

**Supplemental Information:**  
**Cooperative spin transition of monodispersed FeN<sub>3</sub> sites within  
graphene induced by CO adsorption**

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## Computational Details

Spin-polarized Density functional theory (DFT) simulations were implemented in Vienna Ab-initio Simulation Package (VASP)[1, 2] with generalized gradient approximation (GGA) of Perdew-Burke-Ernzerhof (PBE) exchange correlation functional.[3] Strong Coulomb interactions of 3d electron of transition metal were described with DFT+U approach[4, 5], in which Hubbard U and on-site exchange constant J were 4 and 1 eV, respectively.[6] These values were previously reported to demonstrate accurately the spin state in transition metal chemistry.[7, 8] The interaction between core electrons and valence electrons was represented by projected augmented wave (PAW) method[9, 10] and cutoff energy for the plane wave basis set is 500 eV. The van der Waals interaction was considered by DFT-D2 correction of Grimme scheme.[11] The force and energy convergence criterion were set to 0.01 eV/Å and  $10^{-6}$  eV, respectively. Two single Fe atoms were anchored onto  $8 \times 8$  graphene monolayer separated by a vacuum layer of 18 Å and k-point sampling of first Brillouin zone was used with a  $3 \times 3 \times 1$   $\Gamma$ -centered grid for geometry optimizations. [12]

More specifically, in our calculation, to find the magnetic ordering of Fe atoms in magnetic ground state during the stepwise CO and O<sub>2</sub> adsorption, we have considered NM, FM, AFM ordering with different initial value on Fe atoms. We have tried as many initial values of spin moment as possible to find the lowest energy. Then we compare the corresponding total energies of these magnetic states and the lowest one is deemed as the magnetic ground state. Such calculation is widely adopted by theoretical community to explore the most stable magnetic ordering and related magnetic properties in ferromagnetism and ferroelectricity fields, which provides reasonable agreement with experimental findings.[13, 14, 15, 16, 17, 18]

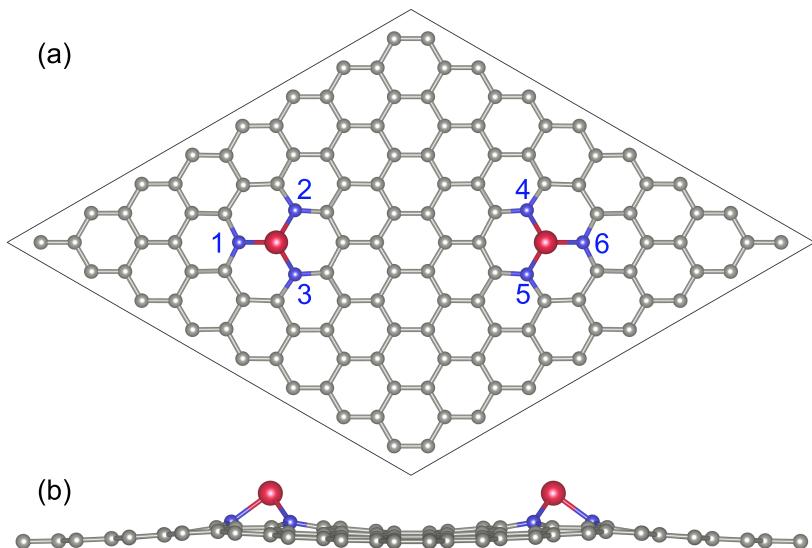


FIG. S1: Top view (a) and side view (b) of the optimized structure of FeN<sub>3</sub>-FeN<sub>3</sub>. The labeled numbers differentiate the N atoms for the purpose of comparing the changes of Fe-N bond length.

Table S1: The geometrical and electronic transformation of FeN<sub>3</sub>-FeN<sub>3</sub> upon various CO adsorption concentration

CO Molecules <sup>a</sup>	Energy <sup>b</sup> (eV)	Spin Moment ( $\mu_B$ )			Fe-N Bond Length <sup>f</sup> (Å)	Fe-C (CO) Bond Length <sup>g</sup> (Å)
		Total <sup>c</sup>	Left <sup>d</sup>	Right <sup>e</sup>		
0	0.00	8.47	3.41	3.41	1.955,1.965,1.965,1.964,1.964,1.954	1.143
1	-1.86	-1.12	2.81	-3.39	1.995,2.053,2.053,1.957,1.945,1.957	1.174
2	-1.70	0.00	2.89	-2.89	1.983,2.034,2.034,2.035,2.035,1.983	1.171,1.171
3	-1.23	4.36	1.03	2.96	1.977,2.035,1.978,2.018,2.018,1.974	1.169,1.168,1.168
4	-1.20	0.00	1.02	-1.03	1.977,2.034,1.977,2.035,1.977,1.977	1.168,1.168,1.168,1.168

