc. Here 0.01 g of purified particles were dispersed in 5 ml of ultrapure deionized H<sub>2</sub>O in a 20-ml Agilent Headspace vial and sealed. Then N<sub>2</sub> was used to purge the air from the vial for 2 min followed by three freeze-pump-thaw cycles to degas the photocatalyst suspension. After degassing, the  ${}^{13}CO_2$  was introduced into the photocatalyst suspension by flushing isotope gas into the vial for 60 s. Then the samples were irradiated using a 300 W Xe lamp (Perkin-Elmer Optoelectronics) at a power of 204 W for 2 min at the same distance as for the photocatalytic activity measurements. The photocatalytic product was measured (Supplementary Fig. 8) using an Agilent Headspace sampler (model 7697 A) connected to an Agilent GC–MS analyser (model 5975C GC–MS, with quadrupole detector, TCD detector, DB-5ms column and helium carrier gas).

**Oxygen isotope tracer measurements.** The  $H_2^{18}O$  (<sup>18</sup>O, 97%) was purchased from Cambridge Isotope Laboratories, Inc and the <sup>18</sup>O<sub>2</sub>/helium mixture (ratio 1:9) from Sigma Aldrich. The <sup>18</sup>O<sub>2</sub>/helium mixture (ratio 1:9) or <sup>16</sup>O<sub>2</sub> was injected for calibration of GC. Then 0.01 g of purified particles were dispersed in 5 ml of  $H_2^{18}O$ in a 20-ml Agilent Headspace vial and sealed. Then N<sub>2</sub> was used to purge the air from the vial for 2 min followed by three freeze–pump–thaw cycles to degas the photocatalyst suspension. After degassing, the <sup>13</sup>Cl<sup>6</sup>O<sub>2</sub> was introduced into the photocatalyst suspension by flushing isotope gas into the vial for 60 s (<sup>13</sup>C was used to change the molecular weight of methanol from 32 to 33 and differentiate it from <sup>16</sup>O<sub>2</sub>). Then the samples were illuminated using a 300 W Xe lamp (Perkin-Elmer Optoelectronics) at a power of 204 W for 2 min at the same distance as for the photocatalytic activity measurements. The oxygen from the photocatalytic reaction was measured (Supplementary Fig. 6) using an Agilent Headspace sampler (model 7697A) connected to an Agilent GC–MS analyser (model 5975C GC–MS, with quadrupole detector, TCD detector, DB-5ms column and helium carrier gas).

**Quantum yield determination.** Actinometry was used to determine the photon flux from the Xe lamp with the same power (204 W) and distance from the optical source to the sample as for the photocatalytic reaction. Since Cu<sub>2</sub>O has a band gap of 2.1 eV (~590 nm), we chose iron(III) oxalate as the standard chemical actinometry following the International Union of Pure and Applied Chemistry (IUPAC)<sup>56</sup>. Iron(III) oxalate hexahydrate (Fe<sub>2</sub>(C<sub>2</sub>O<sub>4</sub>)<sub>3</sub>•6H<sub>2</sub>O), ammonium iron(II) sulfate hexahydrate ((NH<sub>4</sub>)<sub>2</sub>[Fe(SO<sub>4</sub>)<sub>2</sub>]•6H<sub>2</sub>O), 1,10-phenathroline, sodium acetate (CH<sub>4</sub>COONa) and sulfuric acid were purchased from Sigma Aldrich.

Then 1,10-phenathroline was used for photometric determination of Fe<sup>2+</sup> by forming [Fe(phen)<sub>3</sub>]<sup>2+</sup>, which has strong absorption at 510 nm. The molar absorptivity of [Fe(phen)<sub>3</sub>]<sup>2+</sup> was determined by using ammonium iron(II) sulfate and 1,10-phenathroline. A 0.4 mM Fe2+ iron solution was made freshly by dissolving 0.00392 g of (NH<sub>4</sub>)<sub>2</sub>[Fe(SO<sub>4</sub>)<sub>2</sub>]•6H<sub>2</sub>O with 25 ml of 0.05 M H<sub>2</sub>SO<sub>4</sub> solution. Then 0, 0.75, 1.25, 1.75 and 3 ml of 0.4 mM Fe2+ solution were added to a series of 25 ml volumetric flasks and mixed with 1.25 ml of 0.5 M H<sub>2</sub>SO<sub>4</sub>, 6.25 ml of buffer solution (with stock solution of 3.82 g of NaC2H5CO2, 0.4 ml of concentrated H<sub>2</sub>SO<sub>4</sub>, diluted to 40 ml with deionized water) and 5 ml of 0.1% 1,10-phenantroline solution (stored in the dark), before dilution to 25 ml with deionized water. The mixture was kept in the dark for about 60 min when full colour development was achieved. The mixture contained concentrations of [Fe(phen)<sub>3</sub>]<sup>2+</sup> ions ranging from 0 to  $4.8 \times 10^{-5}$  M. The molar absorptivity  $\varepsilon$  of [Fe(phen)<sub>3</sub>]<sup>2+</sup> was determined to be  $0.89 \times 10^4 l \, \text{cm}^{-1} \, \text{mol}^{-1}$  by measuring the concentration-dependent absorbance at 510 nm (Supplementary Fig. 4), which is close to the suggested value of  $1.1 \times 10^{4}$  l cm<sup>-1</sup> mol<sup>-1</sup> (ref. <sup>56</sup>). This measured value was used for calculating the photon flux from the white light source.

A 0.006 M Fe<sup>3+</sup> solution was prepared by dissolving 0.05806 g of  $Fe(C_2O_4)_3 \bullet 6H_2O$  in 2 ml of  $H_2SO_4$  (0.5 M) before dilution with deionized water to 20 ml. Then 3 ml ( $V_1$ ) of the solution was illuminated for certain lengths of time, t, (5, 15, 25, 30, 35 s) under efficient magnetic stirring with the white light source power at 204 W (the same power for the photocatalytic activity measurement). A 1 ml volume ( $V_2$ ) of the illuminated solution was added to a 10 ml ( $V_3$ ) volumetric flask containing a mixture of 5 ml of 0.1% 1,10-phenantroline solution (stored in the dark) and 0.5 ml of buffer solution, which was then diluted to a total volume of 10 ml. A reference sample was prepared in the same way except it was not illuminated. The production of Fe<sup>2+</sup> from Fe<sup>3+</sup> can be summarized by the following photochemical reaction:

$$2\left[Fe(C_2O_4)_3\right]^{3-} + h\nu \to 2\left[Fe(C_2O_4)_2\right]^{2-} + 2CO_2 + C_2O_4^{2-}$$
(1)

Where *h* is the Planck constant and  $\nu$  is light frequency. All the solutions were stored in the dark for about 60 min until full colour development was

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achieved. The absorbance difference between the experimental samples and the reference sample was measured at 510 nm (optical path length l=1 cm,  $\epsilon(510 \text{ nm})=0.89\times10^{11} \text{ cm}^{-1} \text{ mol}^{-1})$ . A 0.2 ml volume ( $V_4$ ) of solution was taken from 10 ml ( $V_3$ ) of fully developed solution and then diluted to 3.2 ml ( $V_5$ ) in an optical cuvette to ensure the absorbance A (510 nm) was within the range 0.3–1.1 (Supplementary Fig. 4). The number of moles, n, of Fe<sup>2+</sup> produced by the photochemical reaction is given by

$$n = \frac{\Delta A V_1 V_3 V_5}{\varepsilon l V_2 V_4} \tag{2}$$

Where  $\Delta A$  is the change of absorbance after illumination with time period, t. The photon flux q (einstein  $s^{-1}$ ) entering the sample cell is

$$q = \frac{n}{t\Phi} = \frac{\Delta A V_1 V_3 V_5}{t\Phi \varepsilon l V_2 V_4} \tag{3}$$