<sup>a</sup> n (n= 0,1,2,3,4) CO molecules is adsorbed onto the Fe active sites orderly, n = 0 represents the bare FeN<sub>3</sub>-FeN<sub>3</sub>;

<sup>b</sup> Taking the energy of FeN<sub>3</sub>-FeN<sub>3</sub> as reference energy, the energy change ( $\Delta E$ ) of each CO adsorption,  $\Delta E = E_{(n+1)CO@FeN_3-FeN_3} - E_{CO} - E_{nCO@FeN_3-FeN_3}$ ;

<sup>c</sup> Total spin moment of FeN<sub>3</sub>-FeN<sub>3</sub> framework shown in the FIG. S1;

<sup>d</sup> Spin moment of the left Fe atom shown in the FIG. S1;

<sup>e</sup> Spin moment of the right Fe atom shown in the FIG. S1;

<sup>f</sup> The Fe-N<sub>m</sub> bond length in the FeN<sub>3</sub>-FeN<sub>3</sub> active site ordered by increasing m values (m = 1,2,3,4,5,6) shown in the FIG. S1;

<sup>g</sup> The Fe-C (CO) bond length is arranged in the sequence of CO adsorption;

Table S2: The geometrical and electronic transformation of FeN<sub>3</sub>-FeN<sub>3</sub> upon various O<sub>2</sub> adsorption concentration

O <sub>2</sub> Molecules	Energy (eV)	Spin Moment ( $\mu_B$ )			Fe-N Bond Length (Å)
		Total	Left	Right	
0	0.00	8.47	3.41	3.41	1.955,1.965,1.965,1.964,1.964,1.954
1	-2.77	1.09	-2.74	3.39	1.955,1.954,2.035,1.950,1.957,1.951
2	-2.59	0.00	2.76	-2.76	1.949,2.025,1.949,2.025,1.949,1.949

Table S3: the calculated spin moment of adjacent Fe atoms by HSE06 functional

CO molecules	Spin Moment ( $\mu_B$ )		
	Total	Left	Right
0	8.47	3.40	3.40
1	2.95	-1.02	3.38
2	0.01	2.95	-2.95
3	-2.34	1.06	-3.05
4	0.00	1.06	-1.06

Compared to PBE+U results, there exists some inconsistency of spin moment on individual Fe atoms in the cases of 1 CO and 3 CO adsorption. Since HSE06 method is quite time-consuming in periodic DFT calculations, an exhaustive search of all possible spin configurations for these two would be computationally prohibitive. Nevertheless, the finished HSE06 results clearly show cooperative communication between two adjacent Fe sites as reflected by continuous cooperative spin transitions of two neighboring Fe sites upon CO adsorption, which is in line with PBE +U results. In addition, we have compared the FM and AFM spin alignment of Fe atoms at FeN3-FeN3@graphene and found that FM state is 9.1 meV more stable than AFM state at HSE06 level, which agrees well with PBE+U calculations (main text, Page 2, Line 15-20). Furthermore, for the stepwise 2 and 4 CO molecules adsorption, the most stable spin alignment on adjacent Fe atoms is AFM state, agreeing with PBE+U calculations. In short, we have compared HSE06 and PBE+U results and argue that our main conclusion about the CO-induced spin transition on adjacent Fe atoms holds.

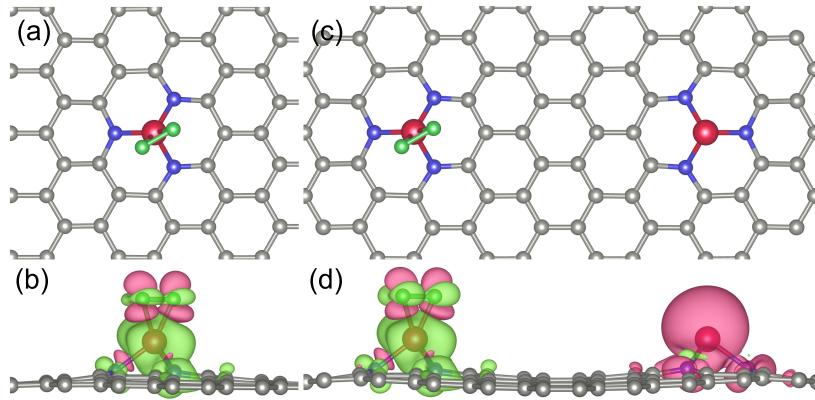


FIG. S2: Top view of the optimized structure of FeN<sub>3</sub> (a) and FeN<sub>3</sub>-FeN<sub>3</sub> (c) after first O<sub>2</sub> molecule is adsorbed, (b) and (d) is the side view of net spin electron density of FeN<sub>3</sub> and FeN<sub>3</sub>-FeN<sub>3</sub> at an isosurface value of  $1 \times 10^{-3} \text{ e}\text{\AA}^{-3}$ .

## Supplementary Note 1

We have also considered the O<sub>2</sub> molecules adsorption on the Fe sites of FeN<sub>3</sub>-FeN<sub>3</sub> as shown in FIG. S2. O<sub>2</sub> molecule bonds to the Fe atom in side-on fashion and the O-O bond length is elongated by 0.20 Å whereas that is 0.18 Å in isolated FeN<sub>3</sub> active site. The calculated adsorption energy is 2.77 eV, 0.11 eV higher than that in FeN<sub>3</sub> system (2.66 eV). These indicate considerable adsorption enhancement if two active sites approach mutually, compared to single Fe site. Upon the O<sub>2</sub> adsorption, the geometrical structure of FeN<sub>3</sub>-FeN<sub>3</sub> adjusted accordingly and all the geometrical symmetry disappeared. As shown in Table S2, the adsorbed O<sub>2</sub> alters not only the Fe-N bond length of the corresponding Fe site, but also that of neighbouring Fe site. Such associated structural transformation upon molecule adsorption cannot be observed in single FeN<sub>3</sub> system and it's indeed the indication of the existence of communicative effect between adjacent Fe sites of FeN<sub>3</sub>-FeN<sub>3</sub>@graphene.

More interestingly, one can see from FIG. S2(d), due to O<sub>2</sub> insertion, the spin moment of Fe atom is decreased to  $2.74 \mu_B$ , the net spin charge localized on Fe atom is reduced with part of them transferred to the adsorbed O<sub>2</sub>. In contrast with CO adsorption, the magnetic interaction between O<sub>2</sub> and Fe center reverses the spin moment of Fe site it attached whereas CO reverses that of adjacent site. Although O<sub>2</sub> adsorption induces comparable structural transformation of Fe active center, it only leads to slight spin moment reduction on the nearby Fe site. As a result, the most energy-favorable spin ordering of FeN<sub>3</sub>-FeN<sub>3</sub> shifts from ferromagnetic to ferrimagnetic with total spin moment of  $1.12 \mu_B$ . These two changes of the neighbouring Fe site signify that the interconnectedness of adjacent Fe centers exists and they communicate and interact indirectly via the graphene matrix, hinting that the mutual influence between the adjacent active sites need to be considered in analogous SAC frameworks. The spin transition upon the O<sub>2</sub> adsorption verifies that O<sub>2</sub> adsorption has, similarly to CO, led to 2D long-range ferrimagnetic orders on Fe sites, and the cooperative interaction between nearby Fe sites mediated by O<sub>2</sub> molecule can tailor the magnetic ordering of 2D material.

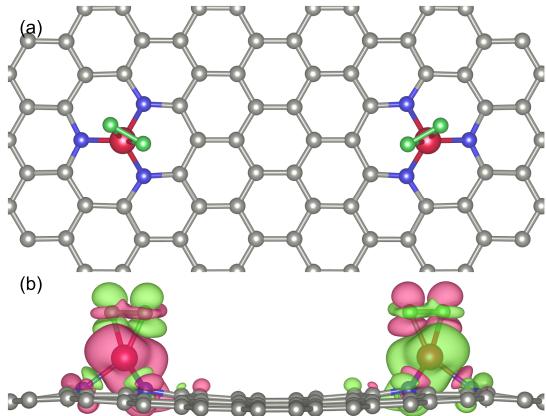


FIG. S3: (a) Top view of the optimized structure of FeN<sub>3</sub>-FeN<sub>3</sub> after two O<sub>2</sub> molecules are adsorbed on the two individual Fe sites, (b) the side view of net spin electron density at an isosurface value of  $1 \times 10^{-3} \text{ e}\text{\AA}^{-3}$ .