The slope of  $\frac{n}{t}$  was measured to be 9.28661×10<sup>-4</sup> mol s<sup>-1</sup> (Supplementary Fig. 4), and since the quantum yield of reaction (1),  $\Phi_{Fe^{2+}}$  (510 nm) = 84.5% (ref. <sup>57</sup>), *q* is calculated to be 1.099×10<sup>-3</sup> einstein s<sup>-1</sup> for the white light source at 204 W. The internal quantum yield (IQY) of methanol production is given by

$$IQY = \frac{6 \times \text{the number of produced methanol molecules}}{qt_{\text{irradiation}}}$$
(4)

When  $t_{\rm irradiation}\,{=}\,120\,{\rm s},$  with 16 mm ol methanol produced (Fig. 1), the IQY is given by

$$IQY = \frac{6 \times 16 \times 10^{-3}}{1.099 \times 10^{-3} \times 120} = 72.8\%$$
(5)

Ultraviolet-visible (UV-vis) absorption was measured using a Perkin-Elmer Lambda 950 UV-vis spectrometer (Supplementary Fig. 21). The wavelength dependences of the internal quantum yield58(the number of electrons for methanol production per absorbed photons) and apparent quantum yield<sup>58</sup> (the number of electrons for methanol production per incident photons) were measured using 0. 01 g of Cu<sub>2</sub>O purified particles in 5 ml of ultrapure deionized H<sub>2</sub>O with continuous CO<sub>2</sub>/H<sub>2</sub>O gas flow. An Xe lamp with the same power (204 W) and distance from the optical source to the sample were used for the photocatalytic reactions with 2 min irradiation time. Illumination was carried out with a 20-cm-long quartz infrared water filter. This cuts off the photons of the infrared region to minimize heating on the sample. Various bandpass filters were applied to select portions of the incident photons. The entire visible light solar spectrum was divided into ten intervals: 400-435 nm, 435-475 nm, 475-500 nm, 500-530 nm, 530-550 nm, 550-570 nm, 570-600 nm, 600-640 nm, 640-680 nm and 680-700 nm. The amount of methanol at certain wavelengths was determined by GC. Wavelength-dependent internal quantum yield and apparent quantum yield are shown in Supplementary Fig. 22. The wavelengthdependent apparent quantum yield correlates well with the wavelength dependence of the photon absorption by Cu<sub>2</sub>O (UV-vis absorption spectrum in Supplementary Fig. 20). On the other hand, the wavelength dependence of the internal quantum yield tells us some useful information about the nature of the photoexcitation process. We find that there are very few losses for the initial condition of reduction of CO2 to methanol with near band-gap excitation (~94% at 620 nm). However, we observe some losses in the high-energy excitation, consistent with the thermalization of photogenerated charges to the band edges. This, in turn, lowers the internal quantum yield to ~50% at 400 nm. The overall internal quantum yield in the visible light region is around 72%. The overall external quantum yield in the visible light region is about 23%. This proves that the reaction is driven by light.

**Solar-to-fuel efficiency calculation from apparent quantum yield data.** The solar-to-fuel efficiency (STF) in the visible light region was calculated from the apparent quantum yield data from Supplementary Fig. 22. This solar energy conversion was evaluated using the same white light source as for the quantum yield measurement with Cu<sub>2</sub>O photocatalyst (0.01 g of photocatalyst in 5 ml of de-ionized water). Illumination was carried out with a 20-cm-long quartz infrared water filter. This cuts off the photons of the infrared region to minimize heating on the sample. Methanol and oxygen were measured as reaction products in the Agilent Headspace sampler as the reactor. The reactor is made of 33–51 coefficient of expansion type I borosilicate glass. A possible chemical formula describing the reaction is the following:

$$\mathrm{CO}_2 + 2\mathrm{H}_2\mathrm{O} \rightarrow \mathrm{CH}_3\mathrm{OH} + \frac{3}{2}\mathrm{O}_2 \tag{6}$$

The Gibbs free energy ( $\Delta G$ ) of this formula to produce CH<sub>3</sub>OH was suggested to be 638.73 kJ mol<sup>-1</sup> (or 6.62 eV per molecule)<sup>39</sup>.

$$STF = \frac{Energy of generated CH_3OH}{Solar energy irriadiating the solution} = \frac{E_{CH_3OH}}{E_{solar}}$$
(7)

$$E_{\rm CH_3OH} = \sum$$
 number of generated CH<sub>3</sub>OH ×  $\Delta$ G (8)

$$\sum_{i} AQY(\lambda_{i}) = \sum_{i} \frac{number \text{ of generated electrons}}{n_{photon,i}}$$
$$= \sum_{i} \frac{6 \times number \text{ of generated CH}_{3}OH}{n_{photon,i}}$$
(9)

Here,  $n_{\text{phonon},i}$  is the number of photons with wavelength  $\lambda$  in interval *i* and AQY( $\lambda_i$ ) is the apparent quantum yield of the *i*th interval in the range.

$$E_{\rm CH_3OH} = \sum_{i} n_{\rm photon,i} \times AQY(\lambda_i) \times \frac{1}{6} \times 6.62 \,\mathrm{eV} \times 1.602 \times 10^{-19} \mathrm{J} \,\mathrm{eV}^{-1}$$
(10)

Thus, we get

$$\text{STF} = \sum_{i} \frac{n_{\text{photon},i} \times \text{AQY}(\lambda_{i}) \times 6.62 \text{ eV} \times 1.602 \times 10^{-19} \text{J eV}^{-1} \times \frac{1}{6}}{E_{\text{solar}}}$$
(11)

In order to derive number of photons per interval *i*,  $n_{\text{photon},i}$ , solar energy irradiating the solution should be divided into intervals  $E_{\mu\nu}$  with  $p_i$  being the fraction of solar power in interval *i*:

$$\sum_{i} p_i \times E_{\text{solar}} = \sum_{i} E_{\lambda_i} \tag{12}$$

$$\sum_{i} E_{\lambda_{i}} = \sum_{i} n_{\text{phonon},i} \times \frac{hc}{\lambda_{i}}$$
(13)

Combining equations (12) and (13), we get

$$n_{\text{phonon,i}} = \frac{p_i \times E_{\text{solar}} \times \lambda_i}{hc} = \frac{p_i \times E_{\text{solar}} \times \lambda_i}{6.626 \times 10^{-34} \times 3 \times 10^8}$$
(14)

Where c is the speed of light. Combining equations (11) and (14), we get

$$\text{STF} = \sum_{i} \frac{p_i \times \lambda_i \times 6.62 \times 1.602 \times 10^{-19} \times \frac{1}{6}}{19.86 \times 10^{-26}} \times \text{AQY}(\lambda_i)$$
(15)

$$STF = \sum_{i} 0.89 \times 10^{6} \times p_{i} \times \lambda_{i} \times AQY(\lambda_{i})$$
(16)

Since the photons of the infrared region were cut off by a 20-cm-long quartz infrared water filter, the visible region of the solar spectrum was divided into ten intervals (i=1, 2, 3, ..., 10). The fraction of solar power in interval *i* was calculated by dividing the area of interval *i* of the spectrum by the entire area of the solar spectrum distribution of the light source. The calculated solar-to-fuel efficiency is about 10.1% as shownsc in Supplementary Table 1. This is less than the theoretical prediction maximum of 35.2% for methanol<sup>39</sup>.

**Operando multimodal single-particle microscopy and nanospectroscopy.** We used a custom gas-flow nanoreactor (Hummingbird Scientific) to examine single  $Cu_2O$  particles under working conditions by in situ ETEM and operando SFXM. The  $Cu_2O$  particles were sandwiched inside the nanoreactor between 50 nm  $Si_3N_4$  membranes through which  $CO_2$  and  $H_2O$  gas were flowed at atmospheric pressure (Fig. 2a). The optical transparency of the  $Si_3N_4$  membranes facilitated excitation of the photocatalyst with a laser whilst probing the structure and chemical state of the  $Cu_2O$  particles (Fig. 2a and Supplementary Fig. 9). This enabled us to study the same photocatalytic particle under working conditions by both in situ ETEM and operando SFXM (Supplementary Figs 6, 12). Nickel grid bars on the bottom membrane served as fiducial marks to identify the same particle in the nanoreactor with both techniques (Fig. 2c).