## Supplementary Note 2

We further consider a second O<sub>2</sub> molecule adsorption on the neighbouring Fe site. The optimized structure after the second O<sub>2</sub> adsorption is shown in FIG. S3(a). The calculated O<sub>2</sub> binding energy is 2.59 eV, 0.18 eV weakening than the first O<sub>2</sub> adsorption. Upon the second O<sub>2</sub> adsorption, the Fe-N bond length readjusted slightly and the overall geometric symmetry (C<sub>2v</sub>) reappeared summarized in Table S2. We further explore the response of spin state to such geometric adjustment. In the ground state after two O<sub>2</sub> molecules interact respectively with the Fe sites, the spin moment of corresponding Fe sites decreases from 3.39 to 2.76  $\mu_B$  whereas that of the neighbouring site with preadsorbed O<sub>2</sub> arises a little by 0.02  $\mu_B$ . The net spin charge is still mainly localized on the Fe sites and adsorbed O<sub>2</sub> molecules if seeing these two sites separately, much resembling to the O<sub>2</sub> adsorption in the single FeN<sub>3</sub> structure. However, if combined, one can see that the net spin moment of each individual Fe site is in the equal value (2.76  $\mu_B$ ) but antiparallel, unambiguously indicating the 2D long-range antiferromagnetic (AFM) orders on Fe sites. This AFM spin ordering is completely different from the FM spin ordering when the Fe active sites are free from O<sub>2</sub> adsorption, although both of the structure are in the same C<sub>2v</sub> symmetry. Meanwhile, one can see clearly in FIG. S3(b), the second adsorbed O<sub>2</sub> reverses not only the spin moment of Fe center it bonds, but also that of its nearby Fe site, which is also different from the second CO adsorption.

This synchronous spin transition of adjacent active sites again indicates inter-site cooperative communication arises when responding to the exterior electronic perturbation induced by CO or O<sub>2</sub> adsorption, if they approach to each other closely to deliver such perturbation effectively. More specifically, it further demonstrates that the communicative effect between two adjacent active sites can be tailored through external surface-molecular spin interaction and the desirable magnetic ordering on the Fe sites can be fine-tuned by controlling the O<sub>2</sub> adsorption.

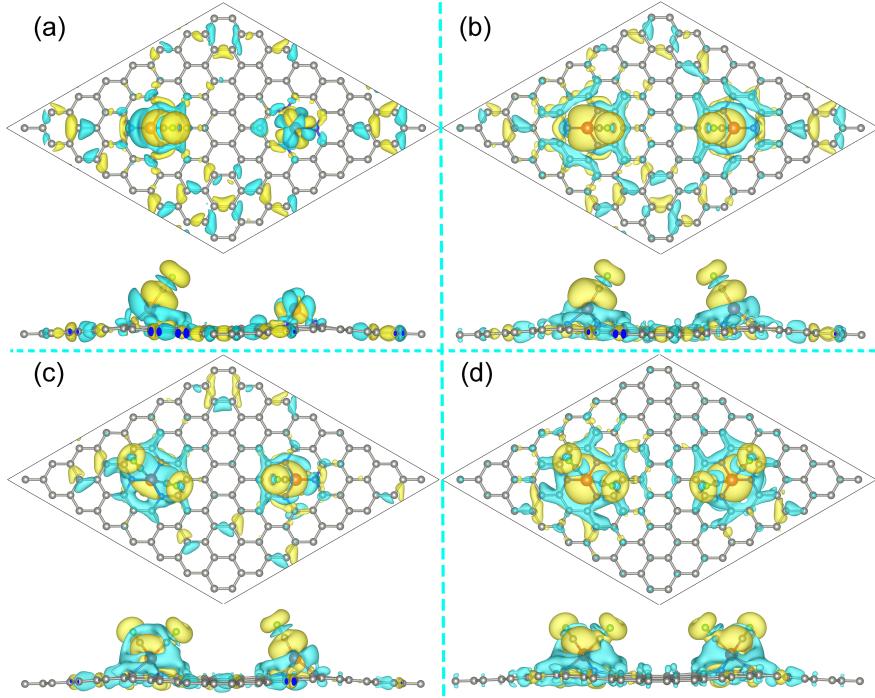


FIG. S4: (a) Top view and side view of charge density difference when one CO molecule adsorbed on the Fe atom of FeN<sub>3</sub>-FeN<sub>3</sub>@graphene. (b),(c) and (d) is for two, three and four CO molecules adsorbed on Fe atoms, respectively. The benchmark is the charge density when no CO molecule is adsorbed on Fe atoms. The isosurface value is  $5 \times 10^{-4} \text{ e}\text{\AA}^{-3}$ . Yellow and blue bubbles represent charge accumulation and depletion, respectively.