Cu<sub>2</sub>O particles were dispersed in ethanol and drop-casted onto the bottom Si chip with a Si<sub>3</sub>N<sub>4</sub> membrane and nickel grid bars. The bottom Si chip was assembled with a top Si chip, which also had a Si<sub>3</sub>N<sub>4</sub> membrane, and separated by spacers. Samples were maintained in the nanoreactor under stable environmental conditions during transfer between the ETEM and SFXM platforms. For the in situ ETEM and operando SFXM measurements, CO2 gas (Airgas, Inc.) and H<sub>2</sub>O vapour gas (500 ppm mol percentage balance with N<sub>2</sub> (Airgas, Inc.), 500 mol H<sub>2</sub>O molecules per 1 million moles of N<sub>2</sub> molecules) were mixed and delivered to the nanoreactor at 0.05 sccm using a gas-flow control box (Hummingbird Scientific). The CO<sub>2</sub> to H<sub>2</sub>O (gas) pressure ratio was about 1:1. ETEM imaging and EELS were performed on a JEOL 2100F electron microscope operated at an accelerating voltage of 200 kV and equipped with a Gatan imaging filter system with an energy dispersion of 0.05 eV per pixel at the Center for Nanoscale Materials (CNM) at Argonne National Laboratory. EELS spectra were recorded in energy-filter TEM mode with an energy resolution of 0.7 eV, a collimated beam and 1 nm probe size. Exposure of the sample to the electron beam was minimized

to avoid damage. Gas flow was started with closed beam valve for 30 min before taking EELS spectra with the gas mixture. An EELS acquisition time of 5 s and zero-loss peak calibration were performed between each acquisition to prevent zero-loss peak drift from microscope instabilities by the introduction of the gas. The threshold dose rate for electron beam damage on metal oxide was suggested to be  $3.125 \times 10^3 \text{ e}^{-s^{-1}} \text{ Å}^{-2}$  (or  $5 \text{ A cm}^{-2}$ )<sup>59</sup>, while the dose rate used in this experiment was about  $1.5 \times 10^2 e^{-s^{-1}} \text{Å}^{-2}$ . This is about an order of magnitude smaller than the threshold of beam damage for metal oxide. This published threshold is in agreement with our experimental observation that no obvious structural damage occurred. SFXM was performed using the Hard X-ray Nanoprobe60 operated by the CNM at beamline 26-ID at the Advanced Photon Source (APS) at Argonne National Laboratory. A Fresnel zone plate lens with a 20 nm finest zone focused a coherent, monochromatic ~9 keV X-ray beam to a ~30 nm spot with an intensity of 109 photons s<sup>-1</sup>. A four-element silicon drift detector (Vortex ME4) collected the emitted characteristic fluorescence X-rays. The sample was illuminated with either a green (532 nm) or blue (445 nm) laser to induce photocatalysis. The X-ray microscopy experiment setup is shown in Supplementary Fig. 9. TEM measurements made on the same particle afterwards show the particle was stable during the X-ray measurements (Supplementary Fig. 11). Reference X-ray absorption spectroscopy data were taken with Cu, Cu<sub>2</sub>O and CuO powder standards at APS beamline 20-ID-B for comparison (Supplementary Fig. 11).

**Operando high-energy X-ray diffraction characterization.** We performed operando high-energy X-ray diffraction characterization during the  $CO_2$  adsorption under dark and illuminated conditions at APS beamline 11-ID-C. A Si(311) single-crystal monochromator, with X-ray beam energy of 105.1 keV, beam size of  $0.2 \times 0.2$  mm<sup>2</sup> and a Perkin-Elmer area detector located at 1.8 m were used to obtain 2D diffraction patterns from the sample. The Cu<sub>2</sub>O ensemble was placed in a glass pipette partially sealed with glass fibres, which was connected to the same gas-delivering system used for in situ single-particle ETEM and operando SFXM. Then the photocatalyst was irradiated using the Xe lamp with continuous  $CO_2/H_2O$  gas flow during the measurement. Diffraction patterns were calibrated using a standard CeO<sub>2</sub> sample, converted to one-dimensional patterns using the Fit2D software package<sup>61</sup>, then analysed using the ranked data analysis software package<sup>62</sup>.

**Ex situ high-resolution synchrotron X-ray powder diffraction.** Synchrotron HRXRD data were collected at APS beamline 11-BM-B, using a wavelength of 0.414230 Å for the CO<sub>2</sub>/H<sub>2</sub>O treated samples (Fig. 5 and Supplementary Fig. 17) and a wavelength of 0.414221 Å for the H<sub>2</sub>O treated samples (Supplementary Fig. 18). All samples were sealed by Varian Torr Seal epoxy in Kapton capillaries inside a N<sub>2</sub> gas-flow glove box. All diffraction patterns were collected at room temperature. Discrete patterns covering an angular range of 4–12° were scanned over a 25° range in 2 $\theta$ , with data points collected every 0.001° using a scan speed of 0.1 s per step.

**Ex situ structural characterization.** SEM images of the photocatalyst particles were taken with a JEOL 7500 scanning electron microscope with an acceleration voltage of 10 kV. Transmission electron microscope and high-resolution TEM images (Supplementary Fig. 23), NBD and EELS measurements were performed with a JEOL 2100F electron microscope operated at an accelerating voltage of 200 kV and equipped with a Gatan imaging filter system at the CNM. Ex situ TEM characterization was performed by drop casting a suspension of the photocatalyst particles in ethanol onto a Formvar-coated Au TEM grid.

**EPR characterization.** Continuous-wave EPR measurements were carried out using a Bruker ELEXYS E580 spectrometer operating in the X-band (9.4 GHz) mode and equipped with an Oxford CF935 helium flow cryostat with an ITC-5025 temperature controller at CNM. The *g* tensor was calibrated for homogeneity and accuracy by comparison to a standard 1,3 Bisdiphenilene 2 phenylallyl benzene (BDPA-benzene) complex, with *g* = 2.00254 ±0.00005. The receiver gain and number of scans were adjusted for each spectrum of a particular sample to enable comparison at the same signal-to-noise ratio. The measurements were carried out at 4K. For the studies of charge transfer, samples were degassed by bubbling in a N<sub>2</sub> flow glove box, frozen at liquid He temperature and illuminated in the EPR cavity using the 300 W Xe lamp with a plano-convex lens (50 cm focal length) with a broadband antireflection coating and water as an infrared cutoff filter.

We used 0.01 g of Cu<sub>2</sub>O in the EPR experiment. The number of Cu<sub>2</sub>O molecules is  $n_e$ .

$$n_a = \frac{m}{M} N_{\rm A} = \frac{0.01\,\rm g}{143\,\rm g\,mol^{-1}} \times 6.02 \times 10^{23} = 4.2 \times 10^{19}$$
(17)

Where  $N_A$  is Avogadro's number. The number of spins from the Cu cluster was determined to be  $3 \times 10^{12}$  (~4.98×10<sup>-12</sup> mol) by comparing the signal to the calibrated weak pitch sample under the same conditions. Assuming that the Cu clusters are produced from the Cu(I) disproportionation reaction under illumination according to the reaction:

$$Cu_2O = Cu^0 + Cu^{2+} + \frac{1}{2}O_2 \tag{18}$$

The number of O<sub>2</sub> molecules produced by disproportionation,  $n_b = 1.5 \times 10^{12}$  O<sub>2</sub> molecules.

With the same amount of illumination, we measured 22  $\mu \rm{mol}$  of  $\rm{O}_2.$  The number of  $\rm{O}_2$  molecules is  $n_c$ 

$$n_c = nN_{\rm A} = 2.2 \times 10^{-5} \times 6.02 \times 10^{23} = 1.32 \times 10^{19}$$
(19)

Since  $n_b$  is seven orders of magnitude less than  $n_c$ , O<sub>2</sub> measured under illumination does not only originate from Cu(1) disproportionation.