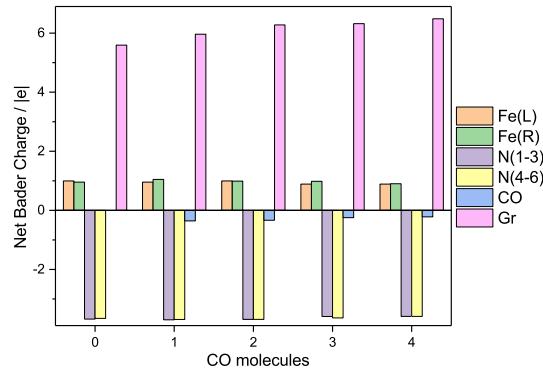


FIG. S5: Net bader charge of each section of the active sites, *i.e.* Fe atoms and three coordinated N atoms as a whole, and graphene matrix and CO, upon stepwise CO adsorption. The Fe and N atom are labeled as the Fig. S1 and the bader charge change of adsorbed CO is averaged.

### Supplementary Note 3

In the Fig. S4, it indicates the inter-site charge redistribution of FeN<sub>3</sub>-FeN<sub>3</sub>@graphene framework upon CO adsorption, which is expected to cause the modification of charge-carrier concentration on graphene matrix hence the inter-site exchange coupling strength. Consequently, adjacent Fe atoms establish characteristic long-range spin ordering.

Clearly, once CO molecule is adsorbed on one Fe active sites, it leads to significant charge transfer in the corresponding site; Further, one can also see the adjoint charge transfer in its adjacent Fe site despite of no direct interaction from the injected CO molecule. Such communicative charge behaviors is the clear evidence of cooperative interaction between active sites. Additionally, the surface electron of graphene also experiences considerable depletion, which, as a bridge, interconnects the

adjacent sites, delivers the electronic perturbation induced by CO from site to site and further mediates effectively the indirect exchange coupling between two adjacent Fe active sites.

Bader charge analysis[19] in the Fig. S5 also indicates the charge transfer of the Fe active sites induced by CO adsorption. It reveals that the Fe atoms still positively charged and three N atoms negatively charged, which facilitates the charge redistribution on the Fe active sites and graphene matrix upon CO adsorption.

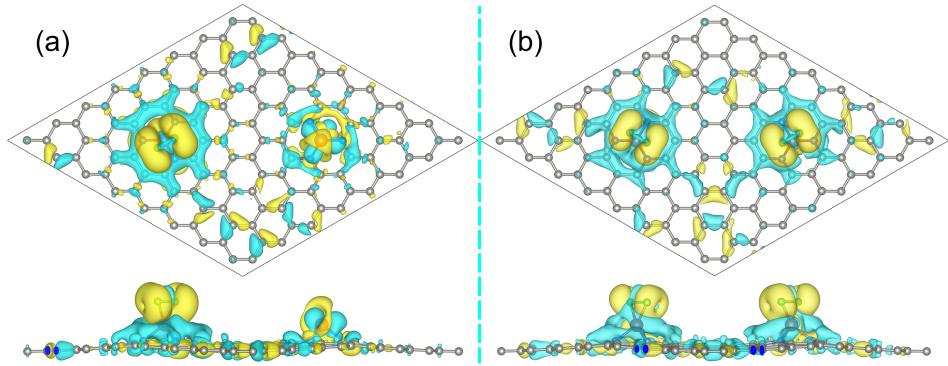


FIG. S6: Top view and side view of charge density difference when one (a) and two (b)  $\text{O}_2$  molecule adsorbed on Fe active sites of  $\text{FeN}_3\text{-FeN}_3@\text{graphene}$ , respectively. The benchmark is the charge density when no  $\text{O}_2$  molecule is adsorbed on Fe atoms. The isosurface value is  $5 \times 10^{-4} \text{ e}\text{\AA}^{-3}$ . Yellow and blue bubbles represent charge accumulation and depletion, respectively.

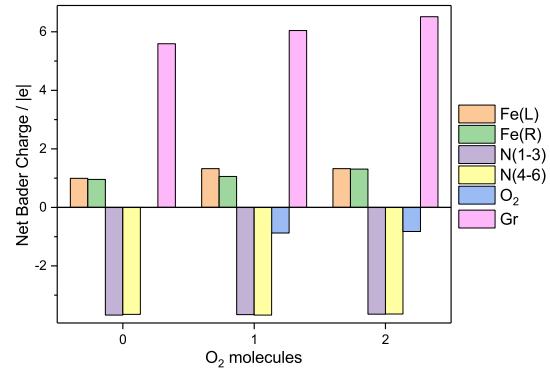


FIG. S7: Net bader charge of each section of the active sites, *i.e.* Fe atoms and three coordinated N atoms as a whole, and graphene matrix and  $\text{O}_2$ , upon stepwise  $\text{O}_2$  adsorption. The Fe and N atom are labeled as the Fig. S1 and the bader charge change of adsorbed  $\text{O}_2$  is averaged.

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