DFT calculations. To study the stability of both the pristine surfaces and adsorption of CO<sub>2</sub> and H<sub>2</sub>O on surfaces of the Cu<sub>2</sub>O particles, we performed DFT calculations with plane-wave basis sets as implemented in the VASP code63,64. All the calculations were spin-polarized and carried out using the gradient corrected exchange-correlation functional of Perdew, Burke and Ernzerhof<sup>85</sup> under the projector augmented wave method, with plane-wave basis set up for a kinetic energy cutoff of 500 eV. The projector augmented wave method66 was used to represent the interaction between the core and valence electrons, and the Kohn-Sham valence states (that is 1s for H, 2s2p for C and O, 3d4s for Cu) are expanded in plane-wave basis sets. For both the Cu<sub>2</sub>O bulk crystal and surface calculations, the GGA + U calculations were followed, with the Hubbard U value of Cu atoms chosen to be U=6.0 eV, according to the reported Cu<sub>2</sub>O compounds in the literature<sup>67</sup>. For the geometry optimization, the convergence criterion of the total energy was set to be within  $1 \times 10^{-5}$  eV for the K-point integration, and all the atoms and geometries were optimized until the residual forces became less than  $1 \times 10^{-2}$  eV Å<sup>-1</sup>. For the Cu<sub>2</sub>O bulk crystal, the calculation was based on a mesh of  $8 \times 8 \times 8$  in the K-point grid. For both the crystalline Cu<sub>2</sub>O surfaces (that is, (110) and (100) surfaces) and adsorption calculations, a K-point grid of 2×2×1 was used. For the surface slab calculations, we adopted a slab/vacuum geometry composed of repeating slabs and vacuum layers. In the supercell of the surface slab calculations, the vacuum thickness was ~18 Å and the thickness of the surface slab was ~22 Å, over which the bottom three layers were fixed in the crystalline bulk position during the geometry optimization. Calculated conformations of linear and bent CO<sub>2</sub> (that is, activated CO<sub>2</sub>)68 on Cu termination of Cu<sub>2</sub>O (110) with the co-adsorption of H<sub>2</sub>O are shown in Supplementary Fig. 14.

**Photoreduction of metal salts on Cu<sub>2</sub>O.** The photoreduction of metal ions on the surface of particles was carried out using HAuCl<sub>4</sub> as precursor and H<sub>2</sub>O as hole scavenger. Specifically, 5 mg of Cu<sub>2</sub>O particles were mixed with 10 mg of HAuCl<sub>4</sub> in 10 ml of de-ionized H<sub>2</sub>O. The suspension was irradiated with a 300 W Xe lamp for 2 min with continuous stirring. Then the suspension was filtered and washed with de-ionized H<sub>2</sub>O and dried in a glove box for 48 h. The Au nanoparticles were deposited randomly on Cu<sub>2</sub>O without preferable facets (Supplementary Fig. 20).

#### Data availability

The data that support the plots within this paper and the findings of this study are available from the corresponding authors upon reasonable request.

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## Author contributions

Y.A.W., I.M., Y.L. and T.R., conceived the ideas. Y.A.W. performed synthesis of samples, ex situ structural characterization, photocatalytic activity characterization, solar-tofuel efficiency measurements and quantum yield measurements. Y.A.W., Z.C., J.R.G., Y.L. and I.M. performed operando SFXM measurements. Y.A.W. and Y.L. performed operando ETEM, electron diffraction and EELS measurements. Y.A.W. and T.R. performed EPR and isotope tracer measurements. Y.A.W., Q.L. and Y.R. performed operando high-energy X-ray diffraction and ex situ high-resolution X-ray powder diffraction measurements. Y.A.W. and C.-J.S. performed X-ray absorption spectroscopy measurements. A.P.P. and V.S. performed XPS measurements. C.L., K.C.L. and L.A.C. performed DFT computations. Y.A.W, IM. and T.R. wrote the paper. All authors contributed to the discussion and revision of the paper.

#### **Competing interests**

The authors declare no competing interests.

#### Additional information

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